

Lecture 12

2/20

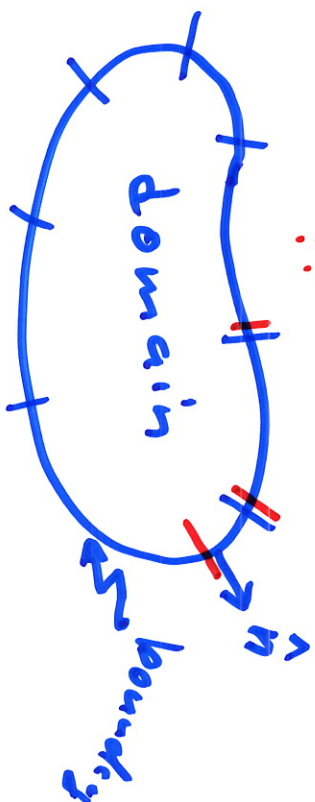
* HW2 posted, due 3/5 (Sunday)

Recap (Lec 11)

End-to-end sol. for Laplace's eq. + Dirichlet b.c. sol. exists and is unique
impose z_c on the boundary

near the end of Lec 11:

Cannot have unique sol. Laplace's eq. + Neumann b.c.

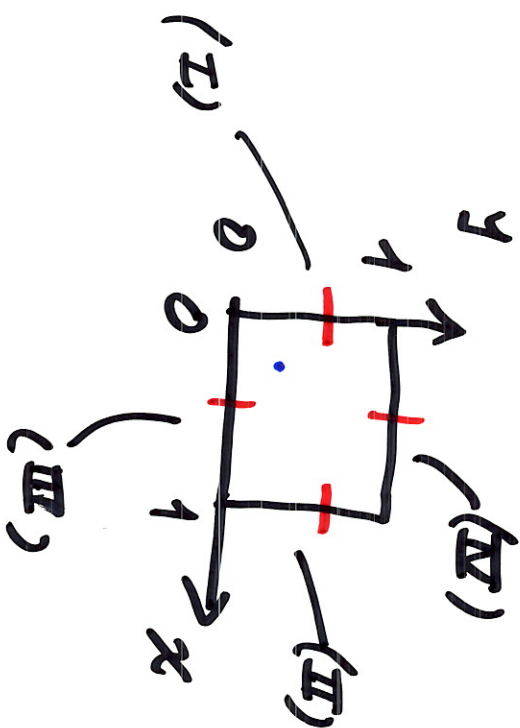


impose derivative of u on the boundary
in the direction normal to boundary

End of Lec 11

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0 \quad (*)$$

- (I) $u_x(0, y) = F_1(y)$
- (II) $u_x(1, y) = F_2(y)$
- (III) $u_y(x, 0) = F_3(x)$
- (IV) $u_y(x, 1) = F_4(x)$



★ has

either infinitely many sols.

(or no solution) ← ?

(if) Say, we find a

solution of ★ $u(x, y)$

Then $\hat{u}(x, y) = u(x, y) + C$ is also a solution of ★

any const!

$\int_0^1 \int_0^1 \psi(x) dx dy$ Area integral of $\psi(x)$

$$\int_0^1 \int_0^1 \left\{ \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right\} dx dy = \int_0^1 \int_0^1 0 dx dy = 0$$

$$\int_0^1 \left[\int_0^1 \frac{\partial^2 u}{\partial x^2} dx \right] dy + \int_0^1 \left[\int_0^1 \frac{\partial^2 u}{\partial y^2} dy \right] dx$$

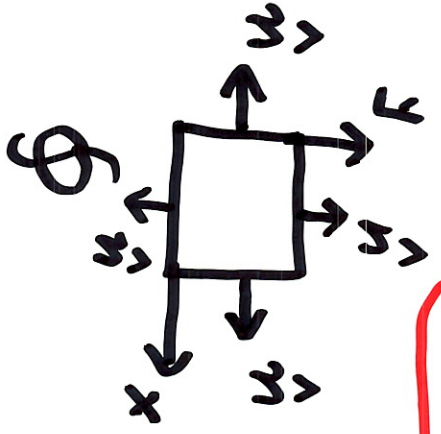
$$= \int_0^1 [u_x(1,y) - u_x(0,y)] dy + \int_0^1 [u_y(x,1) - u_y(x,0)] dx$$

$$= \int_0^1 [F_2(y) - F_1(y)] dy + \int_0^1 [F_4(x) - F_3(x)] dx$$

Solvability condition:

$$\psi = 0$$

\Rightarrow ~~Solution exists~~
infinite many sol.



