

Chiral asymmetry in magnetized stellar matter

Igor Shovkovy

Arizona State University



Chiral effects at $B \neq 0$

- Chiral magnetic effect

$$\langle \vec{j} \rangle = \frac{e^2 \vec{B}}{2\pi^2} \mu_5$$

[Vilenkin, Phys. Rev. D 22, 3080 (1980)]

[Metlitski, Zhitnitsky, hep-ph/0505072]

[Newman, Son, hep-ph/0510049]

- Chiral separation effect

$$\langle \vec{j}_5 \rangle_{\text{free}} = -\frac{e\vec{B}}{2\pi^2} \mu$$

[Kharzeev, Zhitnitsky, arXiv:0706.1026]

[Kharzeev, McLerran, Warringa, arXiv:0711.0950]

[Fukushima, Kharzeev, Warringa, arXiv:0808.3382]

- Chiral magnetic wave

$$\delta\mu \rightarrow \delta j_5 \rightarrow \delta\mu_5 \rightarrow \delta j \rightarrow \delta\mu \rightarrow \dots$$

[Kharzeev, Yee, arXiv:1012.6026]

- etc

[Gorbar, Miransky, Shovkovy, arXiv:1101.4954]

[Burnier, Kharzeev, Liao, Yee, arXiv:1103.1307]

Spectrum at $B \neq 0$

- Dirac equation with massless fermions

$$\left[i\gamma^0 \partial_0 - i\vec{\gamma} \cdot (\vec{\nabla} + ie\vec{A}) \right] \Psi = 0$$

