ARIZONA STATE UNIVERSITY

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Collective modes in chiral (pseudo)relativistic matter

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CHIRAL MATTER

Chiral fermions

- *Massless* Dirac fermions:

$$\left(\gamma^0 p_0 - \vec{\gamma}\cdot\vec{\boldsymbol{p}}\right)\Psi = 0 \quad \Rightarrow \quad \frac{\vec{\Sigma}\cdot\vec{\boldsymbol{p}}}{|\vec{\boldsymbol{p}}|}\Psi = \mathrm{sign}(p_0)\gamma^5\Psi$$

For particles $(p_0 > 0)$: chirality = helicity

For antiparticles $(p_0 < 0)$: chirality = - helicity

- Massive Dirac fermions in *ultrarelativistic* regime
  - High temperature: $T >> m$
  - High density: $\mu >> m$

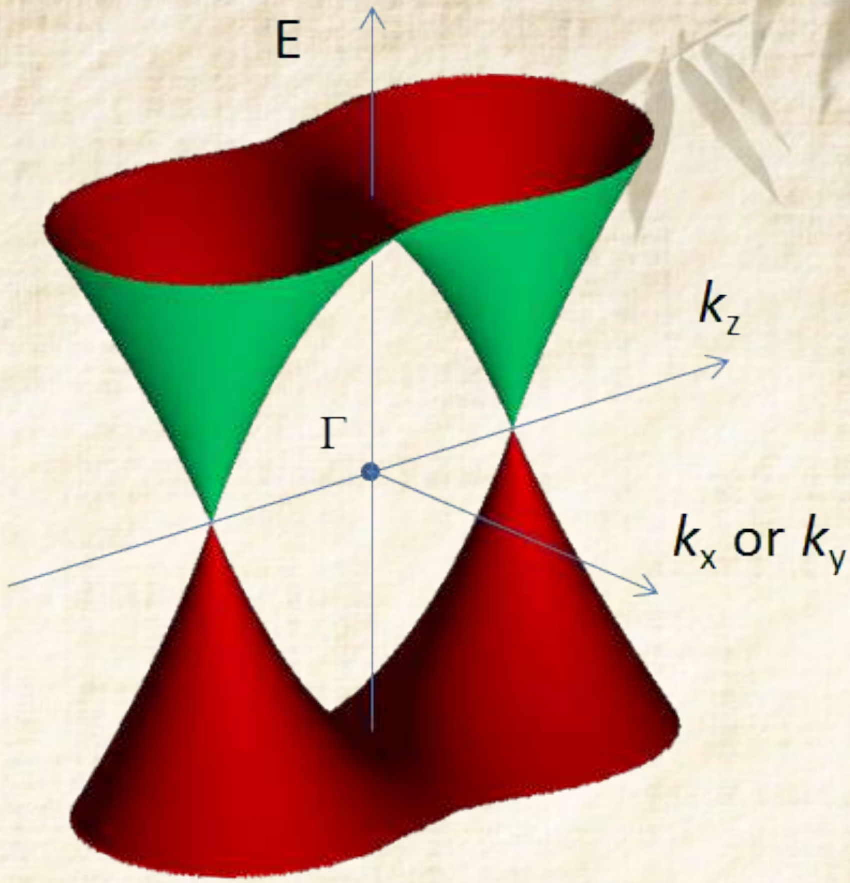
Chiral matter

- Matter made of chiral fermions with $n_\mathrm{L} \neq n_\mathrm{R}$
- Unlike the electric charge ($n_\mathrm{R} + n_\mathrm{L}$), the chiral charge ($n_\mathrm{R}$ - n_L) is **not** conserved

$$\frac{\partial(n_R + n_L)}{\partial t} + \vec{\nabla} \cdot \vec{j} = 0$$

$$\frac{\partial(n_R - n_L)}{\partial t} + \vec{\nabla} \cdot \vec{j}_5 = \frac{e^2 \vec{\boldsymbol{E}} \cdot \vec{\boldsymbol{B}}}{2\pi^2 c}$$

- The chiral symmetry is anomalous in quantum theory

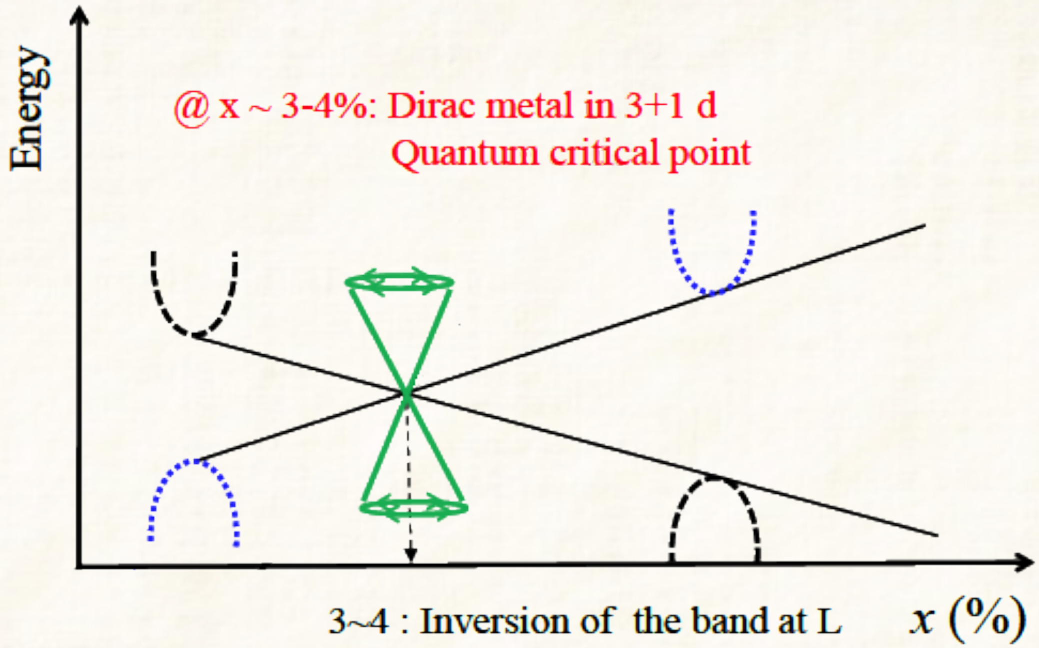


Credits: Borisenko et al., Phys. Rev. Lett. 113, 027603 (2014)

DIRAC & WEYL MATERIALS

Dirac semimetals

- Solid state materials with Dirac quasiparticles:
  - $Bi_{1-x}Sb_x$ alloy



- "New" 3D Dirac materials (ARPES):
  - Na_3Bi (Potassium bismuthide) [Liu et al., Science **343**, 864 (2014)]
  - Cd_3As_2 (Cadmium arsenide) [Neupane et al., Nature Commun. **5**, 3786 (2014)] [Borisenko et al., Phys. Rev. Lett. **113**, 027603 (2014)]

Dirac materials

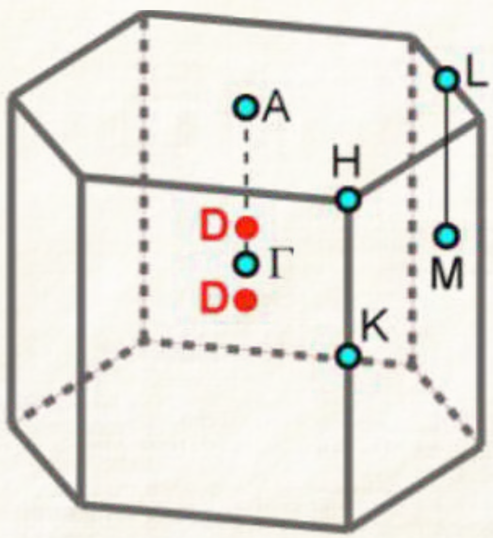
- $Bi_{1-x}Sb_x$ alloy (at $x \approx 4\%$)
- Na_3Bi
- Cd_3As_2
- $ZrTe_5$

