Neutron vs. Quark Stars

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Neutron stars

• Radius: $R \simeq 10 \text{ km}$ • Mass: $1.25M_{\odot} \lesssim M \lesssim 2M_{\odot}$ • Period: $1.6 \text{ ms} \lesssim P \lesssim 12 \text{ s}$ Atmosphere Superhot plasma

> Outer crust Starquakes Crystal lattice: 200 m deep nuclei + electrons

Inner crust Starquakes Crystal lattice: 1 km deep nuclei + electrons + neutron drip

🛏 20 km (12 mi) diameter

Outer core Atomic particle fluid

Inner core Solid block of subatomic particles?

- Surface magnetic field: $10^8 \text{ G} \lesssim B \lesssim 10^{14} \text{ G}$
- Core temperature:

$10 \text{ keV} \lesssim T \lesssim 10 \text{ MeV}$

Dense matter at the core



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Extremely dense matter

Nuclear matter → quark matter

Asymptotic freedom:
 α_s(μ)≪1 when μ≫Λ_{QCD}
 [Gross & Wilczek, 1973; Politzer, 1973]

•High density quark matter is weakly interacting

[Collins & Perry, 1975]

•Note: realistic densities in stars are not large enough...



 $ρ ≤ 10ρ_0$ where $ρ_0 ≃ 0.15$ fm⁻³ $\implies μ ≤ 0.5$ GeV $\implies α_s(μ) ≥ 1$

Ground state of dense matter

•Quarks are fermions (s=1/2)

•Free quarks occupy all states with $k \le k_F$

Real quarks interact





Because of the Cooper theorem, such a degenerate quark system is unstable

The ground state is a (color) superconductor

Many color superconductors

1 quark flavor (spin-1)		2 quark flavors	3 quark flavors
(e.g., only up)		(up & down)	(up, down & strange)
CSL			CFL dud du
Planar			ds sd
A/Polar			
Meissner effect: Yes		Meissner effect: No	Meissner effect: No
Superfluidity: Yes		Superfluidity: No	Superfluidity: Yes

•The actual composition of quark matter depends on its density: q_i is present if $\mu_i > m_i$ •For $\mu \leq 0.5$ GeV, c-, b- and t-quarks have no chance

Color superconductivity in stellar matter

Stellar matter is

(i) neutral (to avoid large Coulomb energy price, $E_{Coulomb} \propto n_Q^2 R^5 \gg M_{\odot}c^2$) (ii) in β -equilibrium: $\mu_d = \mu_u + \mu_e = \mu_s$

Too few d-quarks



Too many d-quarks

β -equilibrium





Unconventional Cooper pairing in stellar matter • Bottom line: Fermi momenta of all quarks are different: $p_{F,u} \neq p_{F,d} \neq p_{F,s}$ (note that $p_{F,u} \approx \mu_u, p_{F,d} \approx \mu_d \& p_{F,s} = \sqrt{\mu_s^2 - m_s^2}$)

Thus, Cooper pairing is "stressed" by the mismatch, $\delta p_F \neq 0$

What happens then?



Gapless phases (2 flavors) [I.S. & M. Huang, Phys. Lett. 564 (2003) 205; Nucl. Phys. 729 (2003) 835.] Strength of pairing (Δ_0) vs. mismatch ($\delta\mu$) $E_{\Delta^+}(p)$ $E_{\Delta^{-}}(p)$ **1. Weak coupling** $\Delta_0 \lesssim \delta \mu$ \Rightarrow normal quark matter phase `δμ+∆ δμ-Δ 2. Strong coupling d $\Delta_0 \gtrsim 2 \, \delta \mu$ \implies "usual" superconducting phase $\delta\mu + \Delta$ 3. Intermediate strength $\delta \mu \lesssim \Delta_0 \lesssim 2 \, \delta \mu$ δμ \implies gapless superconducting phase δμ - Δ μ_{ur} μ μ_{dg}

"No-go" theorem

- Stressed pairing is unavoidable [Schmitt & Rajagopal, PRD 73 (2006) 045003]
- Using graph theory, 511 pairing patterns (including all 148 inequivalent ones) were analyzed
- None of them is stressfree
- So, what does it mean?

Each line in the graph represents an allowed Cooper pairing channel **2SC phase** rs Δ_{33} (bs) Δ_{33} (bd)(bu [adapted from Schmitt & Rajagopal,

Phys. Rev. D 73 (2006) 045003]

Observational data as a tool

1. Neutron star cooling 2. Stellar "glitches" 3. Gravitational waves & r-mode instability 4. Magnetic properties 5. Transient signals from protoneutron stars





[Blaschke et al, Phys.Rev.C71 (2005) 045801]



Future direction: Transport

Conductivities

[I.S. & Ellis, PRC 66 (2002) 015802; ibid. 67 (2003) 048801]

- Heat
- Electric

Viscosities

[Manuel et al, JHEP **0509** (2005) 76] [Sa'd et al, PRD**75** (2007) 065016], [Alford & Schmitt, JPG **34** (2007) 67], [Dong et al, astro-ph/0701104], [Alford et al, nucl-th/0701067]

- Bulk
- Shear

Mean free paths

[Carter & Reddy, PRD **62** (200) 103002], [Kundu & Reddy, PRC **70** (2004) 055803], ...

– Neutrinos – Photons

Emission rates

[Jaikumar et al, PRD 66 (2002) 063003], [Reddy et al, NPA 714 (2003) 337], [Schmitt et al, PRD 73 (2006) 034012], ...



[Sad, I.S. & Rischke, PRD 75 (2007) 065016]



[from Reddy et al, NPA 714 (2003) 337]

Future direction: Thermodynamics

• Equation of state – Pressure

Energy density

[Lugones & Horvath, PRD **66** (2002) 074017], [Alford & Reddy, PRD **67** (2003) 074024], [Baldo et al, 562 (2003) 163], [Banik & Bandyopadhyay, PRD **67** (2003) 123003], ...

Specific heat

[Alford et al, PRD **71** (2005) 114011], ...

Important for cooling
Sensitive to gapless modes





Detour: Atomic systems Dense quark matter may be modeled in a tabletop experiment (using cold gas of ⁶Li or ⁴⁰K atoms)



[from the web page of Ketterle's group]

Current research directions

- Weak processes in various phases of dense quark matter
- Systematic study of transport properties of quark matter
- The study of quark matter in strong external fields
- Analysis of the observational data and search for signatures of new states of matter
- Development of non-perturbative techniques for studying quark matter
- High temperature quark matter (RHIC & LHC)
- Cross-disciplinary insight into quark dynamics (e.g., from physics of cold atoms, graphene, high-T_c superconductivity, etc.)

Summary

- Deconfined quark matter is likely to exist in stars
- β-equilibrium plays and important role in shaping the ground state
- Such matter is an unconventional color superconductor
- Phase structure of dense matter is very rich
- Observational data may help to shed light on the phase diagram