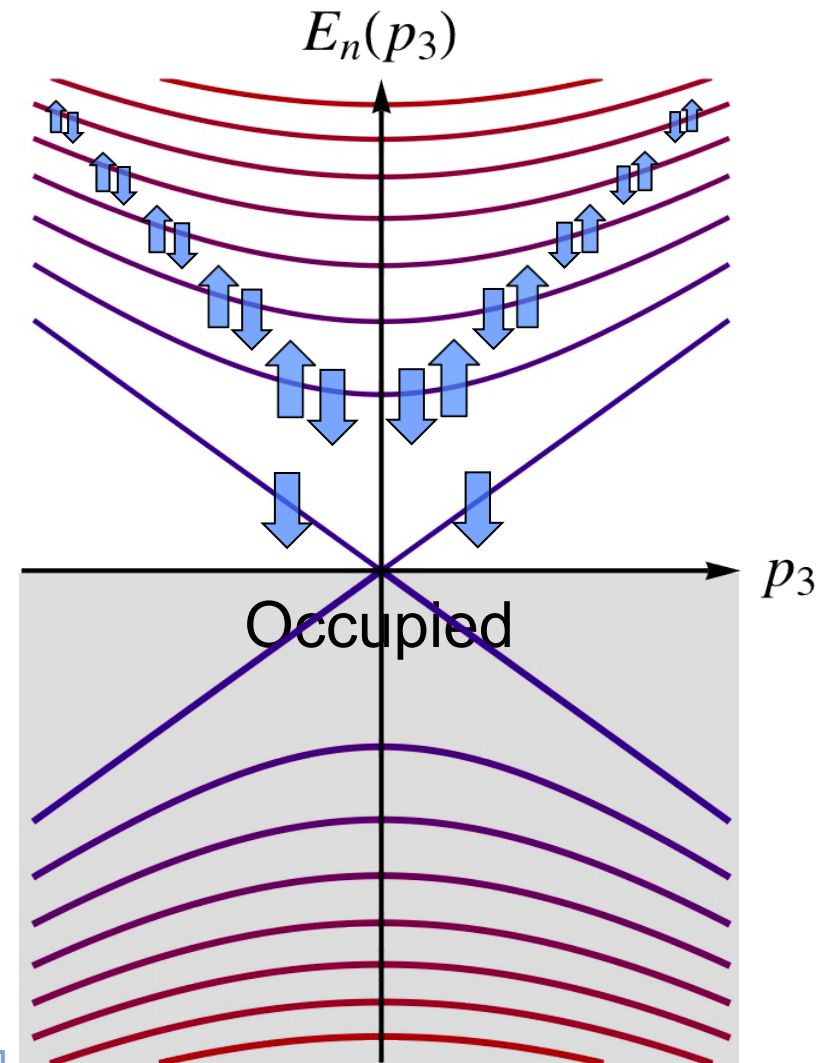
- Energy spectrum

$$E_n^{(3+1)}(p_3) = \pm \sqrt{2n|eB| + p_3^2}$$

$$s = \pm \frac{1}{2} \quad (\text{spin})$$

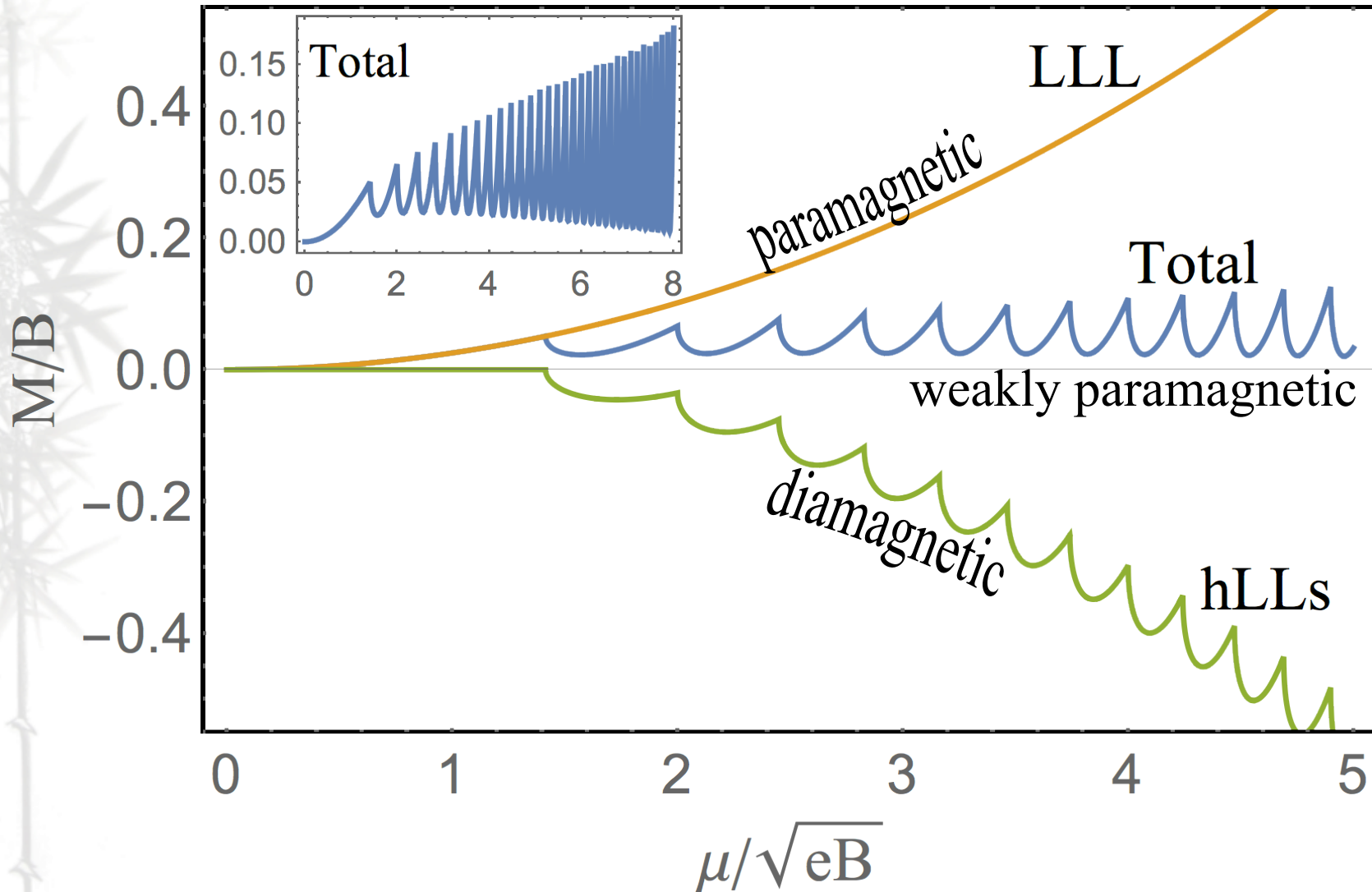
where $n = s + k + \frac{1}{2}$

$$k = 0, 1, 2, \dots \quad (\text{orbital})$$



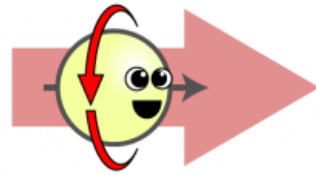
Relativistic magnetization

$$M = \frac{e\mu^2}{4\pi^2} + \frac{e}{2\pi^2} \sum_{n=1}^{n_{\max}} \left(\mu \sqrt{\mu^2 - 2neB} - 4n |eB| \ln \frac{\mu - \sqrt{\mu^2 - 2neB}}{\sqrt{2neB}} \right)$$

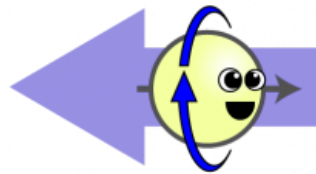


Spin vs. orbital motion

- Helicity/chirality of massless (ultrarelativistic) fermions is (\approx) conserved

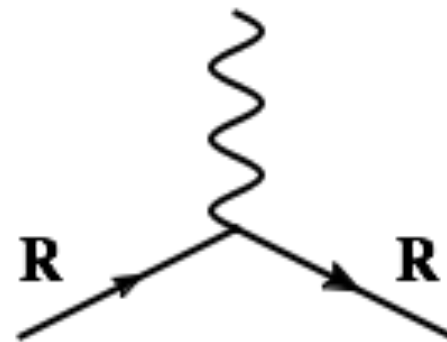
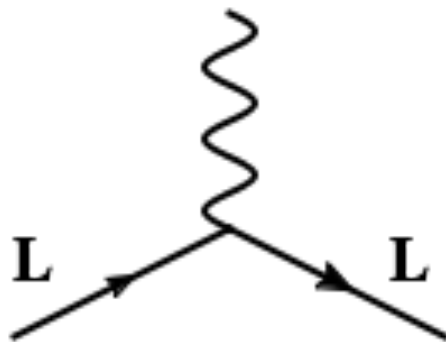