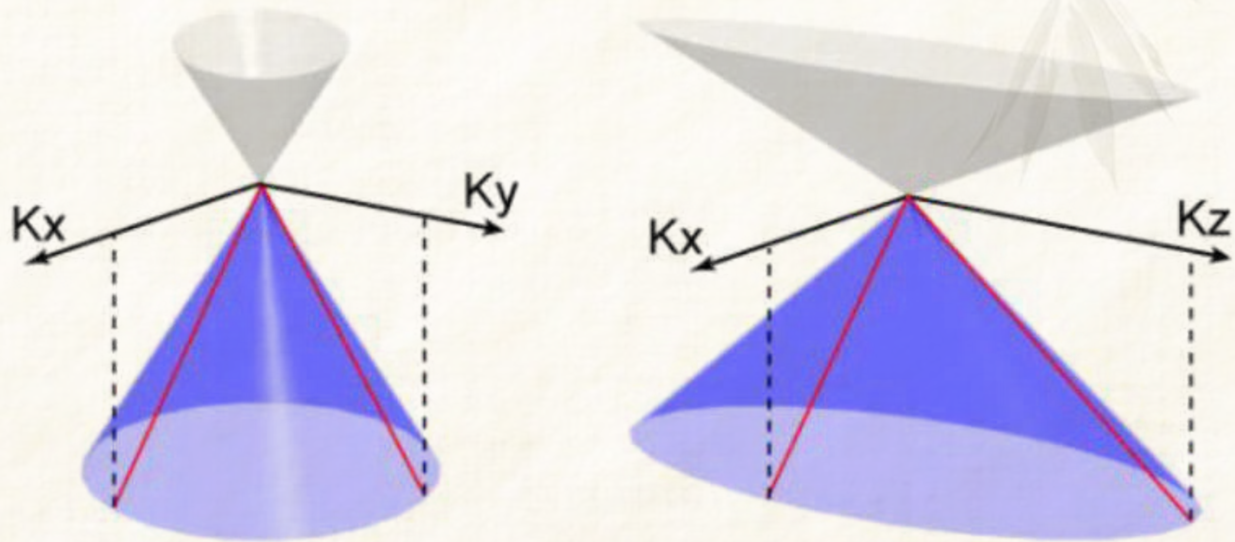
[Z. K. Liu et al., Science **343**, 864 (2014)]

[M. Neupane et al., Nature Commun. **5**, 3786 (2014)]

[S. Borisenko et al., Phys. Rev. Lett. **113**, 027603 (2014)]

[X. Li et al., Nature Physics **12**, 550 (2016)]





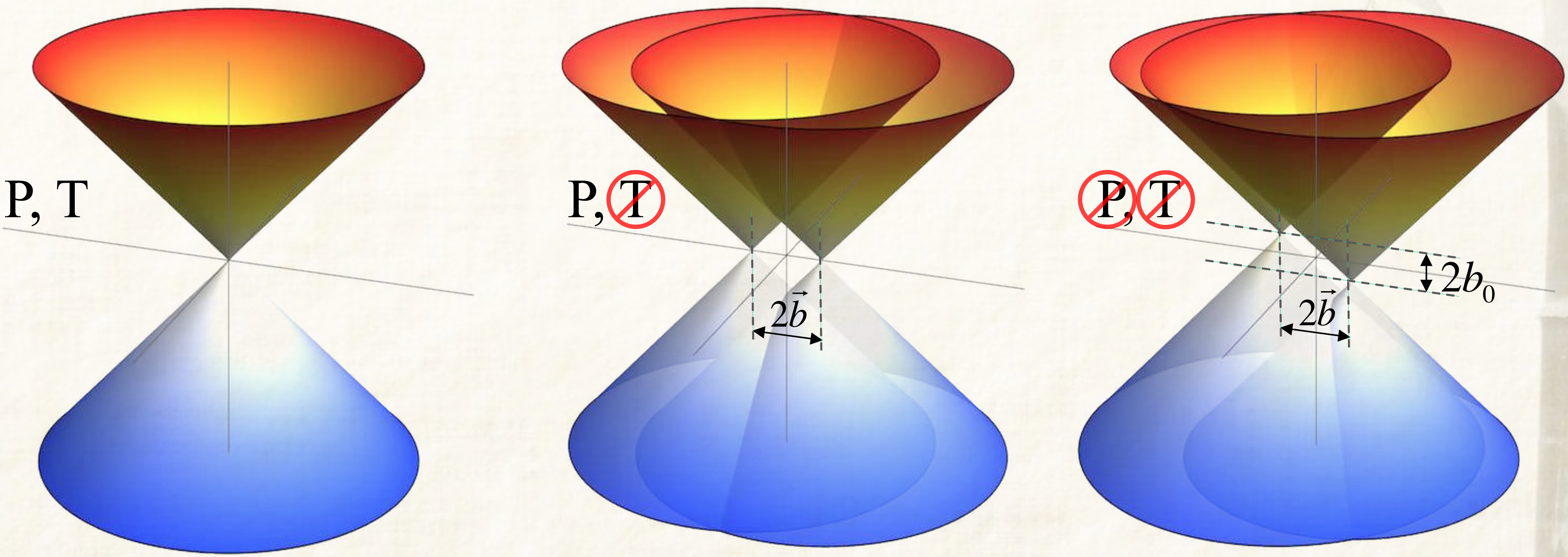
$$E = \sqrt{v_x^2 k_x^2 + v_y^2 k_y^2 + v_z^2 k_z^2}, \quad v_x \approx v_y \approx 3.74 \times 10^5 \,\mathrm{m/s}, \quad v_z \approx 2.89 \times 10^4 \,\mathrm{m/s}$$

Dirac vs. Weyl materials

- Low-energy Hamiltonian of a Dirac/Weyl material

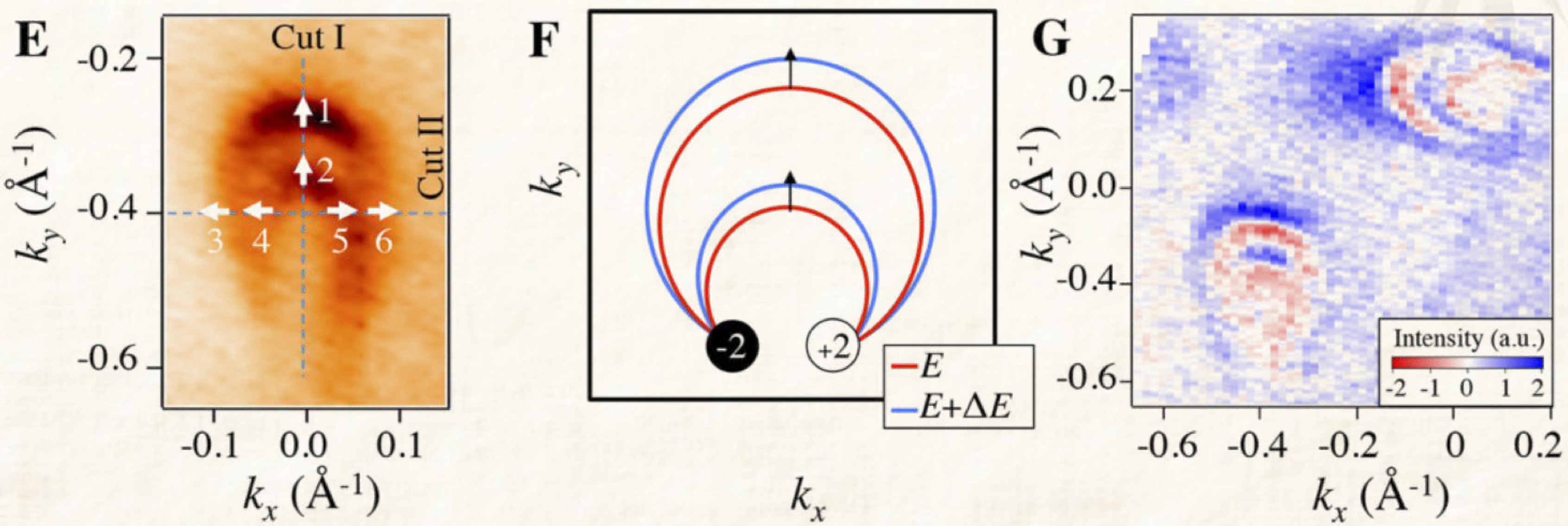
$$H = \int d^3\mathbf{r}\,\bar{\psi}\left[-iv_F\left(\vec{\gamma}\cdot\vec{\mathbf{p}}\right) - \overset{\cancel{T}}{\left(\vec{b}\cdot\vec{\gamma}\right)\gamma^5} + \overset{\cancel{P}}{b_0\gamma^0\gamma^5}\right]\psi$$

Dirac   Weyl



Weyl materials

- TaAs (tantalum arsenide) [S.-Y. Xu et al., Science 349, 613 (2015)] [B. Q. Lv et al., Phys. Rev. X5, 031013 (2015)]
- NbAs (niobium arsenide) [S.-Y. Xu et al., Nature Physics 11, 748 (2015)]
- TaP (tantalum phosphide) [S.-Y. Xu et al., Science Adv. 1, 1501092 (2015)]
- NbP (niobium phosphide) [I. Belopolski et al. arXiv:1509.07465]
- WTe_2 (tungsten telluride) [F. Y. Bruno et al., Phys. Rev. B 94, 121112 (2016)]



Low-energy Hamiltonian

- The Hamiltonian for massless Dirac fermions is given by

$$H_D = \begin{pmatrix} v_F\left(\vec{k}\cdot\vec{\sigma}\right) & 0 \\ 0 & -v_F\left(\vec{k}\cdot\vec{\sigma}\right) \end{pmatrix}$$

