

Fast chemical equilibration via Hagedorn states in heavy ion collisions

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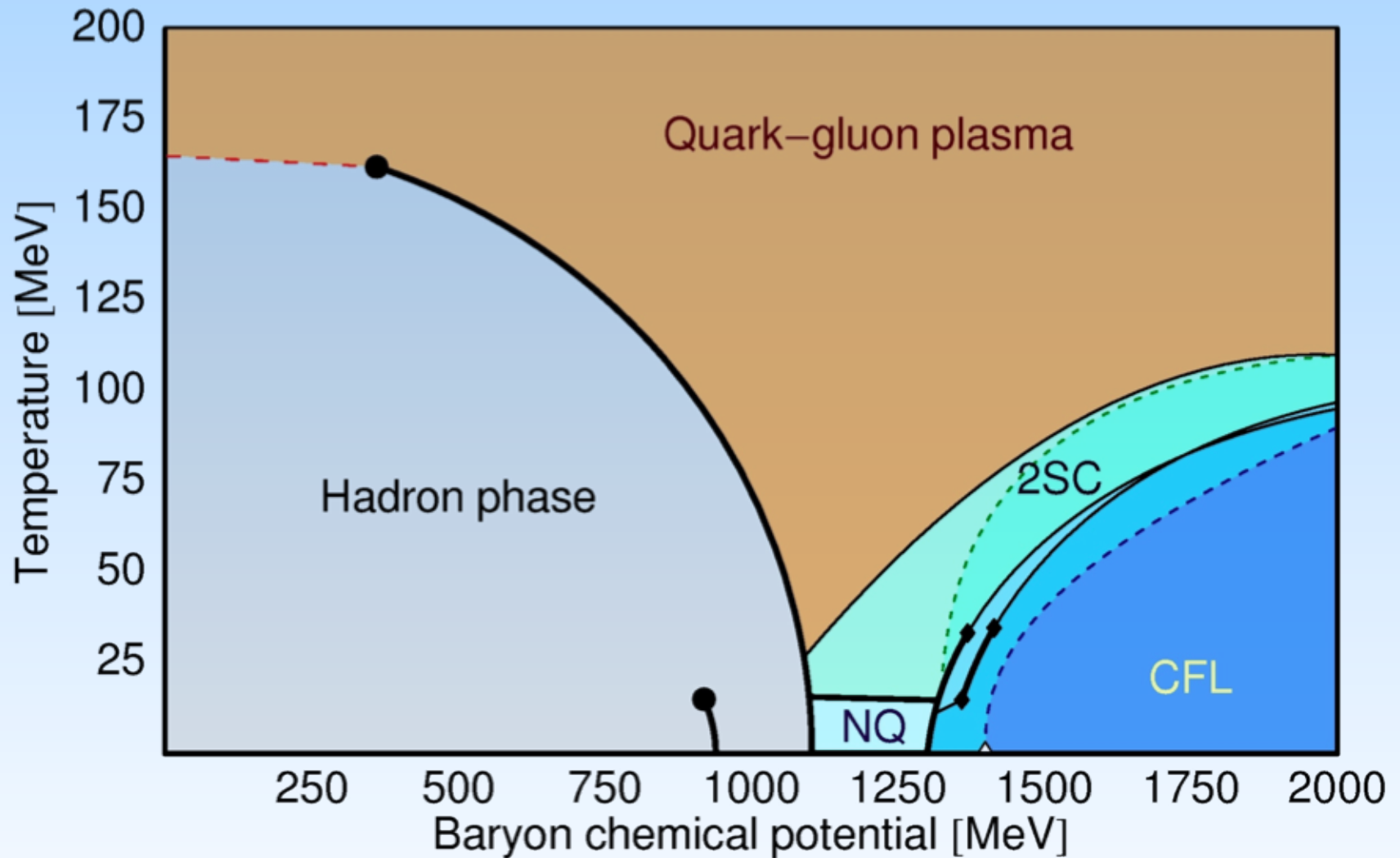
Arizona State University



Symposium on Contemporary Subatomic Physics (JoeFest)
Montreal, Canada, June 12-14, 2012

Phase diagram (1)

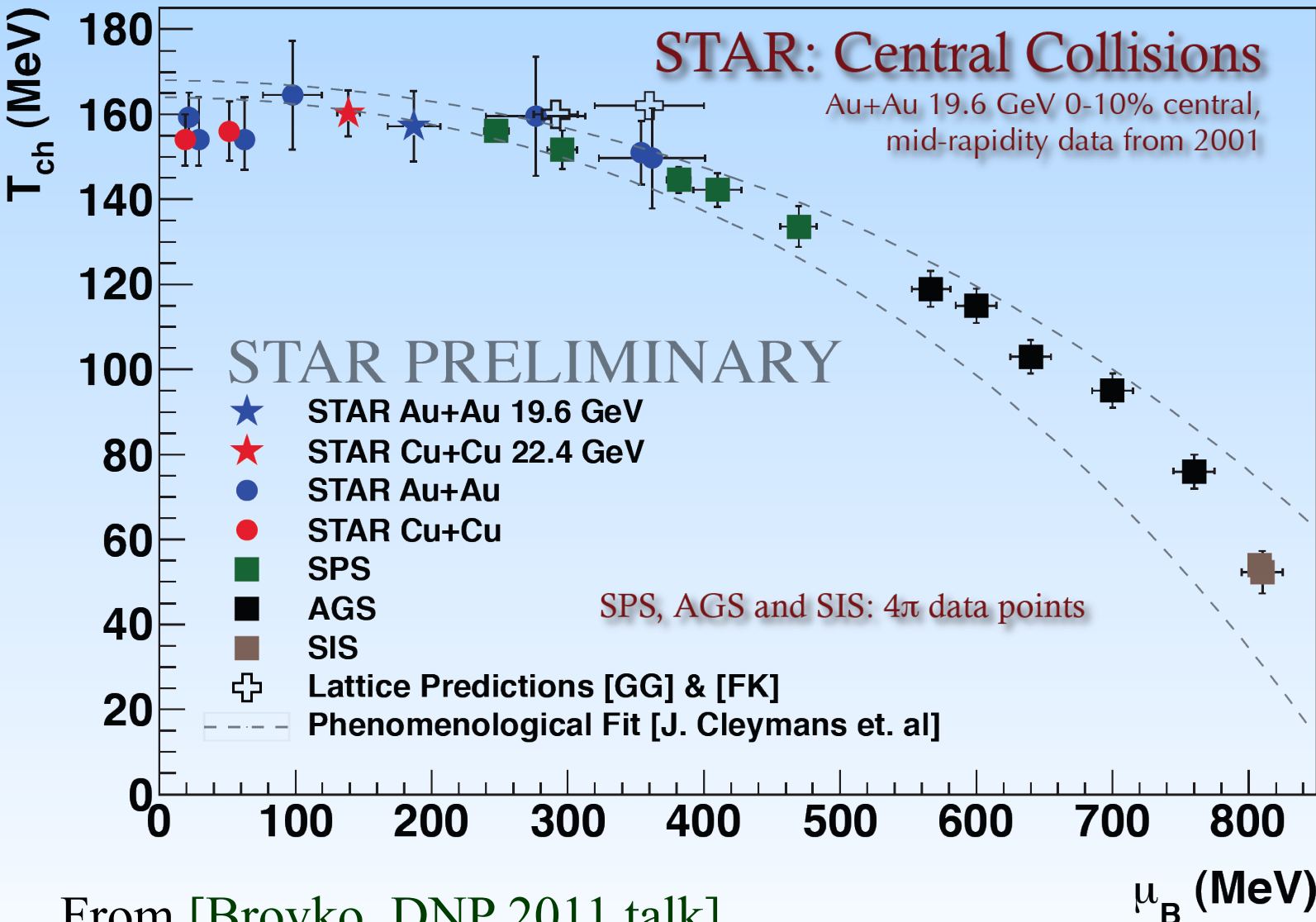
as theoreticians imagine it...



