



ASU college of
technology & innovation

ARIZONA STATE UNIVERSITY

Chiral asymmetry and axial anomaly in magnetized relativistic matter

Igor Shovkovy

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- Relativistic matter & strong magnetic fields ($10^8 - 10^{15}$ G) are common inside compact stars
 - **Electrons** in white dwarfs
 $T \ll m \lesssim \mu$ (i.e., $T \lesssim 1$ keV & $\mu \lesssim 1$ MeV)
 - **Protons** in dense neutron matter
 $T \ll m \lesssim \mu$ (i.e., $T \lesssim 10$ MeV & $\mu \simeq 1$ GeV)
 - **Electrons** in dense neutron matter
 $m \lesssim T \ll \mu$ (i.e., $T \lesssim 10$ MeV & $\mu \lesssim 100$ MeV)
 - **Quark matter** in stellar cores (if formed)
 $T \lesssim m \ll \mu$ (i.e., $T \lesssim 10$ MeV & $\mu \lesssim 500$ MeV)

- Axial vector current in relativistic matter in a magnetic field (3+1 dimensions)

$$\langle j_5^3 \rangle_0 = \frac{-eB}{2\pi^2} \mu_0 \quad (\text{free theory!})$$

[Metlitski & Zhitnitsky, Phys Rev D **72**, 045011 (2005)]

- Is there a dynamical parameter Δ (“chiral shift”) associated with this condensate?

$$\mathcal{L} = \mathcal{L}_0 + \mathcal{L}_\Delta \quad \text{where} \quad \mathcal{L}_\Delta \simeq \Delta \bar{\psi} \gamma^3 \gamma^5 \psi$$

- Note: $\Delta=0$ is not protected by any symmetry

Axial anomaly

- Does the chiral shift modify the axial anomaly relation?
- Using point splitting method, one derives

$$\begin{aligned} \langle \partial_\mu j_5^\mu(u) \rangle &= -\frac{e^2 \epsilon^{\beta\mu\lambda\sigma} F_{\alpha\mu} F_{\lambda\sigma} \epsilon^\alpha \epsilon_\beta}{8\pi^2 \epsilon^2} \left(e^{-is_\perp \Delta \epsilon^3} + e^{is_\perp \Delta \epsilon^3} \right) \\ &\rightarrow -\frac{e^2}{16\pi^2} \epsilon^{\beta\mu\lambda\sigma} F_{\beta\mu} F_{\lambda\sigma} \quad \text{for } \epsilon \rightarrow 0 \end{aligned}$$

[E. V. Gorbar & V. A. Miransky, I.A. Shovkovy, Phys. Lett. B 695 (2011) 354]

- Therefore, the chiral shift does **not** affect the conventional axial anomaly relation

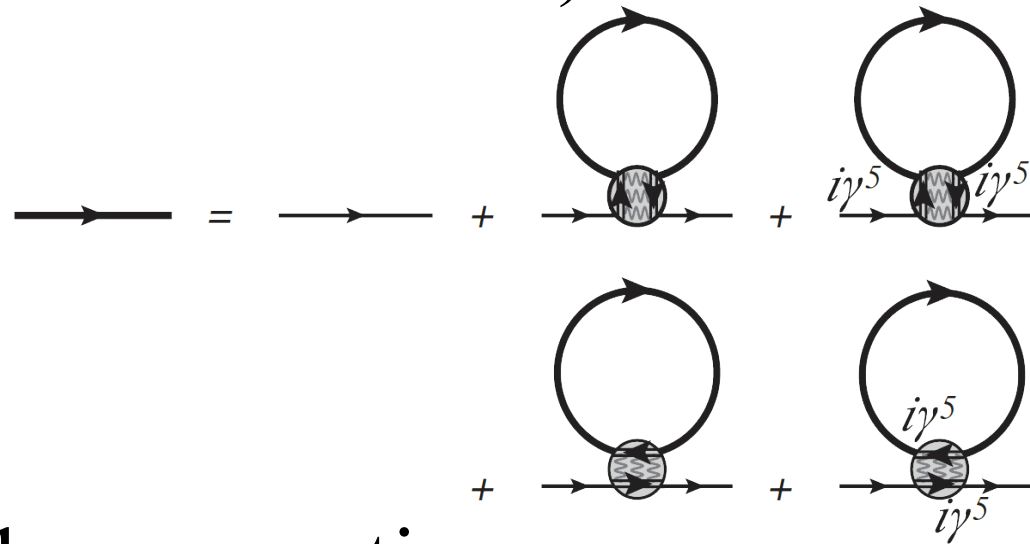
- Does the chiral shift give any contribution to the axial current?
- In point splitting method, one derives