- This can we viewed as a combination of two Weyl fermions

$$H_\lambda = \lambda v_F\left(\vec{k}\cdot\vec{\sigma}\right)$$

where $\lambda = \pm 1$ is a chirality

The Weyl energy eigenstates are given by

$$\psi_k^\lambda = \frac{1}{\sqrt{2\epsilon_k}k_\perp}\begin{pmatrix} \sqrt{\epsilon_k + \lambda k_z}\, k_- \\ \lambda\sqrt{\epsilon_k - \lambda k_z}\, k_\perp \end{pmatrix}$$

They described particles of energy $\epsilon_k = v_F\sqrt{k_x^2 + k_y^2 + k_z^2}$

The mapping $k \to \psi_k^\lambda$ has a nontrivial topology

Berry connection & curvature

- Consider evolution from $\psi_{\boldsymbol{k}}$ to $\psi_{\boldsymbol{k}+\delta\boldsymbol{k}}$:

$$\langle\psi_{\boldsymbol{k}}|\psi_{\boldsymbol{k}+\delta\boldsymbol{k}}\rangle \approx 1 + \delta\boldsymbol{k}\cdot\langle\psi_{\boldsymbol{k}}|\boldsymbol{\nabla}_{\boldsymbol{k}}|\psi_{\boldsymbol{k}}\rangle \approx e^{i\boldsymbol{a}_{\boldsymbol{k}}\cdot\delta\boldsymbol{k}}$$

where $\boldsymbol{a}_{\boldsymbol{k}} = -i\langle\psi_{\boldsymbol{k}}|\boldsymbol{\nabla}_{\boldsymbol{k}}|\psi_{\boldsymbol{k}}\rangle$ is the Berry connection

- The Berry curvature is defined as follows:

$$\boldsymbol{\Omega}_k = \boldsymbol{\nabla}_k \times \boldsymbol{a}_{\boldsymbol{k}}$$

- Note the similarity with gauge fields, but $\boldsymbol{a}_{\boldsymbol{k}}$ and $\boldsymbol{\Omega}_k$ are defined in the momentum space
- It is convenient to define the Chern number (flux of $\boldsymbol{\Omega}_k$)

$$C = \frac{1}{2\pi}\oint \boldsymbol{\Omega}_k \cdot d\boldsymbol{S}_{\boldsymbol{k}}$$

- A nonzero (integer) Chern number indicates a nonzero (topological) charge inside the k-volume surrounded by the closed surface (Gauss's law)

Gauge theory vs. Berry effects

Gauge theory	Berry effects
Local at **coordinate** space	Local at **momentum** space
Gauge field $\vec{\boldsymbol{A}}$	Berry connection $\vec{\boldsymbol{a}}$
Magnetic field $\vec{\boldsymbol{B}} = \vec{\boldsymbol{\nabla}}_{\mathbf{r}} \times \vec{\boldsymbol{A}}$	Berry curvature $\vec{\boldsymbol{\Omega}} = \vec{\boldsymbol{\nabla}}_{\mathbf{k}} \times \vec{\mathbf{a}}$
Aharonov-Bohm phase $\oint d\vec{\boldsymbol{r}}\, \vec{\boldsymbol{A}}(\vec{\boldsymbol{r}})$	Berry phase $\oint d\vec{\mathbf{k}}\, \vec{\mathbf{a}}(\vec{\mathbf{k}})$
Magnetic charge (Dirac monopole) $\int d\,\vec{\boldsymbol{r}}\,(\vec{\boldsymbol{\nabla}}_{\mathbf{r}} \cdot \vec{\boldsymbol{B}}) = const$	Berry monopole $\int d\,\vec{\mathbf{k}}\,(\vec{\boldsymbol{\nabla}}_{\mathbf{k}} \cdot \vec{\boldsymbol{\Omega}}) = const$

Berry curvature for Weyl fermions

- In the case of Weyl fermions,

$$\mathbf{\Omega}_k = \lambda \frac{\vec{k}}{2k^3}$$

(Note: this looks like a field of a monopole at $\vec{k} = 0$)

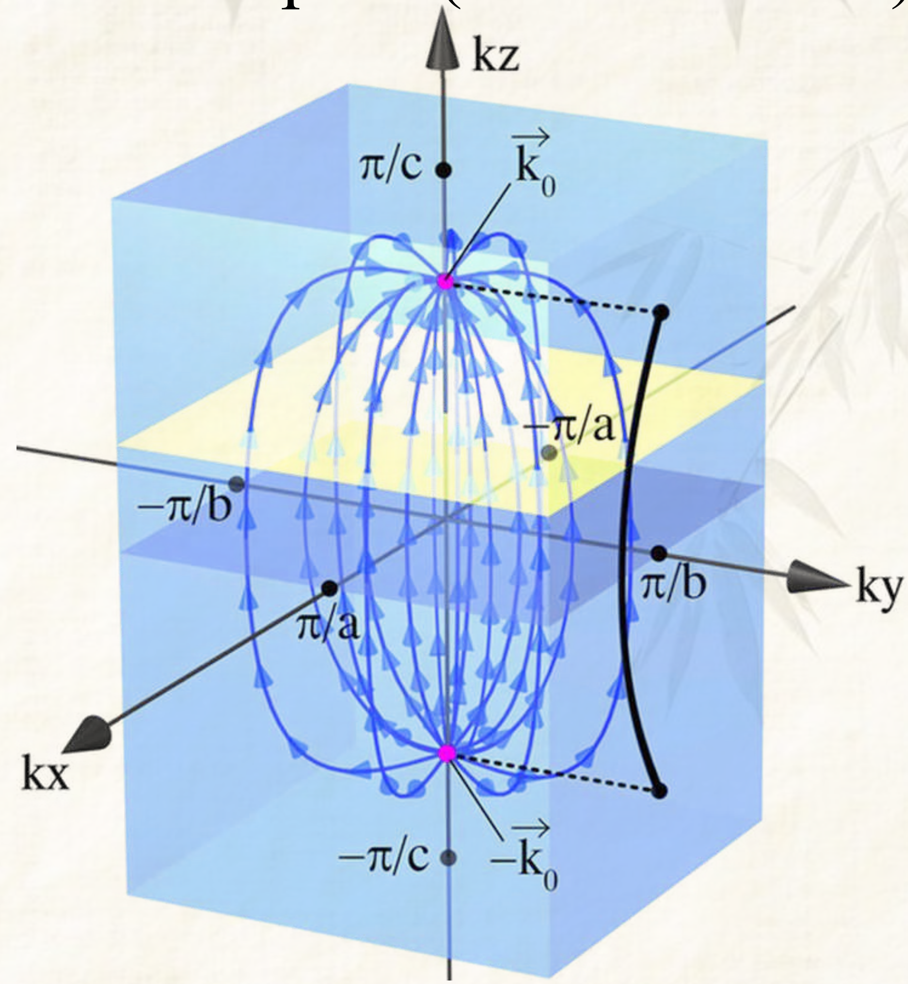
- Let us calculate the total flux of $\mathbf{\Omega}_k$-field through the spherical surface of radius K with the center at $\vec{k} = 0$

$$C = \frac{1}{2\pi} \oint \lambda \frac{\vec{k}}{2k^3} \cdot \frac{\vec{k}}{k} k^2 \sin\theta \, d\theta d\varphi = \lambda = \pm 1$$

- Thus, the electronic structure of massless Weyl fermions is characterized by a topological monopole at $\vec{k} = 0$
- Is the Berry monopole just a mathematical curiosity?
- Are there any observable consequences?