Adapted from [Rüster et al., Phys. Rev. D 72, 034004 (2005)]

Phase diagram (2)

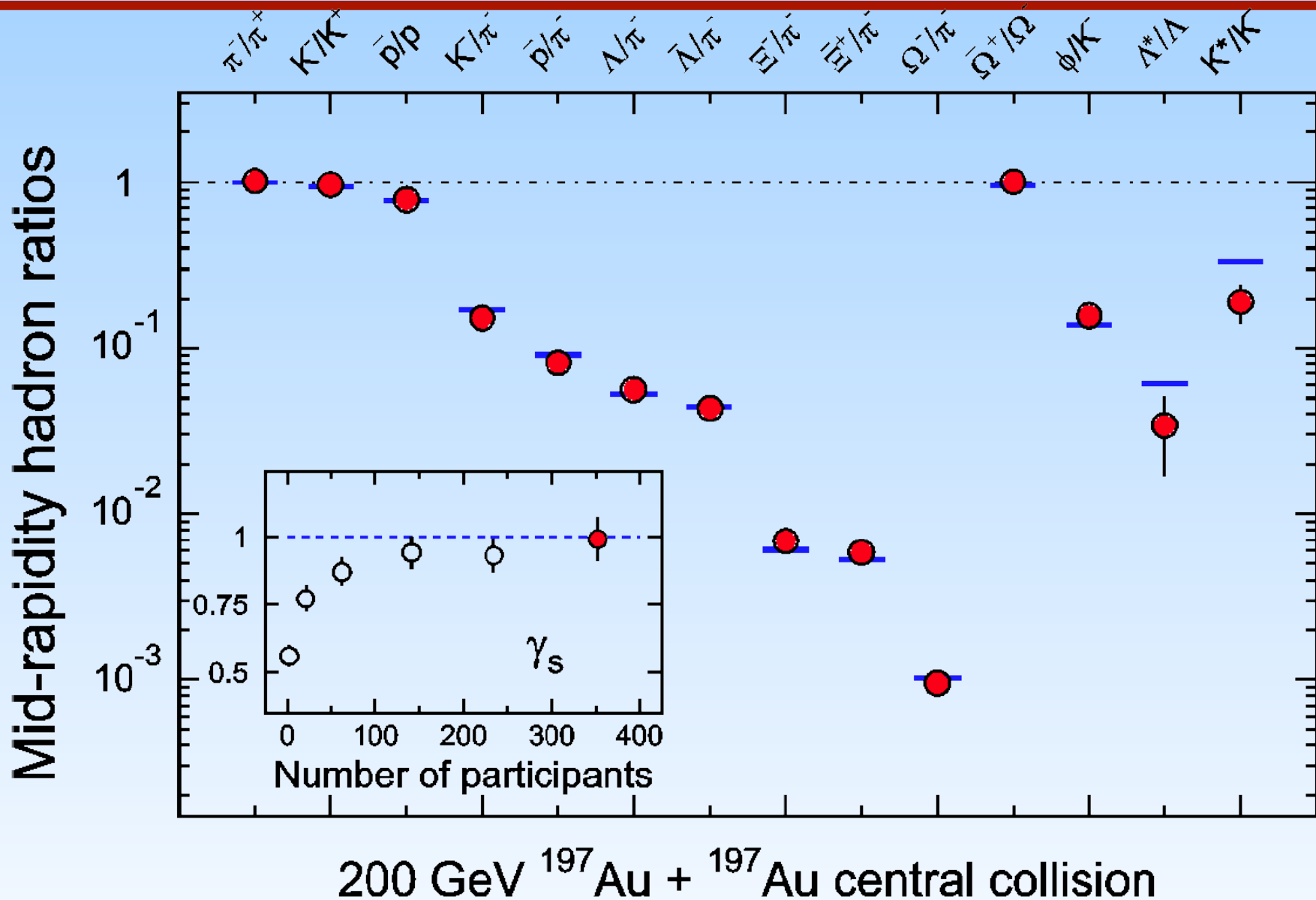
as RHIC physicists get to know it...



p+p 200, Au+Au 200, 130, 62.4 GeV
 [STAR: PhysRevC.79.034909]
 Cu+Cu 62.4, 200 GeV
 [STAR: PhysRevC.83.034910]
 Au+Au 9.2 GeV
 [STAR: PhysRevC.81.024911]
 7.7, 11, 39 GeV: Kumar QM2011
 19.6 Au+Au & 22.4 GeV Cu+Cu:
 Mall SQM2011
 E866/917 PLB476.1.2000
 E895 PRC68.054905.2003
 E802 PRC58.3523.1998
 NA44 PRC66.044907.2002
 NA49 PRC66.054902.2002
 Braun-Munzinger, Heppe, Stachel
 Phys.Lett.B465.15-20. 1999
 Kaneta, Xu, QM04 nucl-th/0405068

From [Brovko, DNP 2011 talk]

Simple observation



From [STAR Collaboration, Nucl. Phys. A 757, 102 (2005)]

Big puzzle

- What is the mechanism of chemical equilibration in heavy ion collisions?
- Rates of hadronic processes are too slow...
- The problem is further amplified for strangeness
- Some suggestions
 - “born into equilibrium” (?)
 - unknown “collective effects” (?)
 - multi-particle interactions (?)
 - or a new incarnation of an old idea...

Hadron gas?

- Two-body reactions (e.g., $\pi\pi \leftrightarrow B\bar{B}$) are not sufficient

[Koch, Muller, Rafelski, Phys. Rev. **142**, 167 (1986)]

- Multi-particle reactions (e.g., $\pi\pi\pi K\bar{K} \leftrightarrow B\bar{B}$) become important

[Rapp and Shuryak, Phys. Rev. Lett. 86, 2980 (2001)]

[Greiner and Leupold, J. Phys. G 27, L95 (2001)]

yielding equilibration times $\sim 1-3$ fm/c at SPS

$$\sigma \approx 50 \text{ mb}, \quad n_B^{eq} \approx \rho_0 \quad \Rightarrow \quad \tau \approx 1 - 3 \text{ fm/c}$$

How about RHIC?