L-handed



R-handed

- Chirality does not change in elementary QED interactions



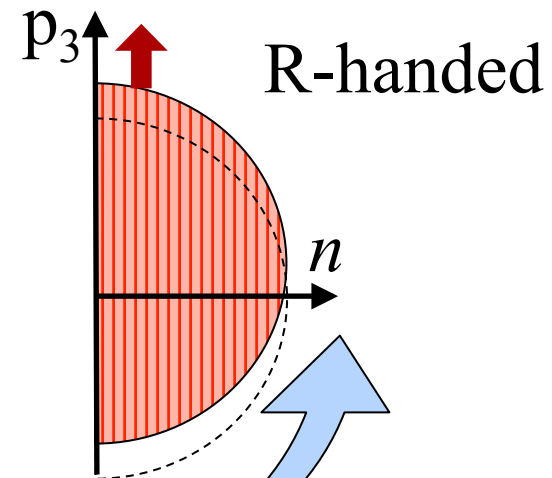
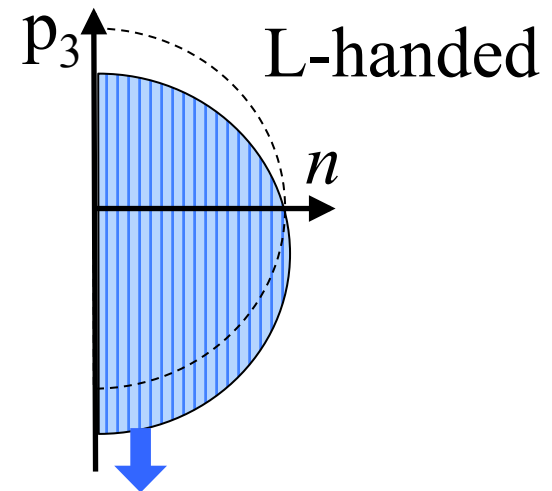
- LLL is spin polarized and chirally asymmetric
 - states with $p_3 < 0$ (and $s = \downarrow$) are L-handed
 - states with $p_3 > 0$ (and $s = \downarrow$) are R-handed
- When scattering off LLL, particles in higher LLs must “flip” both spin & momentum
- B-field \oplus interactions = chiral asymmetry
 - L-handed prefer $s = \downarrow$ and, thus, $p_3 < 0$
 - R-handed prefer $s = \downarrow$ and, thus, $p_3 > 0$

Chiral asymmetry

- Anticipated outcome: L- & R-handed Fermi surfaces shift in p_3 direction

Note: \mathbf{p}_\perp is not well-defined

p_\perp^2 is replaced by $2n|eB|$



- Ground state expectation value of the axial current (CSE)

$$\langle \bar{\psi} \gamma^3 \gamma^5 \psi \rangle = -\frac{eB}{2\pi^2} \mu \quad \text{with} \quad \vec{B} = (0, 0, B)$$

[Metlitski & Zhitnitsky, Phys. Rev. D **72**, 045011 (2005)]

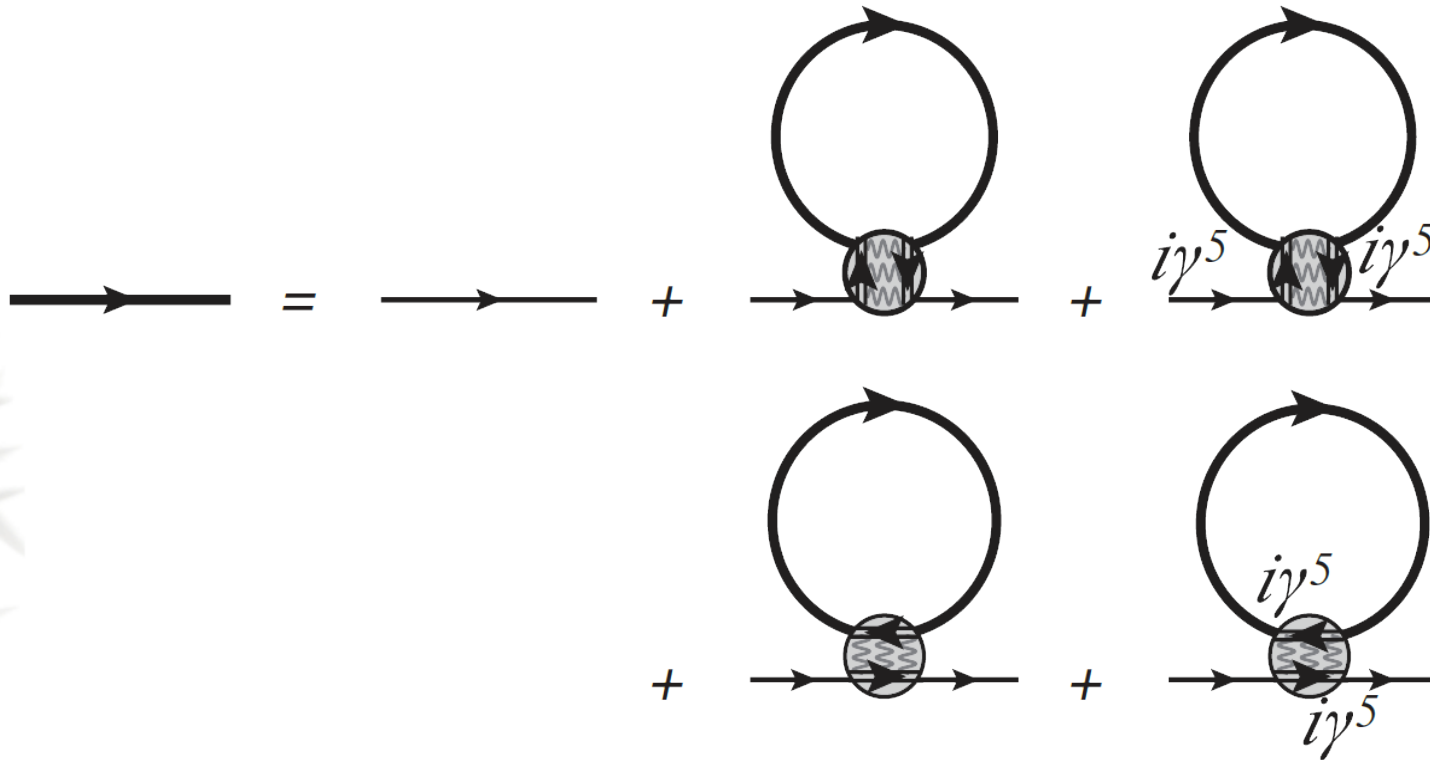
should induce a dynamical (chiral shift) parameter Δ associated with the condensate,

$$\delta L = \Delta \bar{\psi} \gamma^3 \gamma^5 \psi$$

($\Delta=0$ is not protected by any symmetry)