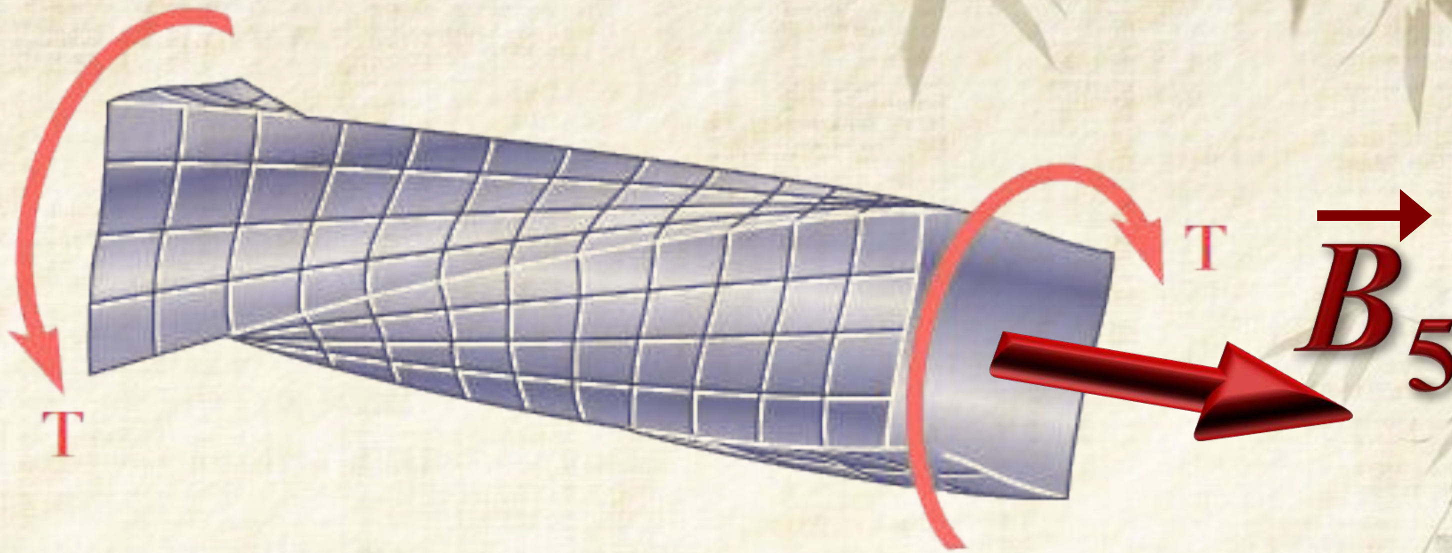
Weyl fermions on a lattice

- In solid state physics, the momentum space (Brillouin zone) is compact
- A closed surface around a single Weyl node is also a closed surface (of opposite orientation) around a the rest of the Brillouin zone
- Flipping surface orientation changes the sign of the flux
- There must be an opposite charge somewhere in the rest of the zone



[Morimoto & Nagaosa, Scientific Reports **6**, 19853 (2016)]

- Thus, Weyl fermions come in pairs of opposite chirality
  [Nielsen & Ninomiya, Nucl. Phys. B **193**, 173 (1981); B **185**, 20 (1981)]



PSEUDO-ELECTROMAGNETIC FIELDS

[Zubkov, Annals Phys. **360**, 655 (2015)]
[Cortijo, Ferreiros, Landsteiner, Vozmediano. Phys. Rev. Lett. **115**, 177202 (2015)]
[Grushin, Venderbos,Vishwanath, Ilan, arXiv:1607.04268]
[Cortijo, Kharzeev, Landsteiner, Vozmediano, arXiv:1607.03491]
[Pikulin, Chen, Franz, Phys. Rev. X **6**, 041021 (2016)]

Strain in Weyl materials

- Strains affect low-energy quasiparticles in Weyl materials

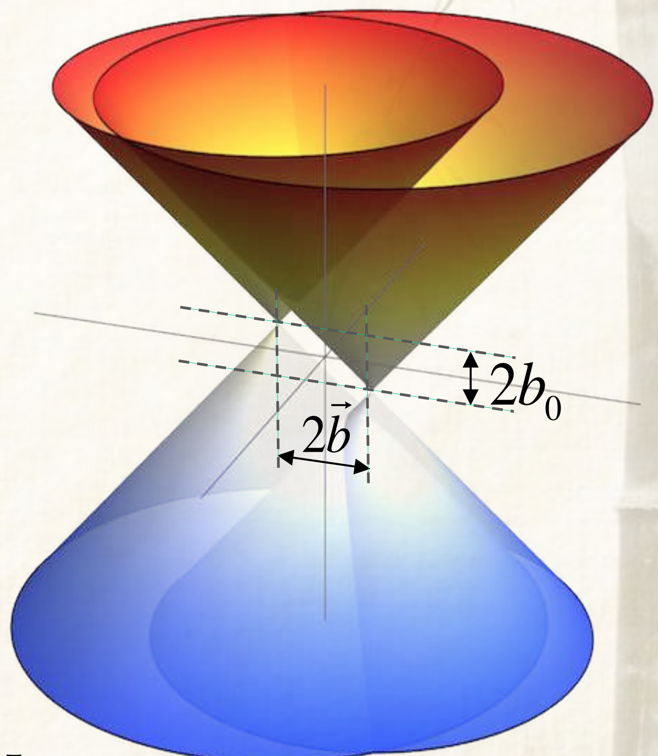
$$H = \int d^3\mathbf{r}\,\bar{\psi}\left[-iv_F\left(\vec{\gamma}\cdot\vec{\mathbf{p}}\right)-\left(\vec{b}+\vec{A}_5\right)\cdot\vec{\gamma}\,\gamma^5+\left(b_0+A_{5,0}\right)\gamma^0\gamma^5\right]\psi$$

where the components of the chiral gauge fields are

$$A_{5,0} \propto b_0\left|\vec{b}\right|\partial_{||}u_{||}$$

$$A_{5,\perp} \propto \left|\vec{b}\right|\partial_{||}u_{\perp}$$

$$A_{5,||} \propto \alpha\left|\vec{b}\right|^2\partial_{||}u_{||}+\beta\sum_i \partial_i u_i$$



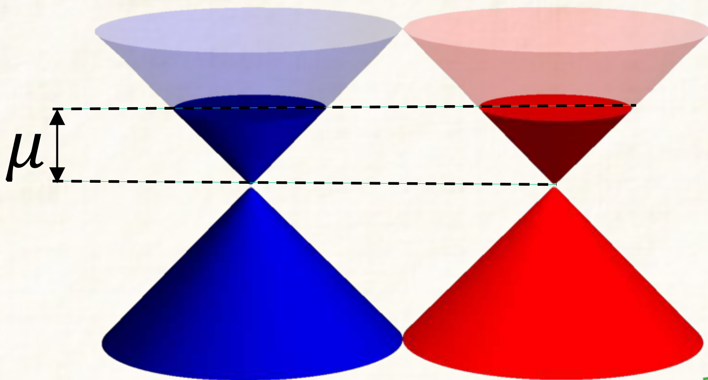
The associated pseudo-EM fields are

$$\vec{B}_5 = \vec{\nabla}\times\vec{A}_5 \quad \text{and} \quad \vec{E}_5 = -\vec{\nabla}A_0 - \partial_t\vec{A}_5$$