- Two-particle annihilation rate

$$\Gamma \approx \langle \sigma v \rangle n_{\bar{B}}$$

Naively, one has

$$\sigma \approx 30 \text{ mb}, \quad n_B^{eq} = n_{\bar{B}}^{eq} \approx 0.04 \text{ fm}^{-3} \Rightarrow \tau_{\Omega} \geq 10 \text{ fm/c}$$

which is about 10 fm/c at RHIC

How about multi-particle processes?

Equilibration times with Joe

- Baryon production rates, e.g.,
[Kapusta & Shovkovy, PRC **68**, 014901 (2003)]

$$r(n\bar{p}) = 2r_+(m_N, m_N) + 2r_-(m_N, m_N)$$

$$r(p\bar{p}) = 2r_+(m_N, m_N) + \frac{82}{81}r_-(m_N, m_N)$$

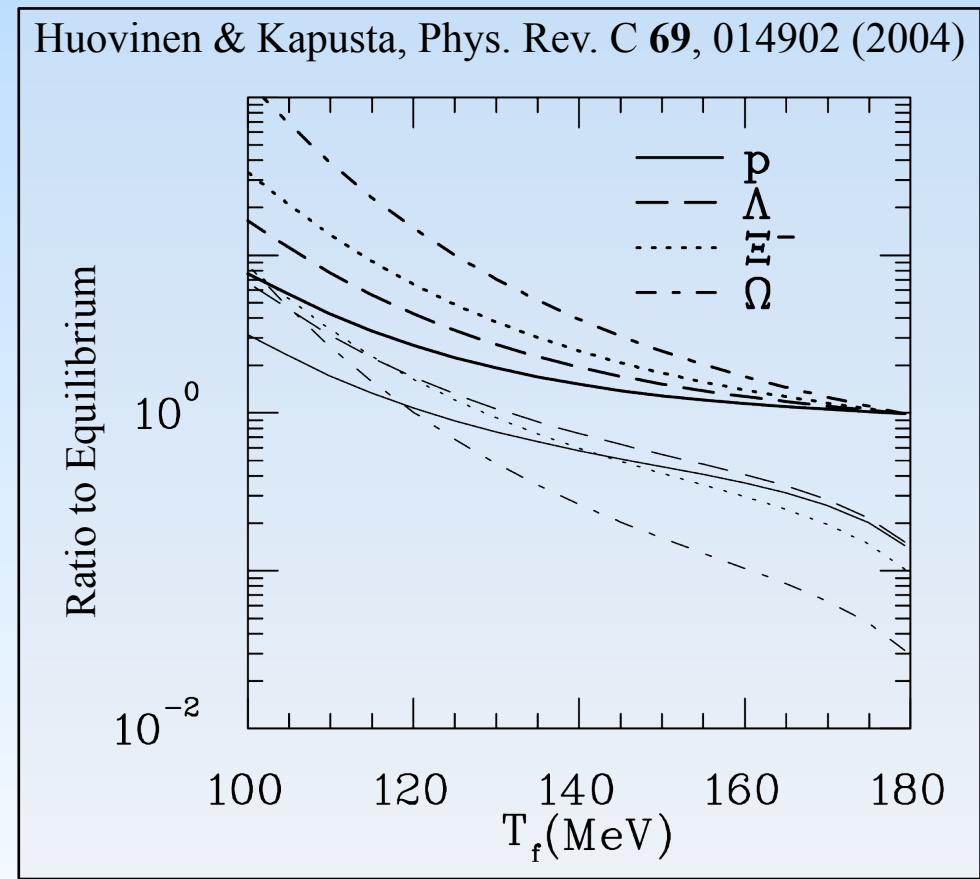
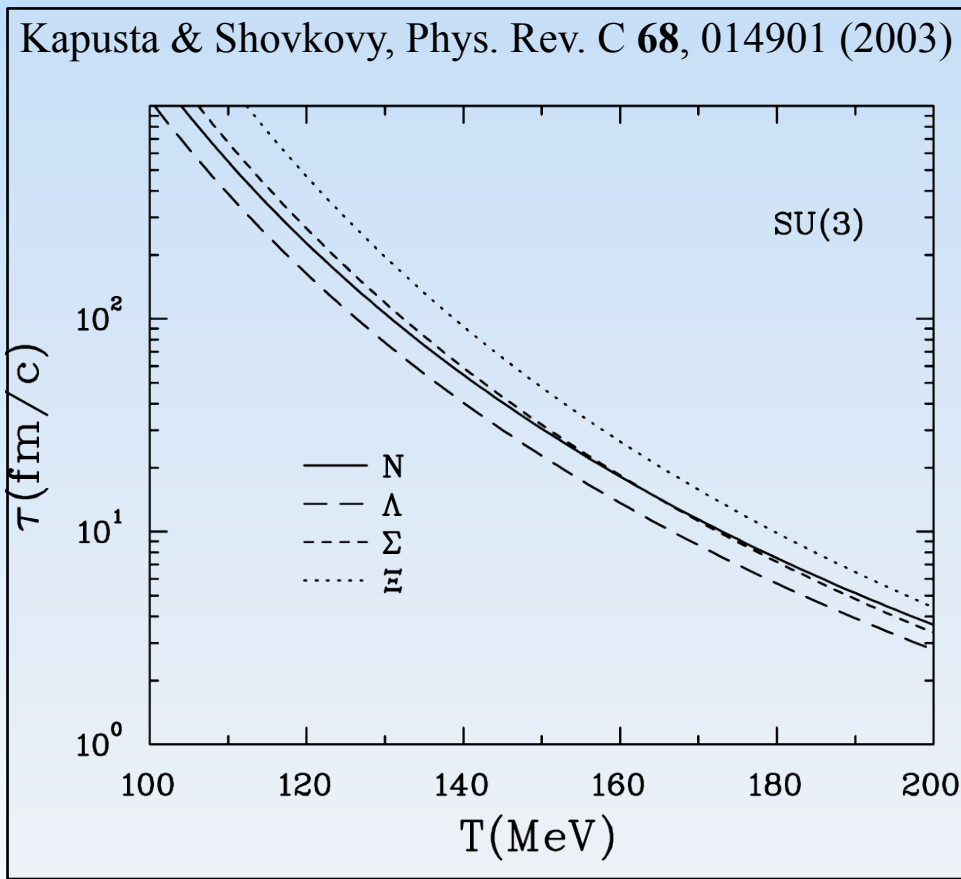
$$r(\Lambda\bar{p}) = 3r_+(m_\Lambda, m_N) + \frac{25}{27}r_-(m_\Lambda, m_N)$$

$$r(\Xi\bar{\Lambda}) = 3r_+(m_\Lambda, m_\Xi) + \frac{1}{27}r_-(m_\Lambda, m_\Xi)$$

$$r_\pm(m_1, m_2) = \frac{2}{(4\pi)^7} \frac{F_{\text{ANN}}^2(s)}{e^{\beta(E_1+E_2)} - 1} \frac{2s^2 - (m_1^2 + m_2^2)s - (m_1^2 - m_2^2)^2 \pm 6m_1m_2s}{f_\pi^4} \left(1 + \frac{\alpha_s}{\pi} \right)$$

Results

- Using fluctuation-dissipation theorem, obtain rates of baryon production



Multi-particle collisions

- Intuition based argument:

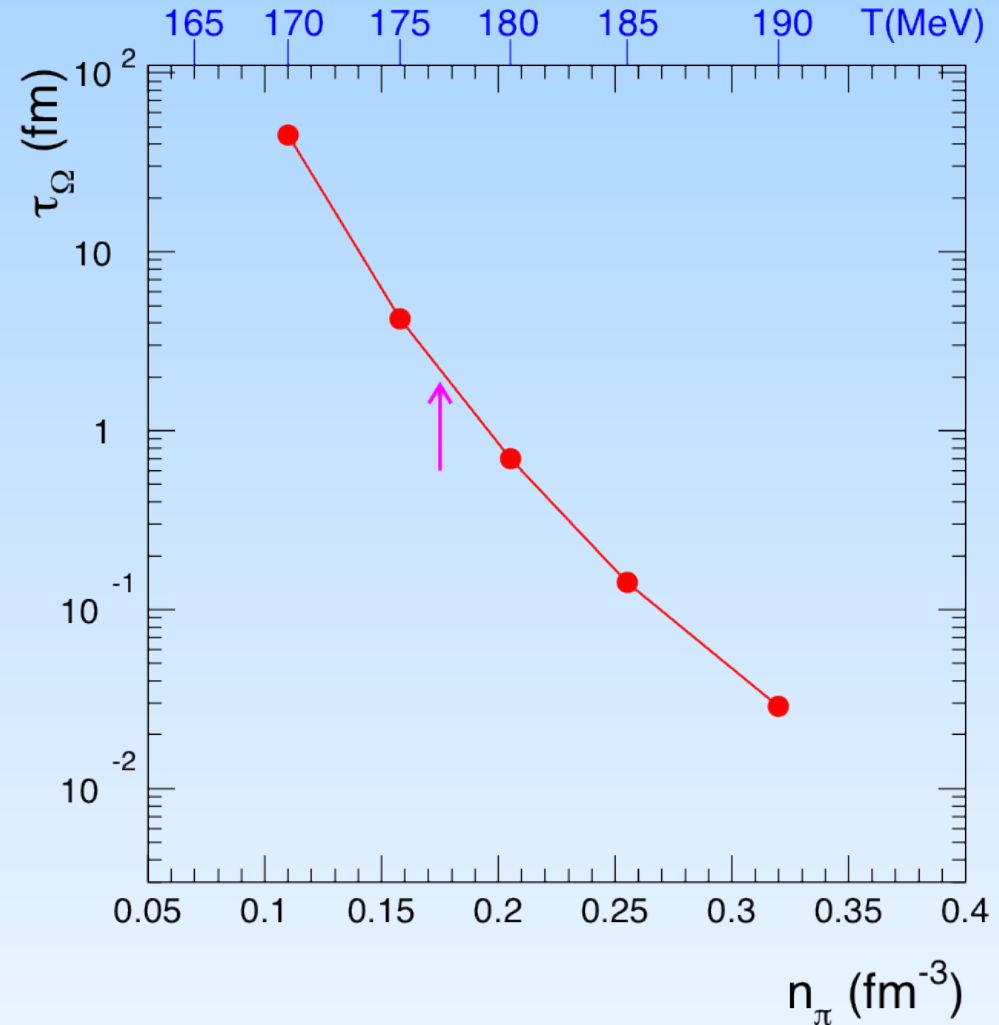
[Braun-Munzinger, Stachel, Wetterich, Phys. Lett. **B 596**, 61 (2004)]

- Large energy density near T_c
- Overpopulation of pions/kaons
- Multi-particle rates overwhelm
- Equilibration is reached



..., but there is a one **LITTLE** problem:

one is left with an overpopulation of (anti-)baryons



What if ...

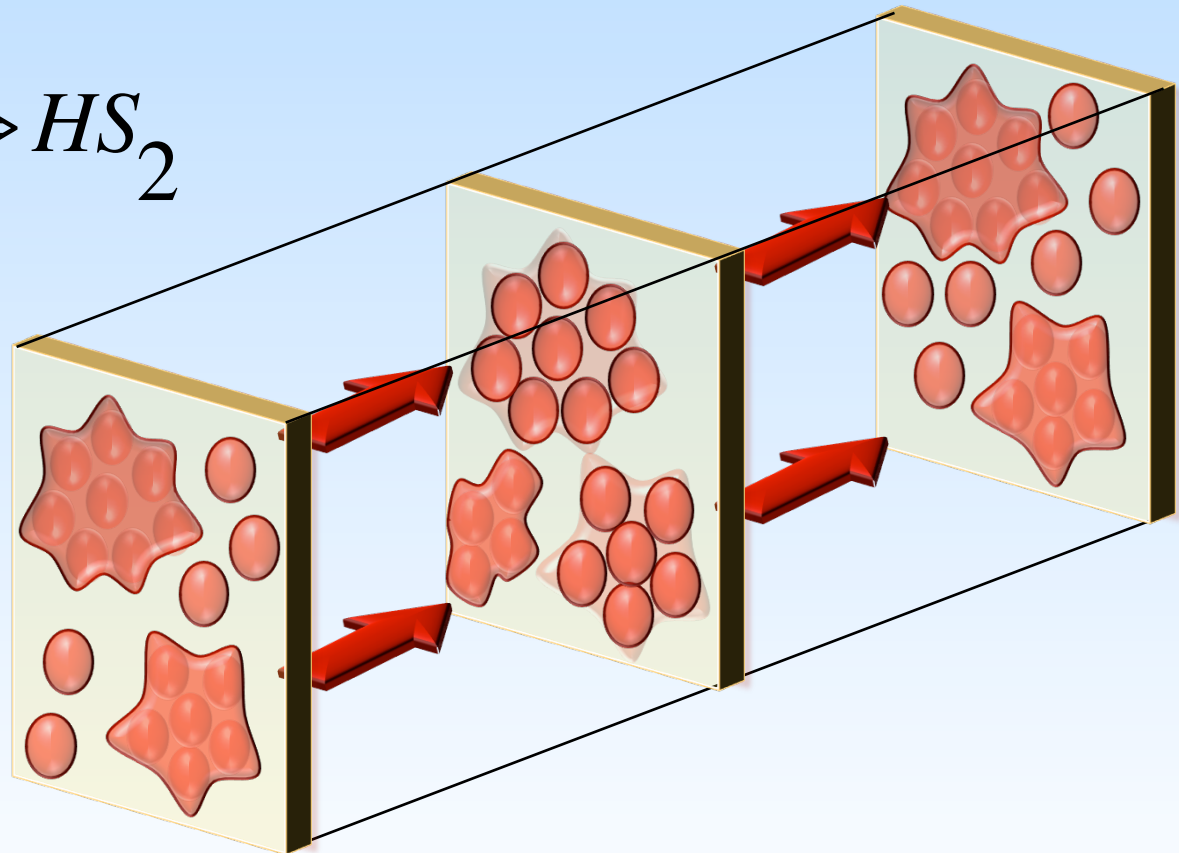
there were Hagedorn states in the spectrum?