$$\langle j_5^\mu \rangle_{\text{singular}} = -\frac{\Delta}{2\pi^2 \epsilon^2} \delta_3^\mu \sim \frac{\Lambda^2 \Delta}{2\pi^2} \delta_3^\mu$$

[E. V. Gorbar & V. A. Miransky, I.A. Shovkovy, Phys. Lett. B 695 (2011) 354]

- This is consistent with the NJL calculations
- Since $\Delta \sim g\mu eB/\Lambda^2$, the axial current is finite

- NJL model (local interaction)



- This leads to three equations:

$$\mu = \mu_0 - \frac{1}{2} G_{\text{int}} \langle j^0 \rangle \quad (\text{“effective” chemical potential})$$

$$m = m_0 - G_{\text{int}} \langle \bar{\psi} \psi \rangle \quad (\text{dynamical mass})$$

$$\Delta = -\frac{1}{2} G_{\text{int}} \langle j_5^3 \rangle \quad (\text{chiral shift parameter})$$

- Magnetic catalysis solution (vacuum state):

$$m^2 \simeq \frac{|eB|}{\pi} \exp\left(-\frac{4\pi^2}{G_{\text{int}}|eB|}\right) \quad \left(|\mu_0| \lesssim \frac{m}{\sqrt{2}}\right)$$

$$\Delta = 0 \quad \& \quad \mu = \mu_0$$

- State with a chiral shift (nonzero density):

$$m = 0 \quad \& \quad \mu \simeq \frac{\mu_0}{1 + g/(\Lambda l)^2}$$

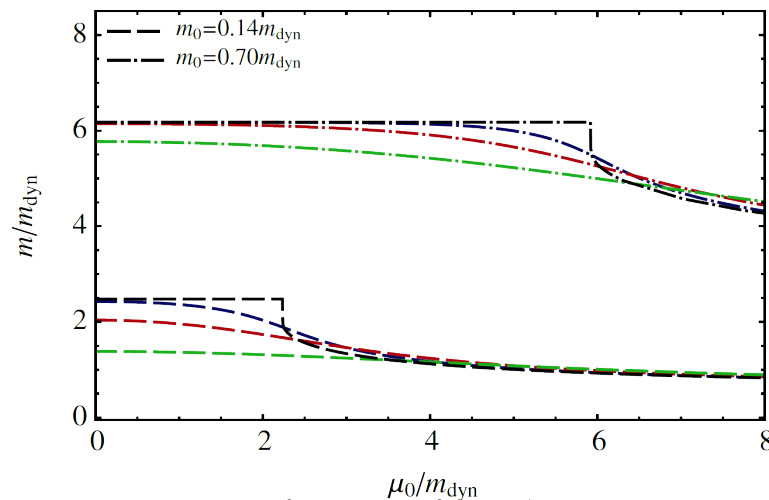
$$\Delta = \frac{gs_{\perp}\mu}{(\Lambda l)^2 + \frac{1}{2}g(\Lambda l)^2} \quad \left(|\mu_0| \gtrsim \frac{m}{\sqrt{2}}\right)$$

Solutions

- Magnetic catalysis solution (vacuum state):

$$m^2 \simeq \frac{|eB|}{\pi} \exp\left(-\frac{4\pi^2}{G_{\text{int}}|eB|}\right)$$