NJL model: quick check

- NJL model (local interaction)



- The equation for the chiral shift

$$\Delta = -\frac{1}{2} G_{\text{int}} \langle \bar{\psi} \gamma^3 \gamma^5 \psi \rangle \approx \frac{G_{\text{int}} eB}{4\pi^2} \mu$$

Chiral shift @ Fermi surface

- Chirality is \approx well-defined at Fermi surface ($|p_3| \gg m$)
- L-handed Fermi surface:

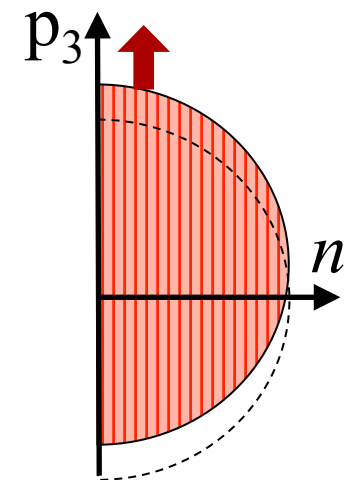
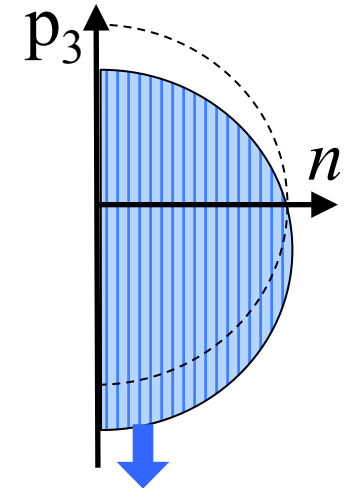
$$n > 0: \quad p_3 = +\sqrt{\left(\sqrt{\mu^2 - 2n|eB|} - s_{\perp} \Delta\right)^2 - m^2}$$

$$p_3 = -\sqrt{\left(\sqrt{\mu^2 - 2n|eB|} + s_{\perp} \Delta\right)^2 - m^2}$$

- R-handed Fermi surface:

$$n > 0: \quad p_3 = -\sqrt{\left(\sqrt{\mu^2 - 2n|eB|} - s_{\perp} \Delta\right)^2 - m^2}$$

$$p_3 = +\sqrt{\left(\sqrt{\mu^2 - 2n|eB|} + s_{\perp} \Delta\right)^2 - m^2}$$



[Gorbar, Miransky, Shovkovy, Phys. Rev. D **83**, 085003 (2011)]

- The result has the form

$$\bar{\Sigma}^{(1)}(p) = \gamma^3 \gamma^5 \Delta(p) + \gamma^0 \gamma^5 \mu_5(p)$$

Near Fermi surface ($p_0 \rightarrow 0$, $|\mathbf{p}| \rightarrow p_F$)

$$\Delta(p) \approx \frac{\alpha e B \mu}{\pi m^2} \left(\ln \frac{m^2}{2\mu(|\mathbf{p}| - p_F)} - 1 \right)$$