[Greiner et al., J. Phys. **G 31**, S725 (2005)]

[Noronha-Hostler et al., Phys. Rev. Lett. **100**, 252301 (2008)]

[Noronha-Hostler et al., Phys. Rev. C **81**, 054909 (2010)]

$$HS_1 \Rightarrow n\pi + B\bar{B} \Rightarrow HS_2$$



Rough estimate...

$$M_{HS} = 3 - 6 \text{ GeV}$$

$$n_{HS} = 0.05 - 0.15 \text{ fm}^{-3}$$

$$\Gamma_{HS}^{tot} \approx 0.5 - 1 \text{ GeV}$$

$$\Gamma_{BB\bar{X}} \approx 100 - 300 \text{ MeV}$$

$$\Gamma_{BB\bar{}}^{prod} \approx n_{HS} \Gamma_{BB\bar{X}} \approx 0.05 \text{ fm}^{-4}$$

$$\tau_{BB\bar{}} \approx 1 \text{ fm}/c$$



Detailed model

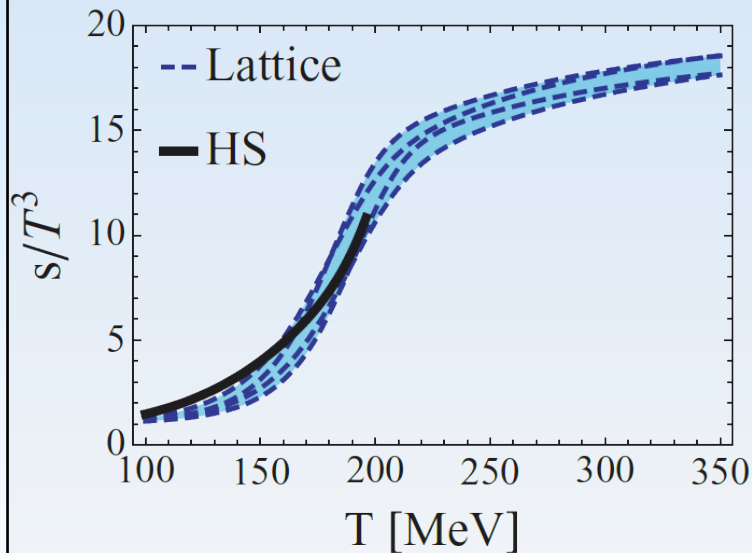
- Hagedorn mass spectrum:

$$\rho = \int_{M_0}^M F(m) e^{\frac{m}{T_H}} dm, \quad \text{where} \quad F(m) = \frac{A}{(m^2 + m_r^2)^{5/4}}$$

Cf. lattice [RBC-Bielefeld]

$$T_H = 196 \text{ MeV}$$

$$A = 0.5 \text{ GeV}^{3/2}$$



$$m_r = 0.5 \text{ GeV}$$

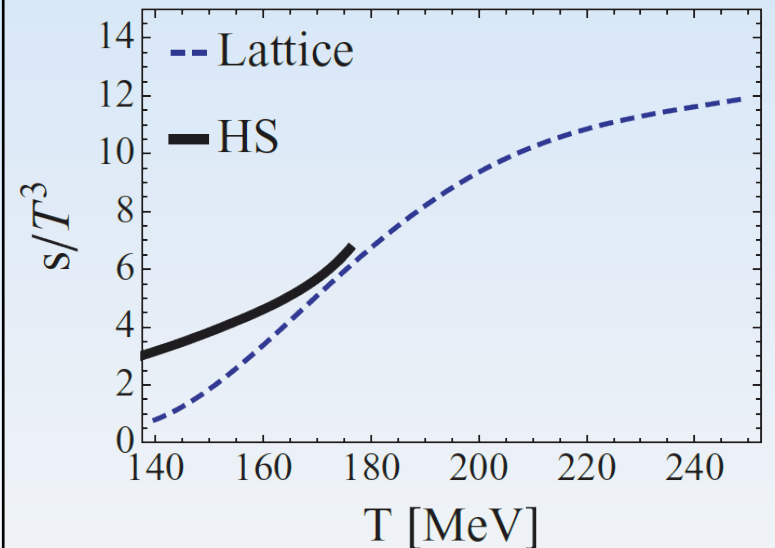
$$M_0 = 2 \text{ GeV}$$

$$M = 12 \text{ GeV}$$

Cf. lattice [Aoki et al.]