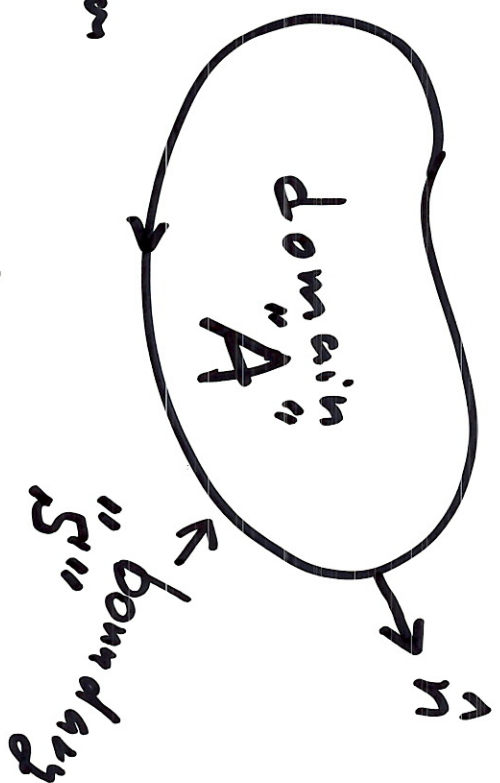
if $\psi \neq 0 \Rightarrow$ NO SOLUTION!

$$\nabla^2 u + \frac{\partial^2 u}{\partial y^2} = 0$$

$$\nabla \cdot \nabla u = 0$$

Laplace's eq.

(cf. Derivation of 2-D Heat eq.)



Neumann b.c.

impose

$$\nabla u \cdot \hat{n} = f$$

heat flux \vec{x} on the boundary
 point in/out of the domain given.

$$\iint_A (***) dA =$$

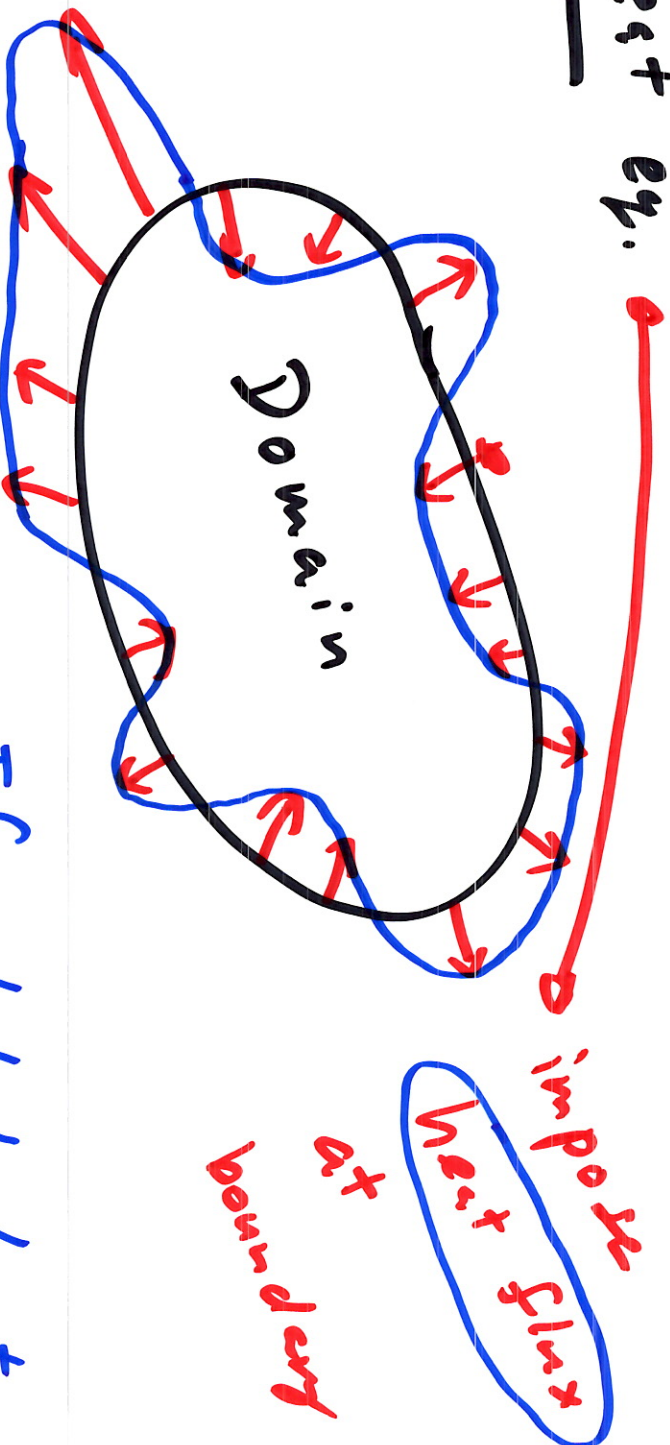
$$\iint_A \nabla \cdot \nabla u dA = 0$$

$$= \oint_S (\nabla u) \cdot \hat{n} dA$$

$$= \oint_S f dA$$

Laplace's eq. governs the steady sol. of

2-D Heat eq.



only when

$$\text{in} = \text{out}$$

⇒ Steady sol.