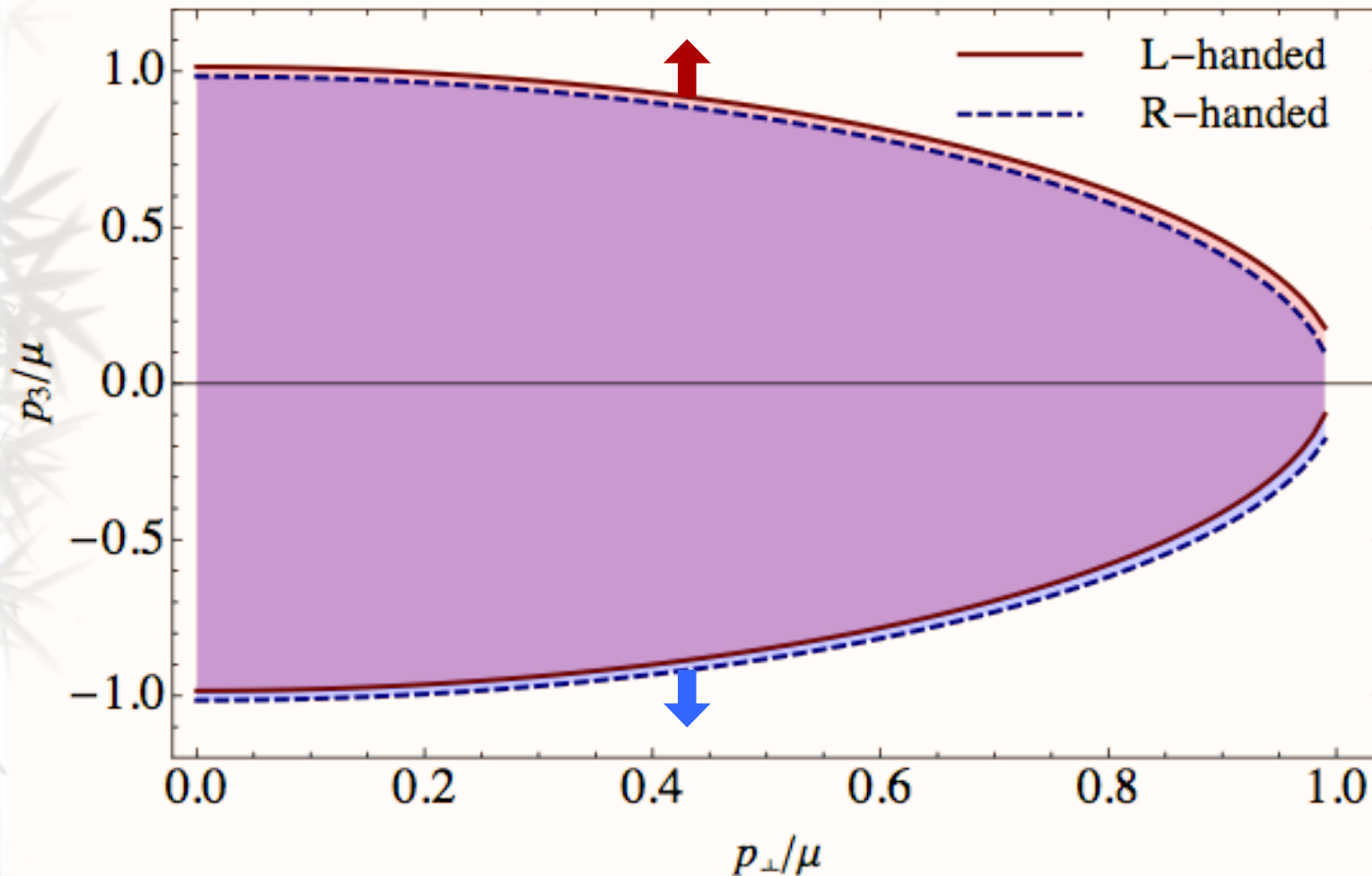
$$\mu_5(p) \approx -\frac{\alpha e B \mu}{\pi m^2} \frac{p_3}{p_F} \left(\ln \frac{m^2}{2\mu(|\mathbf{p}| - p_F)} - 1 \right)$$

[Gorbar, Miransky, Shovkovy, Wang, Phys. Rev. D **88**, 025043 (2013)]

Dispersion relations in QED

- Let us use the condition (for a small B)

$$\text{Det}\left[i\bar{S}^{-1}(p) + \bar{\Sigma}^{(1)}(p)\right] = 0$$



[Gorbar, Miransky, Shovkovy, Wang, Phys. Rev. D **88**, 025043 (2013)]

Self-energy in the Landau-level representation:

$$\bar{\Sigma}(p) = 2e^{-p_{\perp}^2 l^2} \sum_{n=0}^{\infty} (-1)^n \left(-\gamma^0 \delta\mu_n - \gamma^3 \gamma^5 \Delta_n - \gamma^0 \gamma^5 \mu_{5,n} + m_n + \dots \right) \left[P_- L_n - P_+ L_{n-1} \right] - \dots$$

where $\delta\mu_n$, Δ_n , $\mu_{5,n}$, ... are “projections” of the self-energy on the n th Landau level,

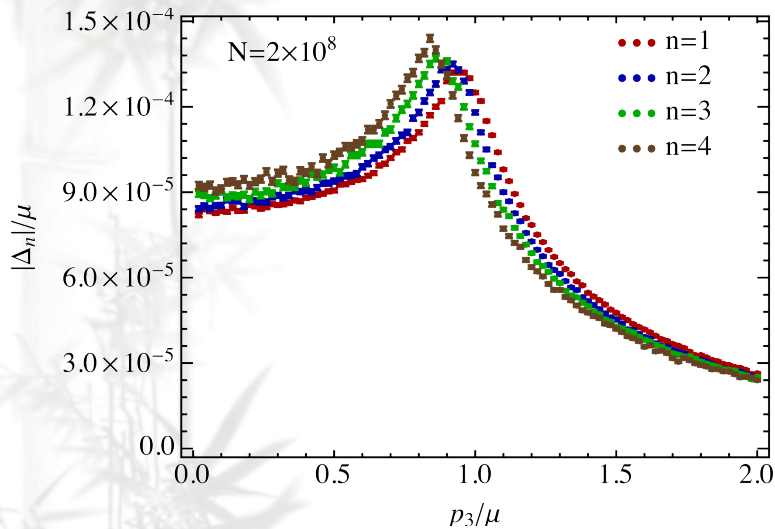
$$\Delta_n(p_0, p_3) = \frac{(-1)^n l^2}{8\pi} \int d^2 p_{\perp} e^{-p_{\perp}^2 l^2} \left[L_n + L_{n+1} \right] \text{Tr} \left[\gamma^0 \bar{\Sigma}(p) \right]$$

where

$$\bar{\Sigma}(p) = -4i\pi\alpha \int \frac{d^4 k}{(2\pi)^4} \gamma^{\mu} \bar{S}(k) \gamma^{\nu} D_{\mu\nu}(k-p)$$