$$T_H = 176 \text{ MeV}$$

$$A = 0.1 \text{ GeV}^{3/2}$$



Master equations

- Processes $n\pi \leftrightarrow HS \leftrightarrow n'\pi + X \bar{X}$

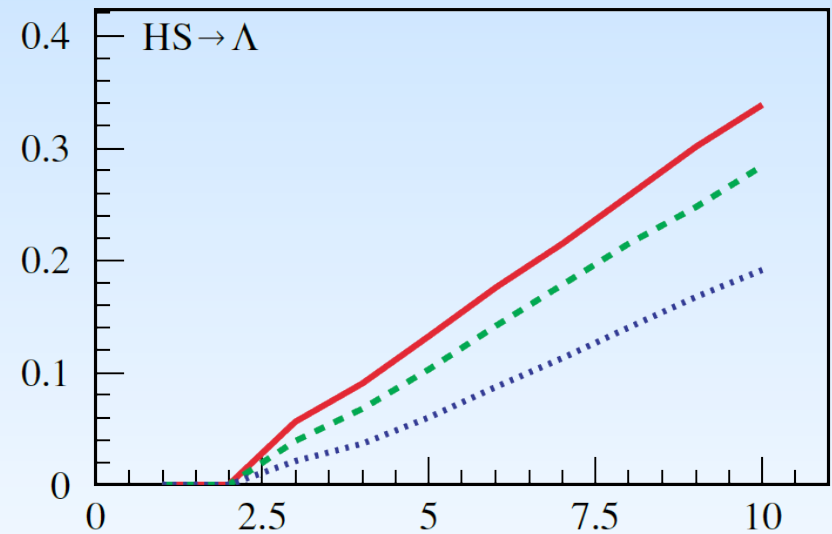
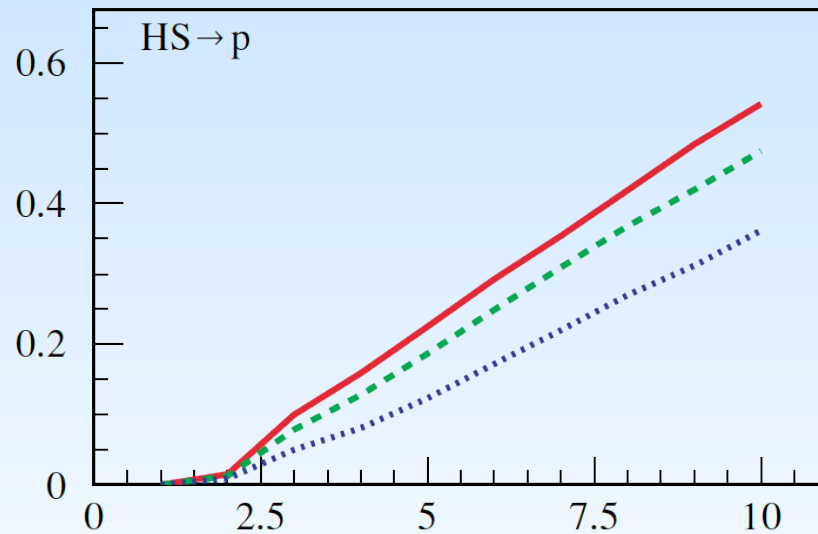
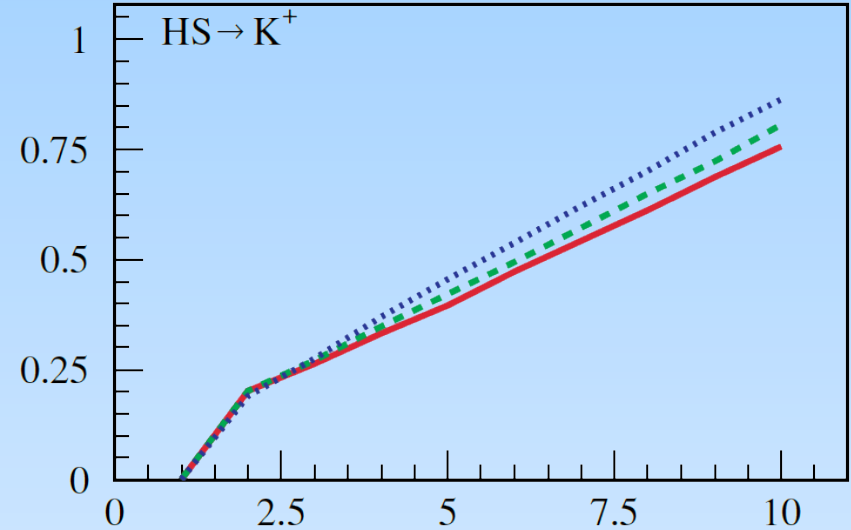
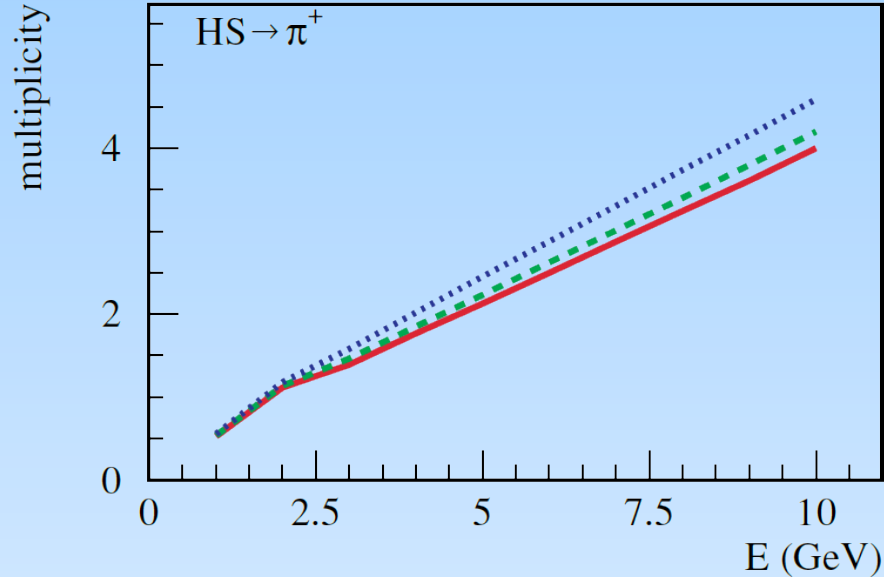
- Schematically $\frac{dN}{dt} = -\text{"loss"} + \text{"gain"}$

$$\frac{dN_{R(i)}}{dt} = -\Gamma_i N_{R(i)} + \sum_n \Gamma_{i,\pi} N_{R(i)}^{eq} \left(\frac{N_\pi}{N_\pi^{eq}} \right)^n B_{i \rightarrow n\pi} + \Gamma_{i,B\bar{B}} \frac{N_{R(i)}^{eq}}{(N_{B\bar{B}}^{eq})^2} \left(\frac{N_\pi}{N_\pi^{eq}} \right)^{\langle n \rangle} (N_{B\bar{B}})^2$$

$$\frac{dN_\pi}{dt} = \sum_i \sum_n \Gamma_{i,\pi} n B_{i \rightarrow n\pi} \left(N_{R(i)} - N_{R(i)}^{eq} \left(\frac{N_\pi}{N_\pi^{eq}} \right)^n \right) + \sum_i \Gamma_{i,B\bar{B}} \langle n \rangle \left(N_{R(i)} - \frac{N_{R(i)}^{eq}}{(N_{B\bar{B}}^{eq})^2} \left(\frac{N_\pi}{N_\pi^{eq}} \right)^{\langle n \rangle} (N_{B\bar{B}})^2 \right)$$

$$\frac{dN_{B\bar{B}}}{dt} = \sum_i \Gamma_{i,B\bar{B}} \langle n \rangle \left(\frac{N_{R(i)}^{eq}}{(N_{B\bar{B}}^{eq})^2} \left(\frac{N_\pi}{N_\pi^{eq}} \right)^{\langle n \rangle} (N_{B\bar{B}})^2 - N_{R(i)} \right)$$

Microcanonical decays of HS



Calculation of Fuming Liu [Greiner et al., J. Phys. **G 31**, S725 (2005)]

Model parameters (1)

- Mass range of Hagedorn states: 2 - 12 GeV
- Branching ratios

– Gaussian distribution: $B_{i \rightarrow n\pi} \approx \frac{1}{\sigma\sqrt{2\pi}} e^{-\frac{(n-\langle n \rangle)^2}{(2\sigma)^2}}$