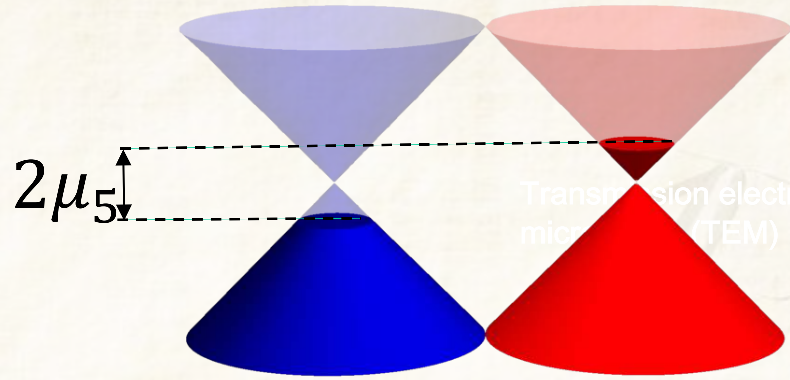
Equilibrium vs. nonequilibrium

$$\boxed{\mu_\lambda = \mu + \lambda\mu_5}$$

1) $\mu_5 = b_0 = 0$

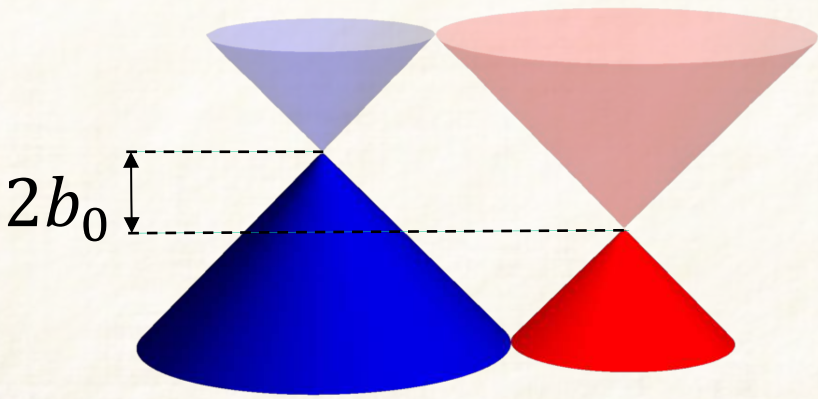


2) $\mu = b_0 = 0$

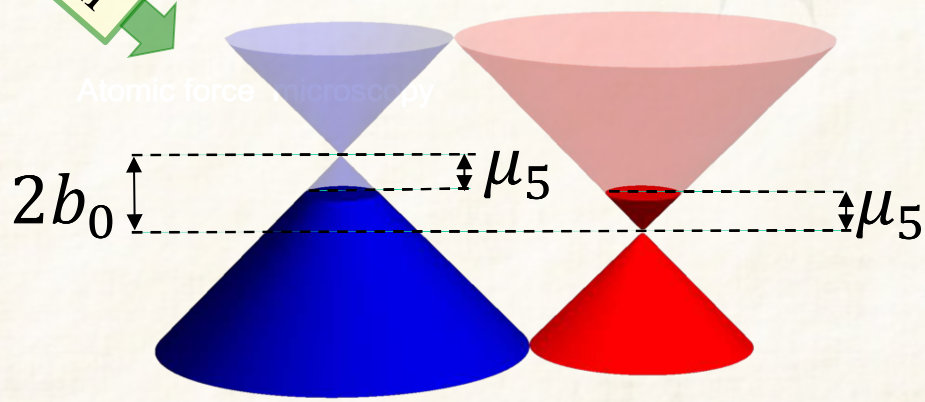


Equilibrium

3) $\mu = \mu_5 = 0$



4) $\mu_5 = b_0,\ \mu = 0$



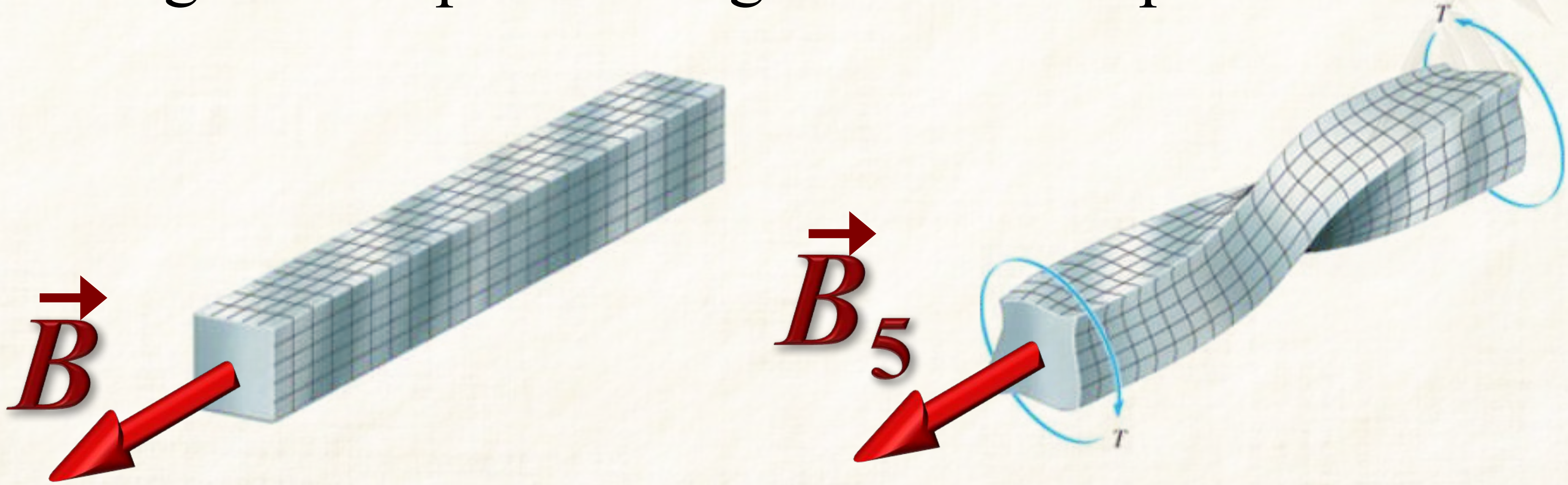
Chiral effects in Weyl materials

- Any qualitative properties of Weyl materials directly sensitive to b_0 and $\vec{b}$?
- Some proposals:
  - Anomalous Hall effect
  - Anomalous Alfven waves
  - Strain/torsion induced CME
  - Strain/torsion induced quantum oscillations
  - Strain/torsion dependent resistance
  - etc.
- Spectrum of chiral (pseudo-)magnetic plasmons

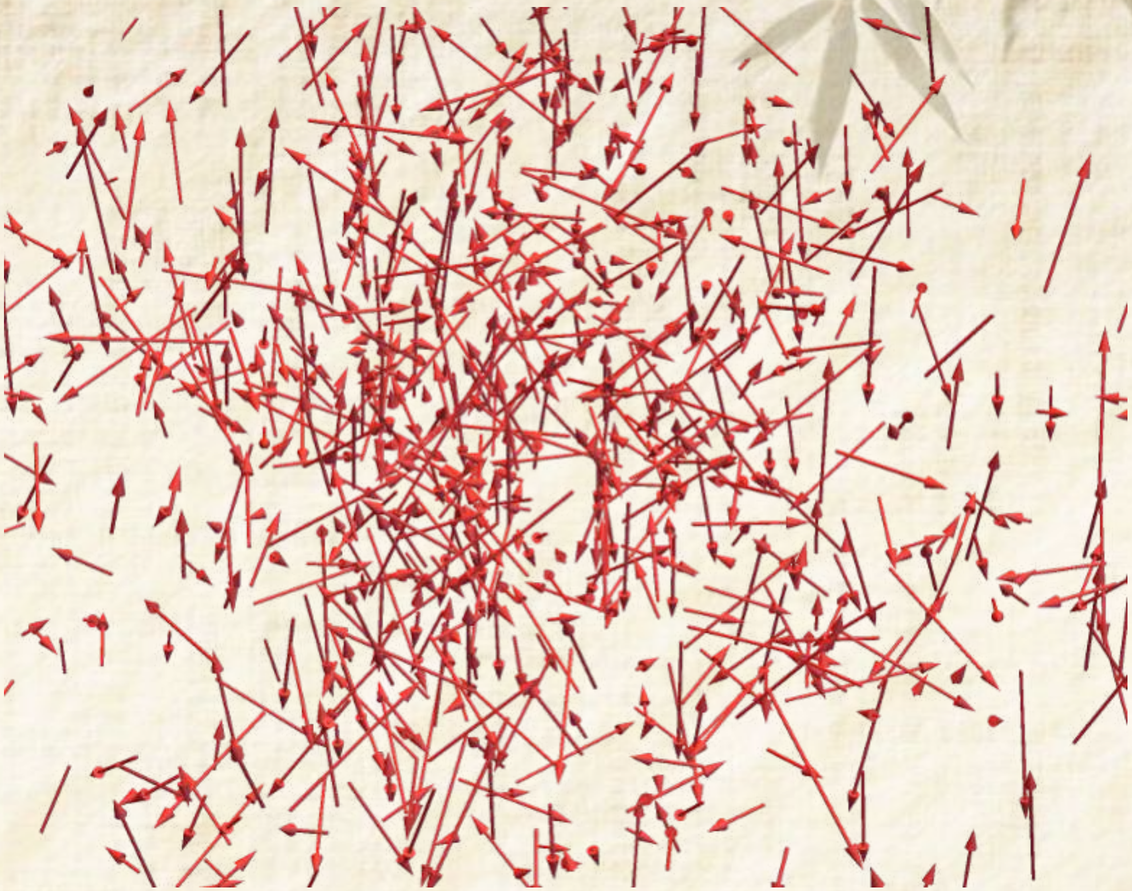
[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. Lett. **118**, 127601 (2017)]
[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. B **95**, 115202 (2017)]

General question

- What are the properties of plasmons in magnetized chiral material with $b_0 \neq 0$ and $\vec{b} \neq 0$?
- Chiral matter ($\mu_R \neq \mu_L$)
  - This is the case in equilibrium when $b_0 \neq 0$ ($\mu_5 = -eb_0$)
- Magnetic or pseudomagnetic field is present



- In general, $\mathbf{E}_\lambda = \mathbf{E} + \lambda \mathbf{E}_5$ and $\mathbf{B}_\lambda = \mathbf{B} + \lambda \mathbf{B}_5$



MODEL & METHOD

[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. Lett. **118**, 127601 (2017)]
[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. B **95**, 115202 (2017)]

Chiral kinetic theory

- Kinetic equation:

[Son and Yamamoto, Phys. Rev. D **87**, 085016 (2013)]
[Stephanov and Yin, Phys. Rev. Lett. **109**, 162001 (2012)]

$$\frac{\partial f_\lambda}{\partial t} + \frac{\left[e\tilde{\mathbf{E}}_\lambda + \frac{e}{c}(\mathbf{v}\times\mathbf{B}_\lambda) + \frac{e^2}{c}(\tilde{\mathbf{E}}_\lambda\cdot\mathbf{B}_\lambda)\boldsymbol{\Omega}_\lambda\right]\cdot\boldsymbol{\nabla}_{\mathbf{p}} f_\lambda}{1+\frac{e}{c}(\mathbf{B}_\lambda\cdot\boldsymbol{\Omega}_\lambda)} + \frac{\left[\mathbf{v} + e(\tilde{\mathbf{E}}_\lambda\times\boldsymbol{\Omega}_\lambda) + \frac{e}{c}(\mathbf{v}\cdot\boldsymbol{\Omega}_\lambda)\mathbf{B}_\lambda\right]\cdot\boldsymbol{\nabla}_{\mathbf{r}} f_\lambda}{1+\frac{e}{c}(\mathbf{B}_\lambda\cdot\boldsymbol{\Omega}_\lambda)} = 0$$

where $\tilde{\mathbf{E}}_\lambda = \mathbf{E}_\lambda - (1/e)\boldsymbol{\nabla}_{\mathbf{r}}\epsilon_{\mathbf{p}}$, $\quad \mathbf{v} = \boldsymbol{\nabla}_{\mathbf{p}}\epsilon_{\mathbf{p}}$,

$$\epsilon_{\mathbf{p}} = v_F p\left[1 - \frac{e}{c}(\mathbf{B}_\lambda\cdot\boldsymbol{\Omega}_\lambda)\right]$$

and $\boldsymbol{\Omega}_\lambda = \lambda\hbar\frac{\hat{\mathbf{p}}}{2p^2}$ is the Berry curvature