If total heat flux crossing

the boundary $\neq 0$

⇒ total energy $\rightarrow +\infty$ or $\rightarrow -\infty$

⇒ **NO STEADY SOL!**

⇒ Laplace's eq. has **NO SOL!**

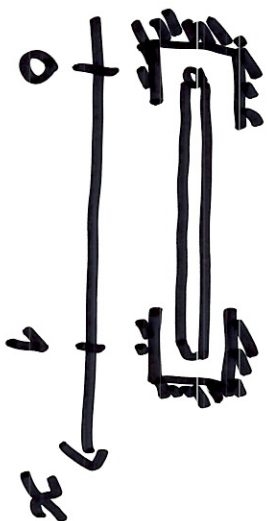
Recall (Lec 7?)

1-D Heat eq. (impose heat flux in x-dir)

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \quad \leftarrow (*)$$

$$\left. \begin{aligned} u_x(0, t) &= 0 \\ u_x(1, t) &= 0 \end{aligned} \right\}$$

$$u(x, 0) = p(x)$$



heat flux = 0

at $x=0, 1$

total energy

$$E(t) \equiv \int_0^1 u(x, t) dx$$

Calc 101

(*)

$$\frac{dE}{dt} \equiv \frac{d}{dt} \int_0^1 u dx = \int_0^1 \frac{\partial u}{\partial t} dx = \int_0^1 \frac{\partial^2 u}{\partial x^2} dx$$

$$E(t) = \text{const} \quad \text{b.c.} \rightarrow = u_x(1, t) - u_x(0, t)$$

$$= E(0) = E(\infty)$$

$$\int_0^1 \frac{\partial^2 u}{\partial x^2} dx = \dot{Q}_0$$

$$= 0$$

What if

$$\left\{ \begin{array}{l} \frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \\ u_x(0, t) = A \\ u_x(1, t) = B \\ u(x, 0) = p(x) \end{array} \right.$$

Steady sol.
 $u_s(x)$
 $\frac{d^2 u_s}{dx^2} = 0$
 $u_s'(0) = A$ $u_s'(1) = B$

\rightarrow replace eq.

$$\Rightarrow \frac{dE}{dt} = \dots = u_x(1, t) - u_x(0, t) = B - A$$

$$E(t) = E(0) + (B - A)t$$

* If $B > A \Rightarrow E(t) \rightarrow \infty$ as $t \rightarrow \infty$

$B < A \Rightarrow E(t) \rightarrow -\infty$ "

No Steady sol. !!

Solvability cond.

Only

$$B = A \Rightarrow$$

Steady sol.

$$B - A = 0$$

in 1-D

Laplace's eq + Dirichlet b.c.

→ sol. exists & is unique

Laplace's eq + Neumann b.c.

→ ~~so~~ infinite many sol.

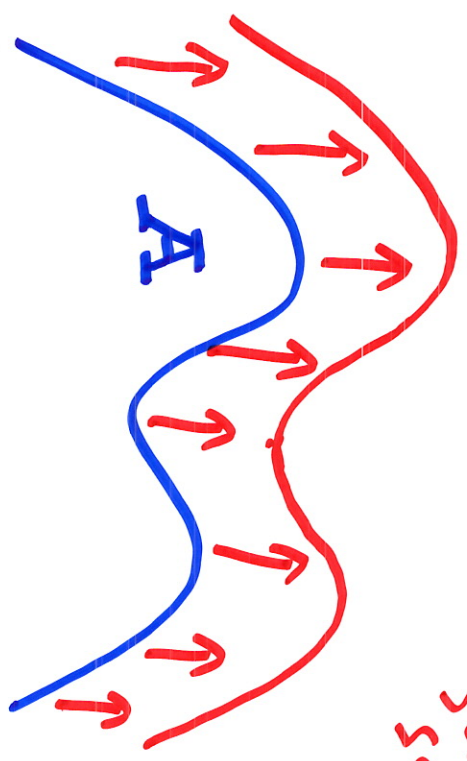
if solvability
condition ✓

NO SOLUTION

if solvability
condition ✗

(all differ
by a
const.)

$u(x,y)$



same slope
same curvature
lift up by 10

1 sol.

Mixed b.c.'s ?

(e.g. HW2-Q3)

→ find the full sol.
you will see!

HW2
Q2-b

Even in the case of
pure ~~Dir~~ Neumann or Dirichlet b.c.

you will see as well
by brute force solution

If multiple sol.

(by separation of var.
etc.)

Exist, then some expansion

coefficients will remain undetermined.