$$\sigma_i = 0.5 \frac{m_i}{m_p} \approx 0.9 - 22,$$

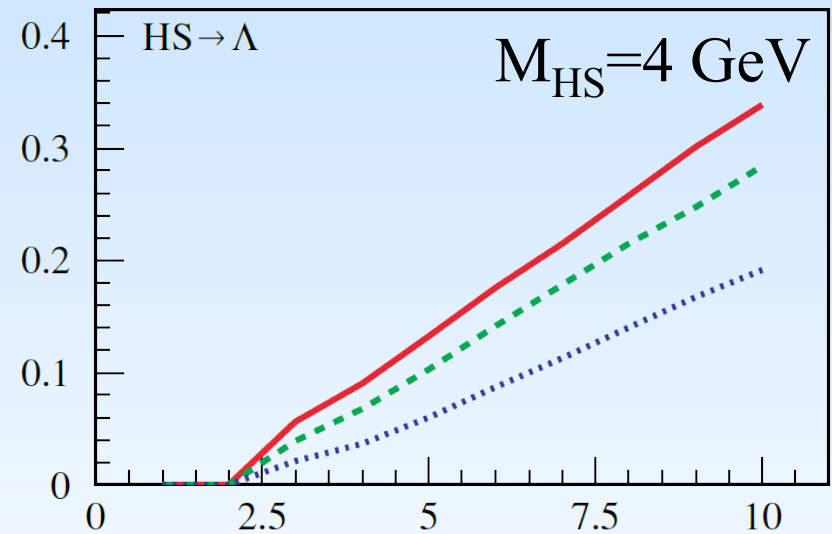
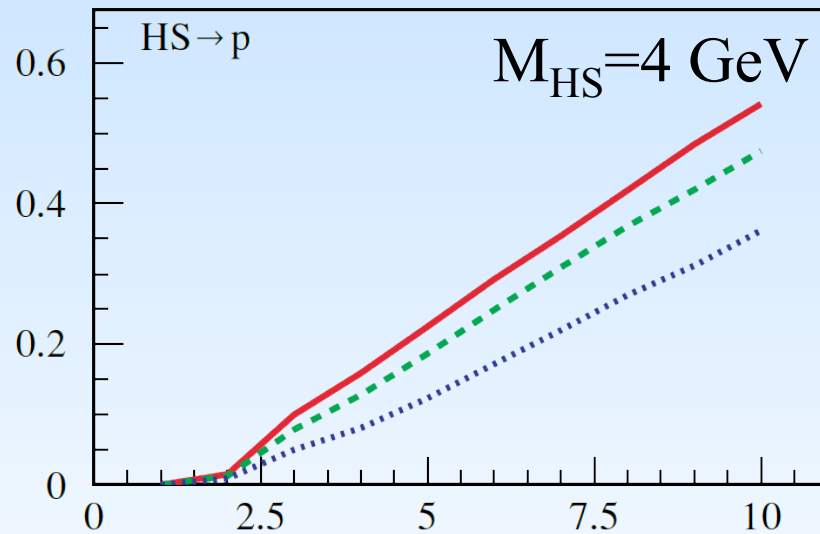
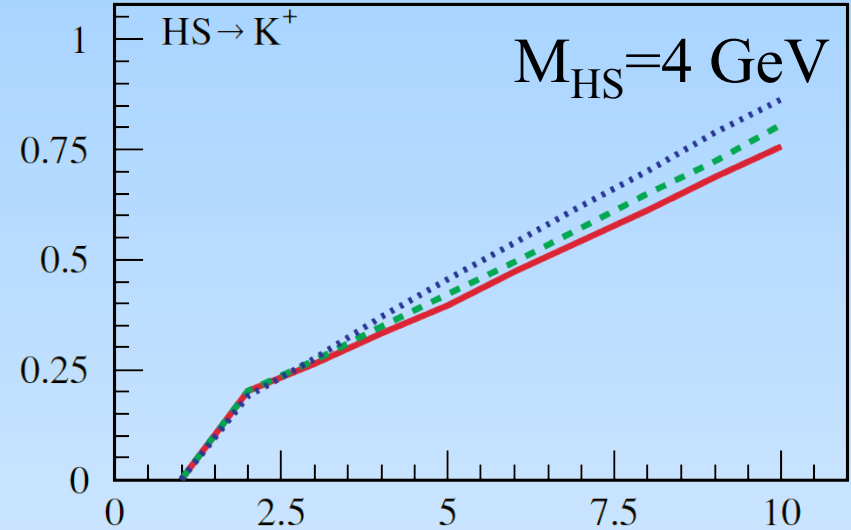
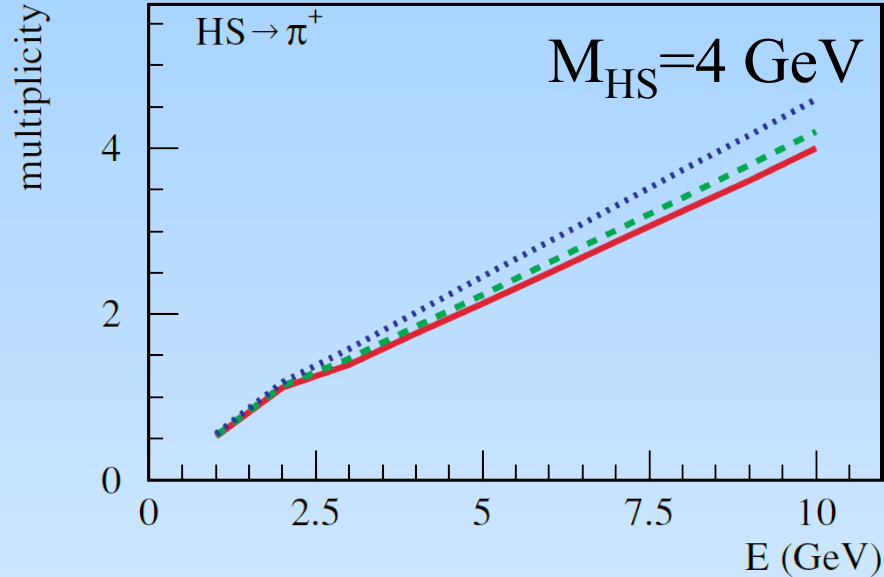
$$HS \leftrightarrow n\pi: \quad \langle n_i \rangle = 0.9 + 1.2 \frac{m_i}{m_p} \approx 3 - 34$$

$$HS \leftrightarrow n\pi + X\bar{X}: \quad \langle n_{i,x} \rangle = \frac{2.7}{1.9} \left(0.9 + 0.37 \frac{m_i}{m_p} \right) \approx 2 - 7$$

- Decay widths

$$\Gamma_i = \Gamma_{i,\pi} + \Gamma_{i,X\bar{X}} = 0.15 m_i - 27 \text{ MeV} \approx 250 - 1800 \text{ MeV}$$

Microcanonical decays of HS



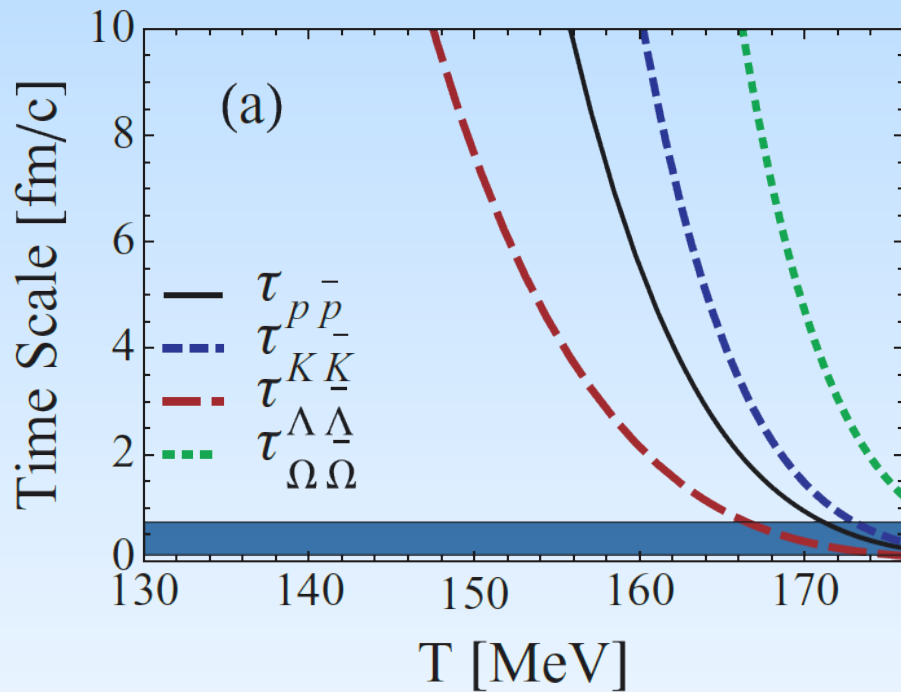
Calculation of Fuming Liu [Greiner et al., J. Phys. **G 31**, S725 (2005)]