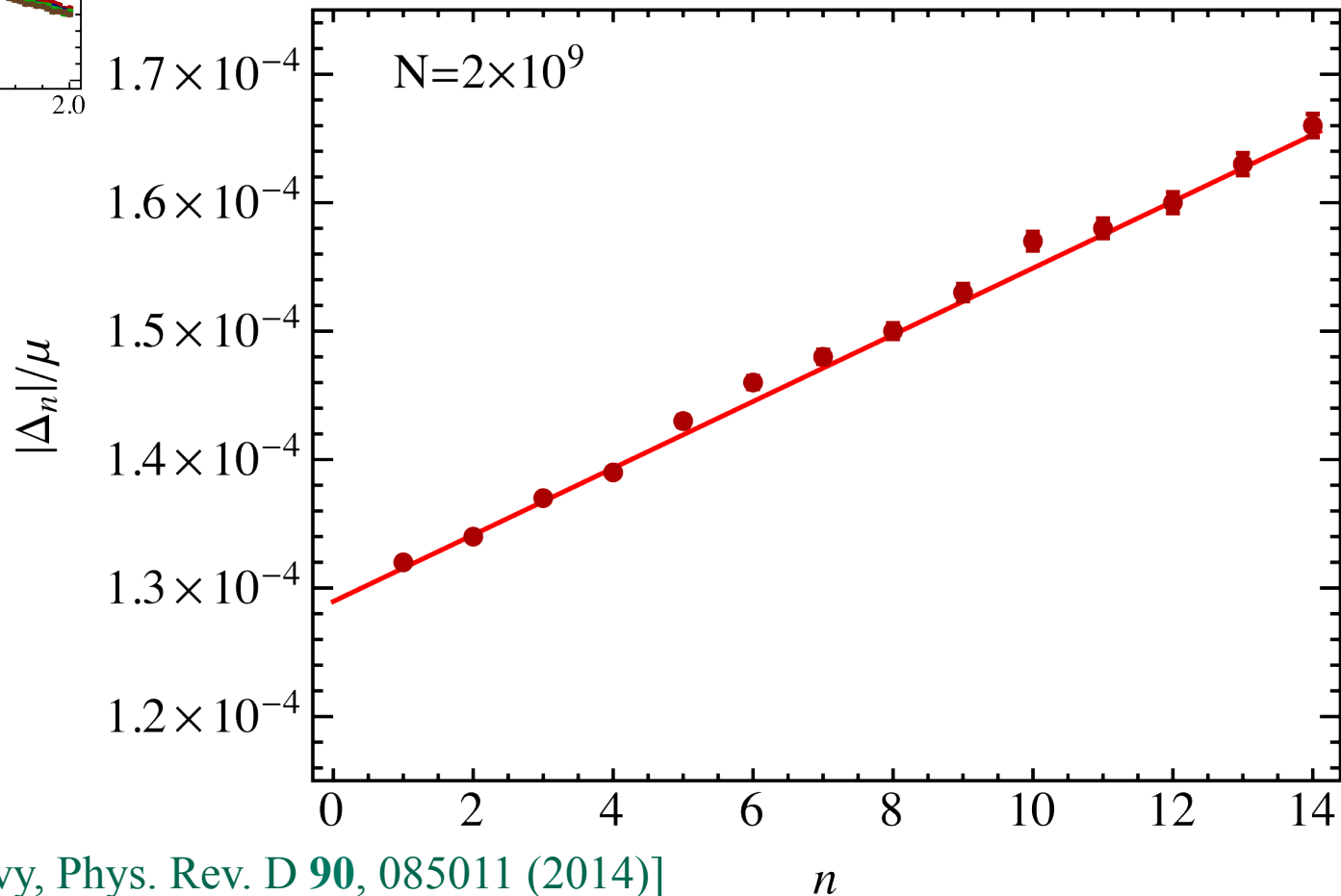
[Gorbar, Miransky, Shovkovy, Wang, Phys. Rev. D **88**, 025025 (2013)]

QED in strong field: Δ_n



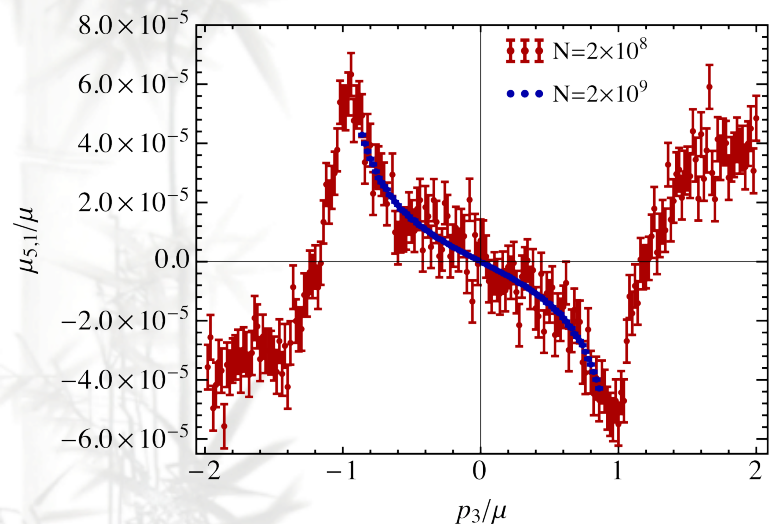
Model fit:

$$\Delta_n = -\frac{\alpha |eB|}{\mu} \left(0.53 + 0.32 \frac{|eB| n}{\mu^2} \right)$$