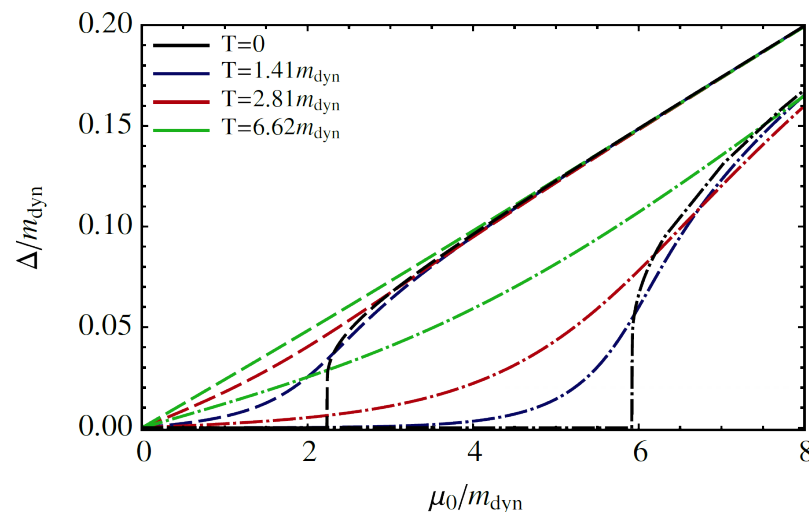
$$\Delta = 0 \quad \& \quad \mu = \mu_0$$



- State with a chiral shift (nonzero density):

$$m = 0 \quad \& \quad \mu \simeq \frac{\mu_0}{1 + g/(\Lambda l)^2}$$

$$\Delta = \frac{gs_{\perp}\mu}{(\Lambda l)^2 + \frac{1}{2}g(\Lambda l)^2}$$



Chiral shift and Fermi surface

- Chirality is approx. well defined at Fermi surface ($|k^3| \gg m$)
- L-handed Fermi surface:

$$n = 0 : k^3 = +\sqrt{(\mu - s_{\perp}\Delta)^2 - m^2}$$

$$n > 0 : k^3 = +\sqrt{\left(\sqrt{\mu^2 - 2n|eB|} - s_{\perp}\Delta\right)^2 - m^2}$$