Current and chiral anomaly

- The definitions of density and current are

$$
\begin{aligned}
\rho_\lambda &= e\int \frac{d^3p}{(2\pi\hbar)^3}\left[1+\frac{e}{c}(\mathbf{B}_\lambda\cdot\boldsymbol{\Omega}_\lambda)\right]f_\lambda,\\
\mathbf{j}_\lambda &= e\int \frac{d^3p}{(2\pi\hbar)^3}\left[\mathbf{v}+\frac{e}{c}(\mathbf{v}\cdot\boldsymbol{\Omega}_\lambda)\mathbf{B}_\lambda+e(\tilde{\mathbf{E}}_\lambda\times\boldsymbol{\Omega}_\lambda)\right]f_\lambda\\
&\quad +e\boldsymbol{\nabla}\times\int\frac{d^3p}{(2\pi\hbar)^3}f_\lambda\epsilon_{\mathbf{p}}\boldsymbol{\Omega}_\lambda,
\end{aligned}
$$

They satisfy the following anomalous relations:

$$
\frac{\partial\rho_5}{\partial t}+\boldsymbol{\nabla}\cdot\mathbf{j}_5 = \frac{e^3}{2\pi^2\hbar^2 c}\Big[(\mathbf{E}\cdot\mathbf{B})+(\mathbf{E}_5\cdot\mathbf{B}_5)\Big] \quad ✔
$$

$$
\frac{\partial\rho}{\partial t}+\boldsymbol{\nabla}\cdot\mathbf{j} = \frac{e^3}{2\pi^2\hbar^2 c}\Big[(\mathbf{E}\cdot\mathbf{B}_5)+(\mathbf{E}_5\cdot\mathbf{B})\Big] \quad ✘
$$

Consistent definition of current

- Additional Bardeen-Zumino term is needed,

$$\delta j^\mu = \frac{e^3}{4\pi^2\hbar^2 c}\epsilon^{\mu\nu\rho\lambda} A^5_\nu F_{\rho\lambda}$$

- In components,

$$\delta\rho = \frac{e^3}{2\pi^2\hbar^2 c^2}(\mathbf{A}^5\cdot\mathbf{B})$$

$$\delta\mathbf{j} = \frac{e^3}{2\pi^2\hbar^2 c} A^5_0\mathbf{B} - \frac{e^3}{2\pi^2\hbar^2 c}(\mathbf{A}^5\times\mathbf{E})$$

- Its role and implications:
  - Electric charge is conserved locally ($\partial_\mu J^\mu = 0$)
  - Anomalous Hall effect is reproduced
  - CME vanishes in equilibrium ($\mu_5 = -eb_0$)

Collective modes

We search for plane-wave solutions with

$$\mathbf{E}' = \mathbf{E}e^{-i\omega t + i\mathbf{k}\cdot\mathbf{r}}, \;\; \mathbf{B}' = \mathbf{B}e^{-i\omega t + i\mathbf{k}\cdot\mathbf{r}}$$

and the distribution function $f_\lambda = f_\lambda^{(\mathrm{eq})} + \delta f_\lambda$,

where

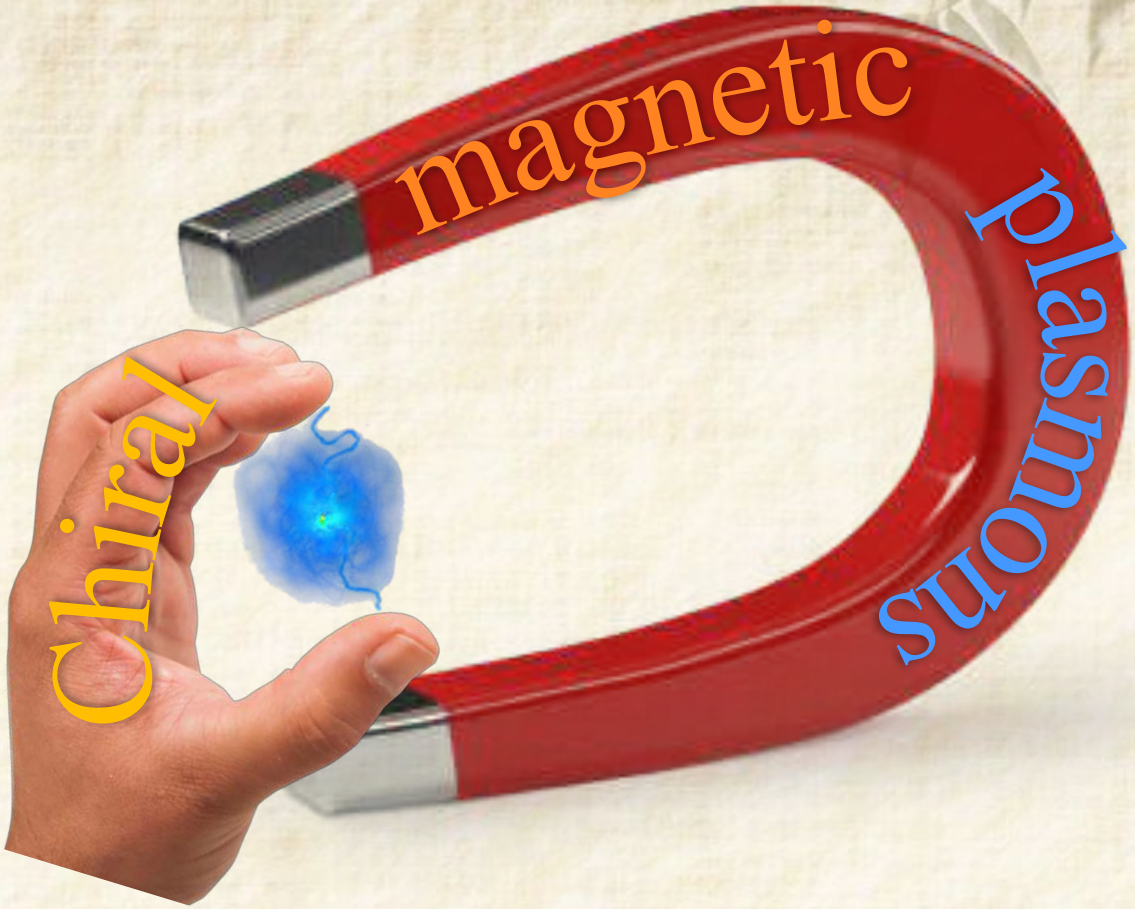
$$\delta f_\lambda = f_\lambda^{(1)} e^{-i\omega t + i\mathbf{k}\cdot\mathbf{r}}$$

The polarization vector & susceptibility tensor:

$$P^m = i\frac{J'^m}{\omega} = \chi^{mn}E'^n$$

The plasmon dispersion relations follow from

$$\det\left[\left(\omega^2 - c^2k^2\right)\delta^{mn} + c^2k^mk^n + 4\pi\omega^2\chi^{mn}\right] = 0$$



RESULTS

[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. Lett. **118**, 127601 (2017)]
[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. B **95**, 115202 (2017)]

Chiral magnetic plasmons

Non-degenerate plasmon frequencies @ k=0:

$$\omega_l = \Omega_e, \qquad \omega_{\mathrm{tr}}^{\pm} = \Omega_e \sqrt{1 \pm \frac{\delta\Omega_e}{\Omega_e}}$$

where the Langmuir frequency is

$$\Omega_e \equiv \sqrt{\frac{4\alpha}{3\pi\hbar^2}\left(\mu^2 + \mu_5^2 + \frac{\pi^2 T^2}{3}\right)}$$