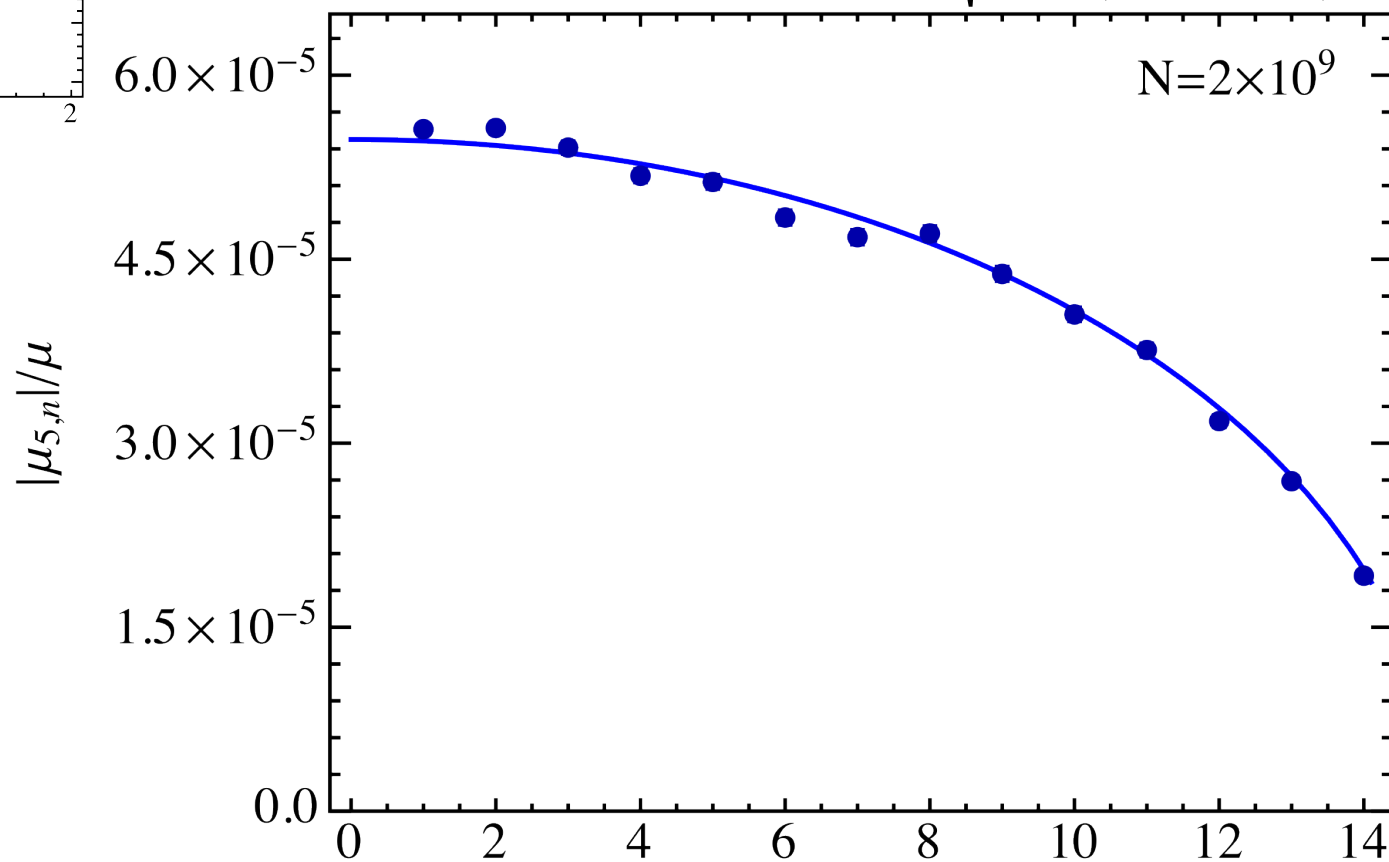
[Xia, Gorbar, Miransky, Shovkovy, Phys. Rev. D **90**, 085011 (2014)]

QED in strong field: $\mu_{5,n}$



Model fit:

$$\mu_{5,n} = -0.225 \frac{\alpha |eB|}{\mu} \sqrt{1 - \left(\frac{2n |eB|}{\mu^2} \right)^2}$$