Model parameters (2)

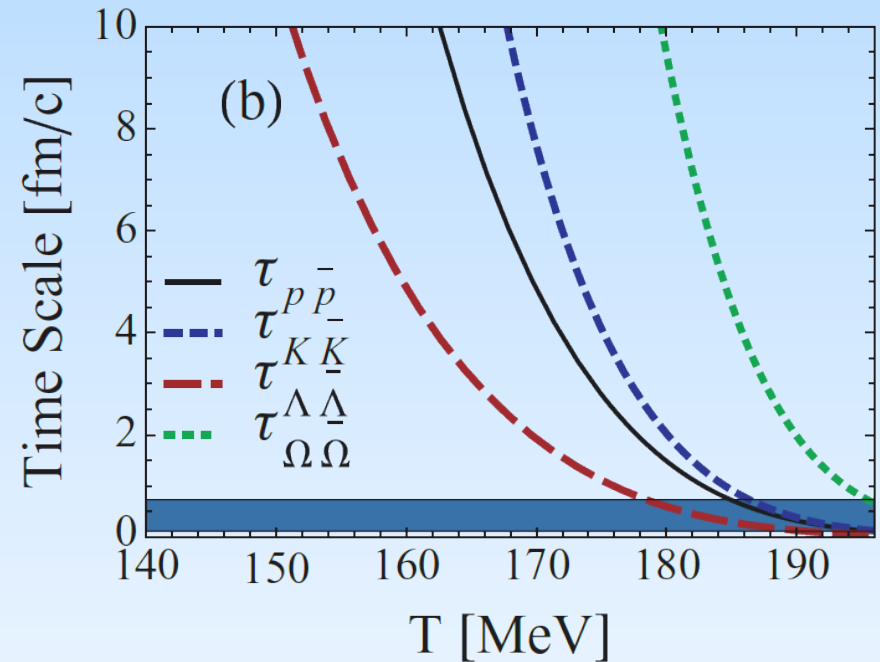
- Protons $\Gamma_{i,p\bar{p}} = 3 - 1000 \text{ MeV}$
- Kaons $\Gamma_{i,K\bar{K}} = 50 - 1700 \text{ MeV}$
- Lambdas $\Gamma_{i,\Lambda\bar{\Lambda}} = 3 - 250 \text{ MeV}$
- Omegas $\Gamma_{i,\Omega\bar{\Omega}} = 0.01 - 4 \text{ MeV}$

Chemical equilibration time

- $T_H=176$ MeV

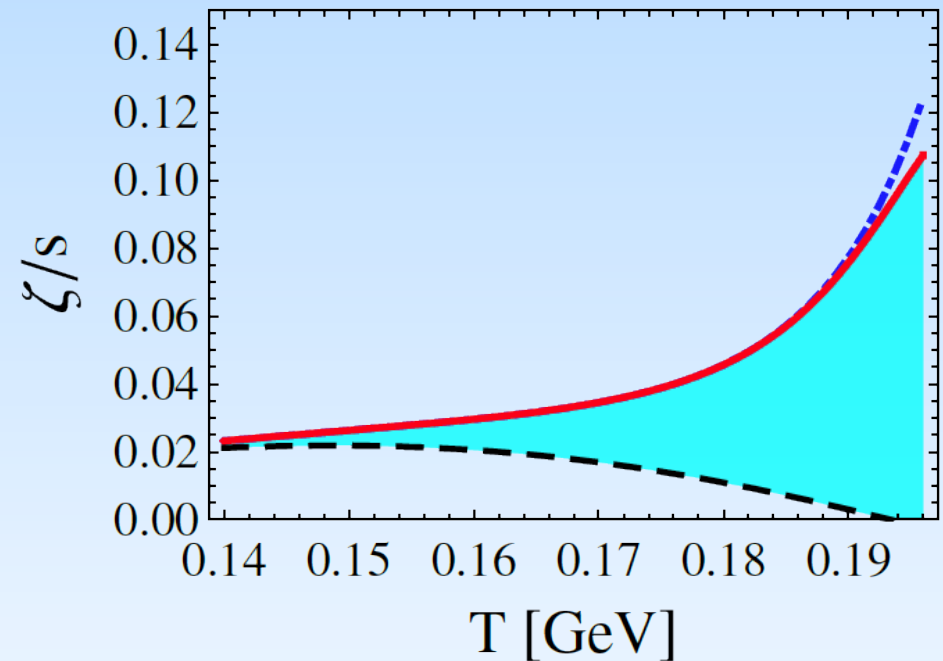
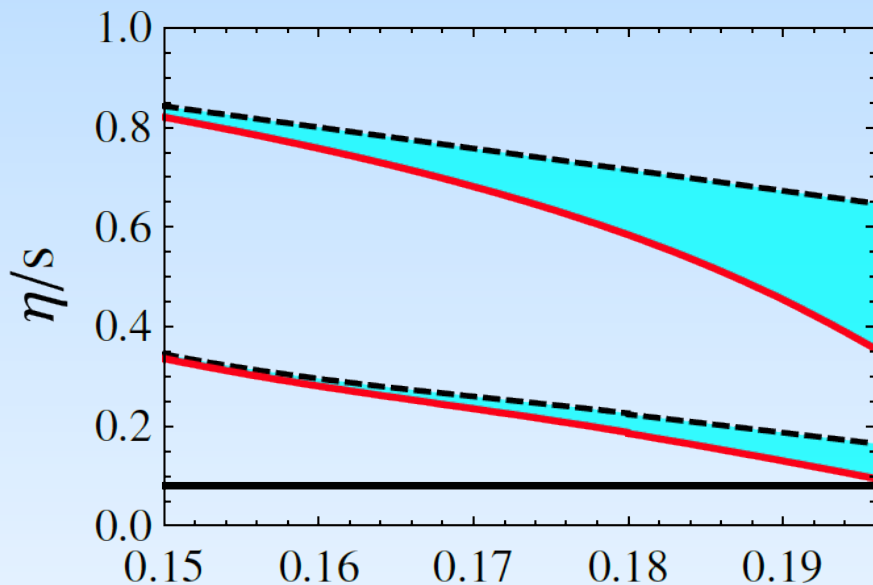


- $T_H=196$ MeV



Extra bonus

- Hagedorn states also affect viscosity



[Noronha-Hostler, Noronha, Greiner, Phys. Rev. Lett. **103**, 172302 (2009)]

Summary

- Collective behavior near critical temperature is important
- Chemical equilibration via Hagedorn states is a feasible mechanism
- Perhaps, there is no need for artificial “born in equilibrium” scenario
- Hagedorn states can drive viscosity low too
- Other implications (?)