## Summary of Chapter 5

(When do we have orthogonal eigenfunctions for our boundary value problem?)

Key: A Sturm-Lioville problem has orthogonal eigenfunctions

Sturm-Liouville (eigenvalue) problem:

$$\frac{d}{dx}\left[P(x)\frac{du}{dx}\right] + Q(x)u - \lambda R(x)u = 0 , \qquad (1)$$

for u(x) defined on  $x \in [a, b]$ , plus homogeneous b.c.'s

A 
$$u(a) + B u'(a) = 0$$
 (u'  $\equiv du/dx$ ) (I)  
C  $u(b) + D u'(b) = 0$ . (II)

Remarks:

- (1) The forms of the ODE and b.c.'s above are general enough that many physical problems can be converted to a standard Sturm-Liouville problem => Orthogonality of eigenfunctions
- (2) It is crucial that the b.c.'s are homogeneous. If they are not, there may not be orthogonal eigenfunctions for the system.

Proof of orthogonality...

**Step 1**: Define the operator *L* as

$$\boldsymbol{L}\{\boldsymbol{u}\} \equiv \frac{d}{dx} \left[ P(\boldsymbol{x}) \frac{d\boldsymbol{u}}{d\boldsymbol{x}} \right] + Q(\boldsymbol{x})\boldsymbol{u} ,$$

such that the original Eq. (1) in the Sturm-Liouville system can be written as

$$L\{u\} = \lambda R(x) u \tag{2}$$

Let *u* and *v* be two solutions (need not be eigenfunctions at this point) to the Sturm-Liouville problem, then

$$\int_{a}^{b} v L\{u\} - u L\{v\} dx = \int_{a}^{b} v\{\frac{d}{dx} [P\frac{du}{dx}] + Qu\} - u\{\frac{d}{dx} [P\frac{dv}{dx}] + Qv\} dx$$

$$= \int_{a}^{b} v \frac{d}{dx} [P\frac{du}{dx}] - u \frac{d}{dx} [P\frac{dv}{dx}] dx$$

$$= \int_{a}^{b} \frac{d}{dx} [P(v\frac{du}{dx} - u\frac{dv}{dx})] dx$$

$$= \left[ P(v\frac{du}{dx} - u\frac{dv}{dx}) \right]_{a}^{b}$$

$$= P(b)[v(b)u'(b) - u(b)v'(b)] - P(a)[v(a)u'(a) - u(a)v'(a)]$$

$$= 0$$
(3)

See next page for an explanation why the green-colored expression is identically zero.

(Addendum to the derivation in previous page)

Since *u* and *v* are two solutions to the Sturm-Liouville system, they both satisfy the b.c.'s (I) and (II),

A u(a) + B u'(a) = 0(I-u)A v(a) + B v'(a) = 0(I-v)C u(b) + D u'(b) = 0.(II-u)C v(b) + D v'(b) = 0.(II-v)

From (I-u) and (I-v), we have

$$u(a) = (-B/A) u'(a)$$
, and  $v'(a) = (-A/B) v(a) \implies v(a)u'(a) - u(a)v'(a) = 0$ 

similarly, from (II-u) and (II-v) we can establish that v(b)u'(b) - u(b)v'(b) = 0.

**Step 2**: Now, consider that *u* and *v* are two eigenfunctions,  $u = \phi_m$ ,  $v = \phi_n$ , of the Sturm-Liouville problem corresponding to eigenvalues  $\lambda_m$  and  $\lambda_n$ . Then,

$$L\{\phi_m\} = \lambda_m R(x) \phi_m$$
,  $L\{\phi_n\} = \lambda_n R(x) \phi_n$ .

Using Eq. (3), we have

$$0 = \int_a^b \phi_n \boldsymbol{L} \{\phi_m\} - \phi_m \boldsymbol{L} \{\phi_n\} dx = \int_a^b (\lambda_m - \lambda_n) \phi_m \phi_n \boldsymbol{R}(x) dx .$$

Therefore, as long as  $\lambda_m \neq \lambda_n$  , we have the orthogonality relation

$$\int_{a}^{b} \phi_{m} \phi_{n} R(x) dx = 0 \quad .$$

## **Remarks:**

In addition to orthogonality of eigenfunctions, it can be shown that

- The eigenvalues of the Sturm-Liouville system are <u>discrete</u> and <u>real</u>, and they have a lower bound (but no upper bound); The eigenvalues can be ordered as  $\lambda_1 < \lambda_2 < \lambda_3 < ...$ , with  $\lambda_1$  the smallest eigenvalue
- There is a one-to-one correspondence between an eigenvalue and an eigenfunction
- The eigenfunctions form a <u>complete basis</u> for piece-wise continuous functions defined on [*a*, *b*], meaning that any function *f*(*x*) that is piece-wise continuous can be represented by the eigenfunction expansion,

$$f(x) \approx \sum_{n} a_{n} \phi_{n}$$

See p. 163 in textbook for further detail.