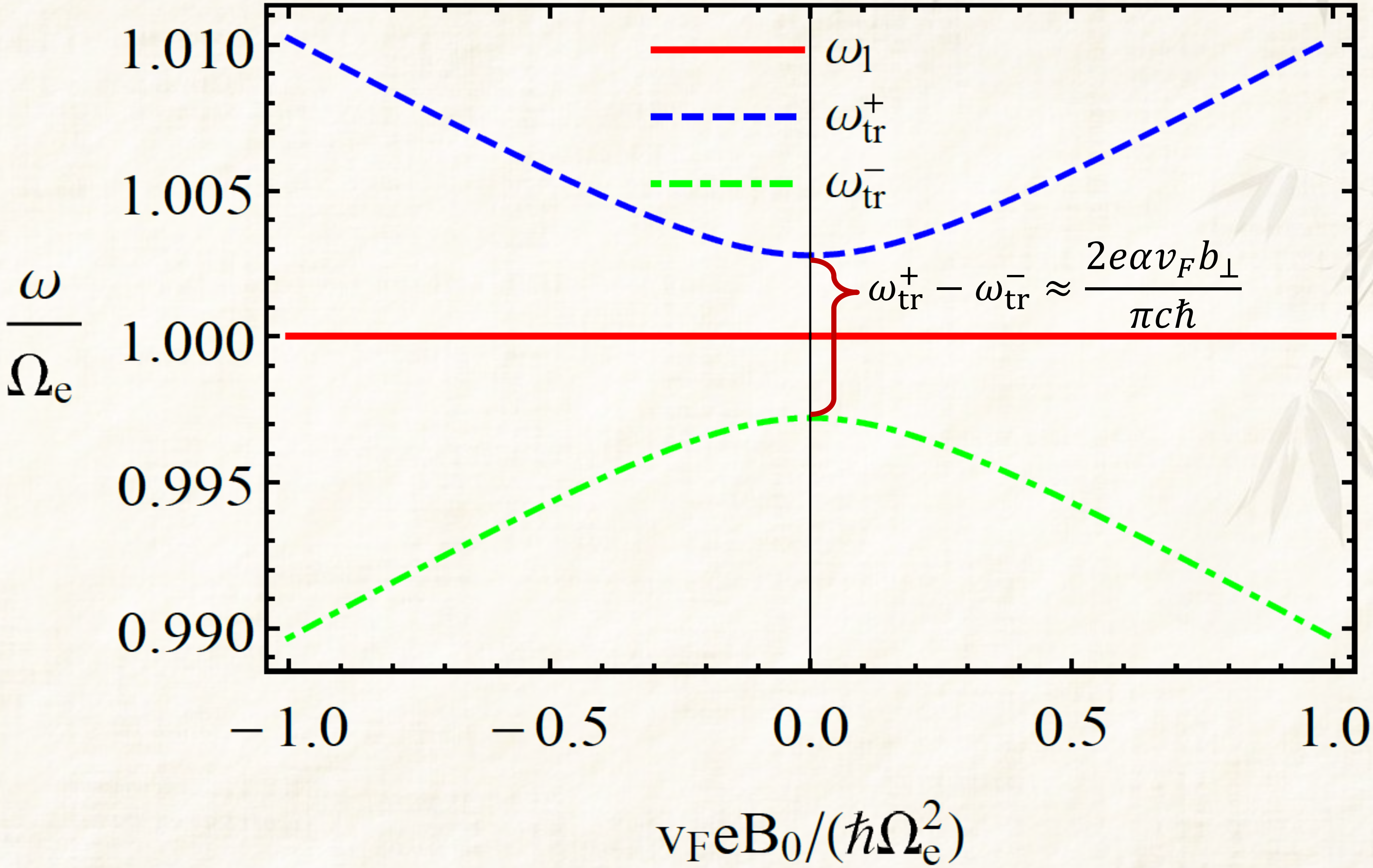
and

$$\delta\Omega_e = \frac{2e\alpha v_F}{3\pi c\hbar^2}\left\{9\hbar^2 b_\perp^2 + \left[\frac{2v_F}{\Omega_e^2}(B_0\mu + B_{0,5}\mu_5) - 3\hbar b_\parallel - \frac{v_F\hbar^2}{4T}\sum_{\lambda=\pm} B_{0,\lambda} F\left(\frac{\mu_\lambda}{T}\right)\right]^2\right\}^{1/2}$$

[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. Lett. **118**, 127601 (2017)]
[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. B **95**, 115202 (2017)]

Plasmon frequencies, $\vec{B} \perp \vec{b}$

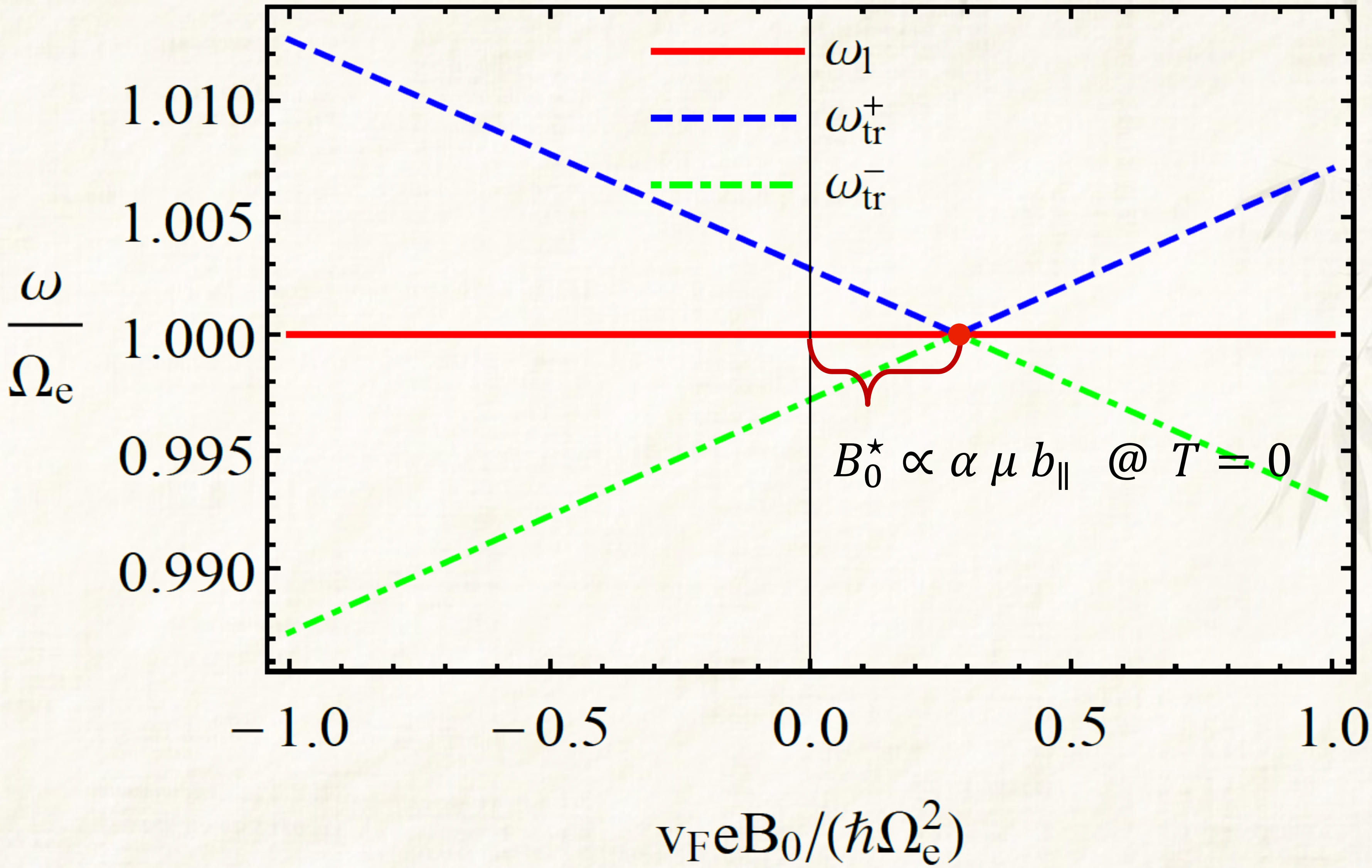
$$b_\perp = 0.2\hbar\Omega_e/e$$



[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. Lett. **118**, 127601 (2017)]