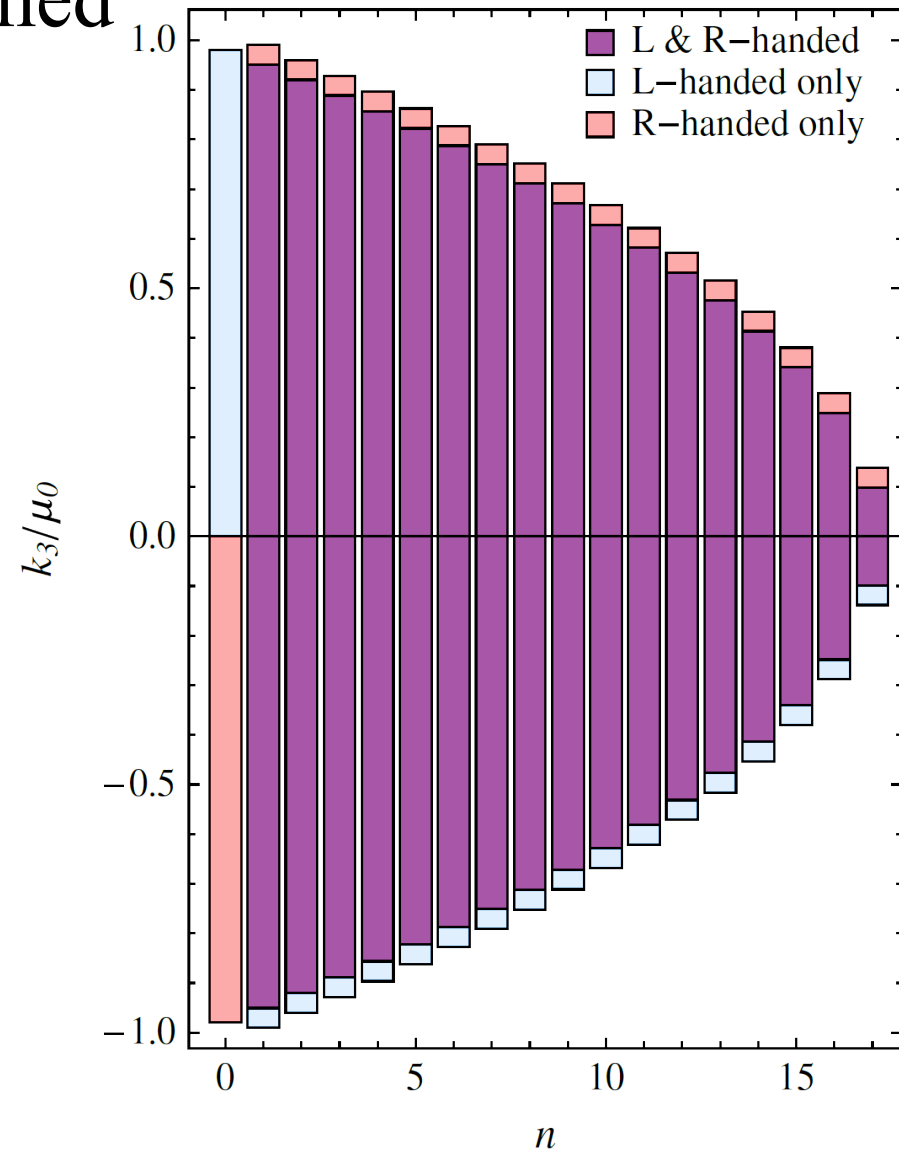
$$k^3 = -\sqrt{\left(\sqrt{\mu^2 - 2n|eB|} + s_{\perp}\Delta\right)^2 - m^2}$$

- R-handed Fermi surface:


$$n = 0 : k^3 = -\sqrt{(\mu - s_{\perp}\Delta)^2 - m^2}$$

$$n > 0 : k^3 = -\sqrt{\left(\sqrt{\mu^2 - 2n|eB|} - s_{\perp}\Delta\right)^2 - m^2}$$

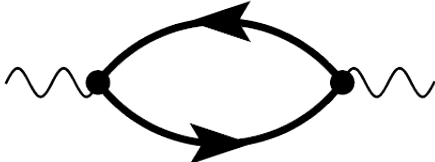
$$k^3 = +\sqrt{\left(\sqrt{\mu^2 - 2n|eB|} + s_{\perp}\Delta\right)^2 - m^2}$$



- Chiral shift is also generated in QED

$$(\text{---})^{-1} = (\text{---})^{-1} + \text{---} \text{---} \text{---}$$


The diagram shows a fermion propagator (a solid line with an arrow) equal to the free propagator plus a correction term. The correction term is a fermion loop (a solid line with an arrow forming a circle) attached to the propagator at two vertices.

$$(\text{~})^{-1} = (\text{~})^{-1} + \text{~} \text{---} \text{~}$$


The diagram shows a photon propagator (a wavy line) equal to the free propagator plus a correction term. The correction term is a fermion loop (a solid line with an arrow forming a circle) attached to the propagator at two vertices.

- Screening effects are important
- Large density estimate for the chiral shift:

$$\Delta^{QED} \simeq -\frac{\alpha eB}{\pi M_D^2} \mu, \quad \text{where} \quad M_D = \frac{\alpha \mu}{2\pi}$$

- Numerically, $\Delta \simeq 1 - 100 \text{ MeV}$ ($B \simeq 10^{13} - 10^{16} \text{ G}$)

- New dynamical parameter (chiral shift) is generated in magnetized dense matter
- Chiral shift induces a chiral asymmetry at the Fermi surface
- Potential applications:
 - Pulsar kicks (?)
 - Quark stars
 - Facilitation of supernova explosions (?)
 - Axial current in QGP
 - modified CME (chiral magnetic effect)