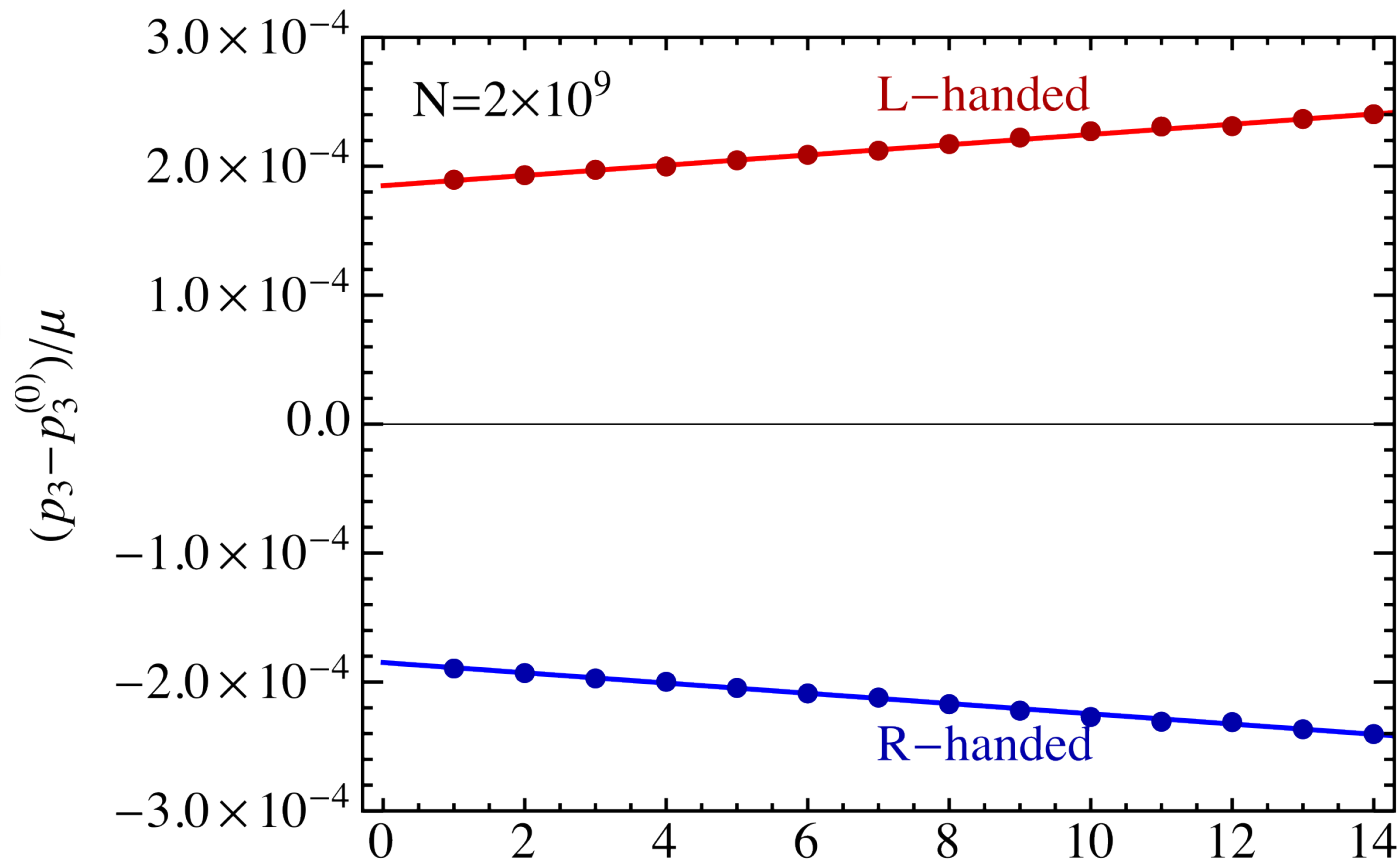


[Xia, Gorbar, Miransky, Shovkovy, Phys. Rev. D **90**, 085011 (2014)]

n

Model fit:

$$p_3 - p_3^{(0)} = \pm \frac{\alpha |eB|}{\mu} \left(0.76 + 0.49 \frac{|eB|n}{\mu^2} \right)$$



[Xia, Gorbar, Miransky, Shovkovy, Phys. Rev. D **90**, 085011ⁿ (2014)]

How large is the asymmetry?

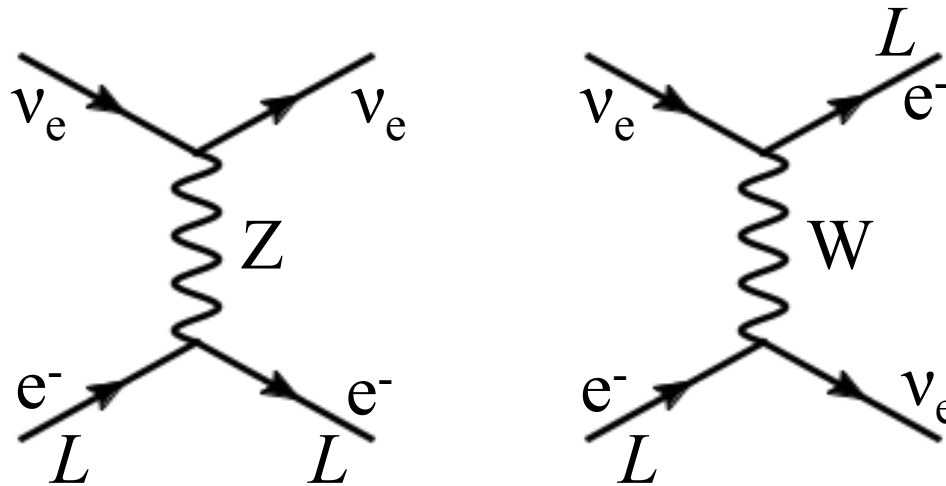
In QED:

$$\frac{\alpha |eB|}{\mu} \approx 0.4 \left(\frac{B}{10^{18} \text{ G}} \right) \left(\frac{100 \text{ MeV}}{\mu} \right) \text{ MeV}$$

In QCD:

$$\frac{\alpha_s |eB|}{\mu} \approx 10 \left(\frac{B}{10^{18} \text{ G}} \right) \left(\frac{400 \text{ MeV}}{\mu} \right) \text{ MeV}$$

- Neutrinos scatter off L-handed fermions only



- An asymmetric L-handed Fermi surface with

$$\delta p_3 \sim \alpha |eB|/\mu$$

should scatter ν_e 's more preferably in the direction opposite to \mathbf{B}

- Total momentum carrier by L-handed fermions

– QED ($B=10^{18}$ G and $\mu=100$ MeV):

$$P \sim N \delta p \sim 10^{55} \frac{\alpha |eB|}{\mu} \sim (1 \text{ km/s}) M_{Sun}$$

– QCD ($B=10^{18}$ G and $\mu=400$ MeV):

$$P \sim N \delta p \sim 10^{55} \frac{\alpha_s |eB|}{\mu} \sim (30 \text{ km/s}) M_{Sun}$$

- Pulsar kicks? Possible, but questions remain...

- LLL chiral asymmetry plus **interactions** generate chiral shift/asymmetry in higher LLs
- Chiral asymmetry shifts the L-handed and R-handed **Fermi surfaces** along **B**-field direction
- Chiral asymmetry can produce asymmetric **neutrino emission** and generate **pulsar kicks**
- The mechanism is more promising for quark stars, but may even affect all compact stars