Plasmon frequencies, $\vec{B} \parallel \vec{b}$

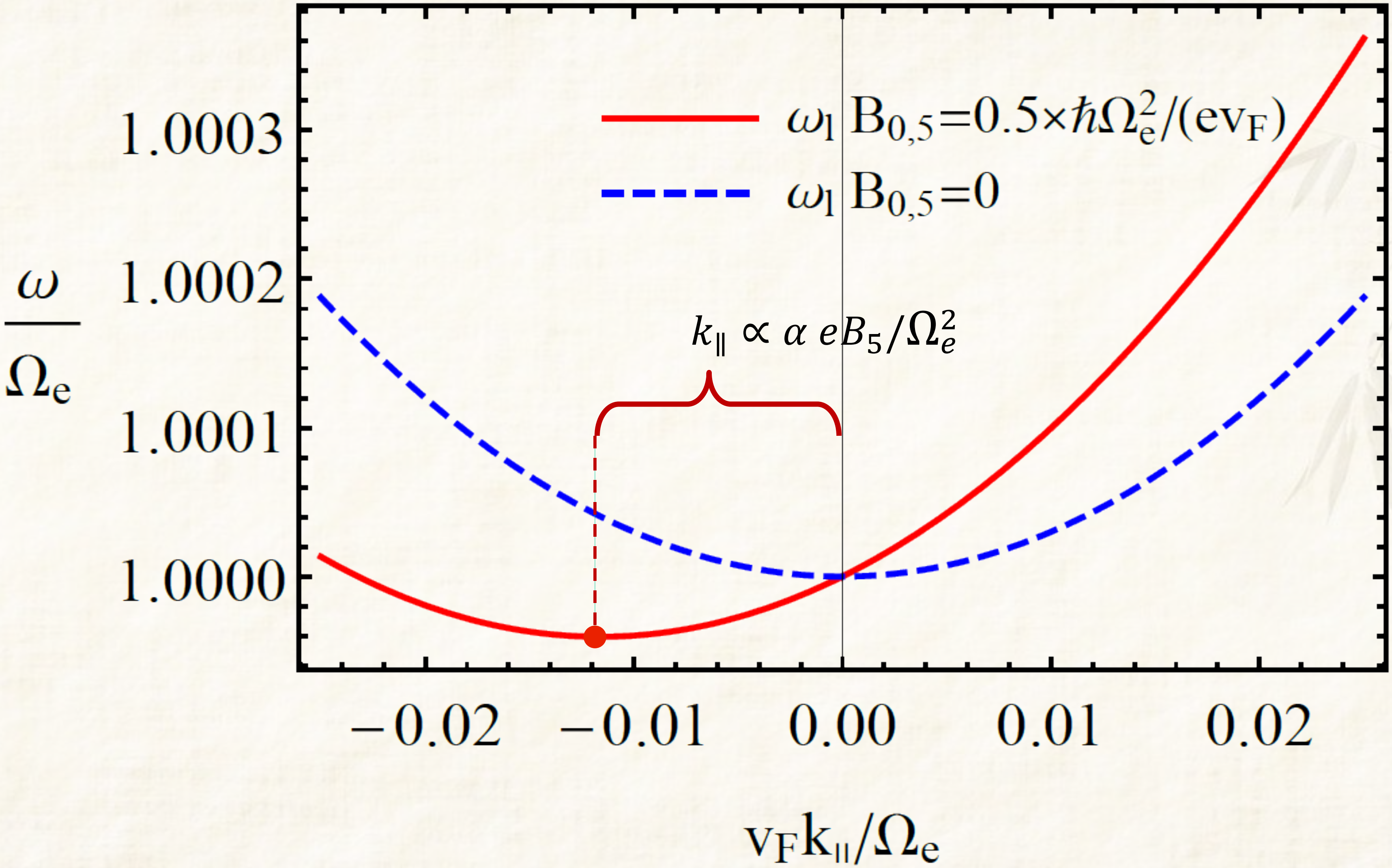
$b_{\parallel} = 0.2\hbar\Omega_e/e$



[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. Lett. **118**, 127601 (2017)]

Plasmons with $\vec{k} \neq 0$, $\vec{k} \parallel \vec{B}_5$

- The longitudinal mode is sensitive to $\vec{B}_5$

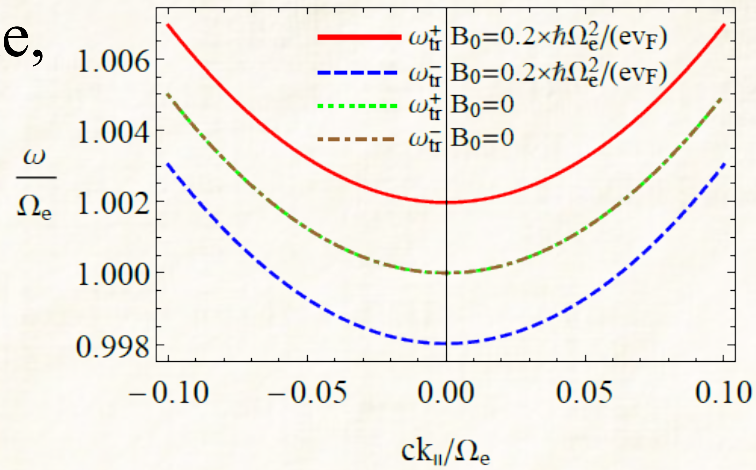


[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. B **95**, 115202 (2017)]

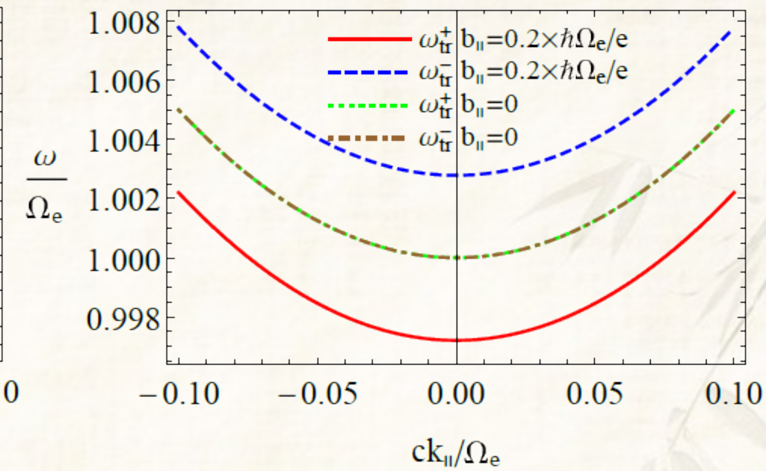
Plasmons with $\vec{k} \neq 0$, $\vec{k} \parallel \vec{B}, \vec{B}_5$

The transverse modes split (in different ways) when (i) $\vec{B} \neq 0$ & $\mu \neq 0$, or (ii) $\vec{B}_5 \neq 0$ & $\mu_5 \neq 0$, or (iii) $b_{\parallel} \neq 0$, or (iv) $b_{\perp} \neq 0$.

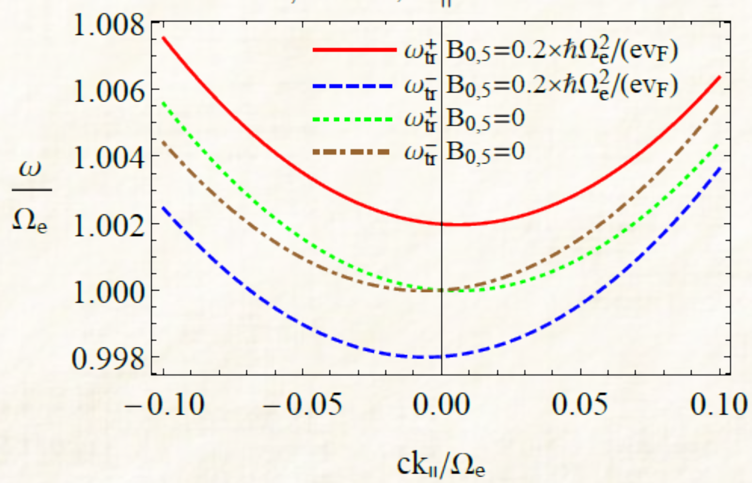
For example,



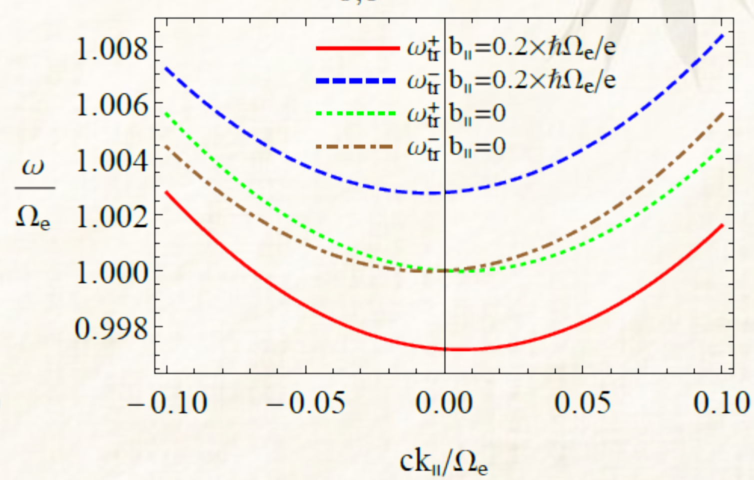
(a) $\mu = \sqrt{3\pi/(4\alpha)}\hbar\Omega_e$, $\mu_5 = 0$, $B_{0,5} = 0$, $b_{\parallel} = 0$



(c) $\mu = \sqrt{3\pi/(4\alpha)}\hbar\Omega_e$, $\mu_5 = 0$, $B_0 = 0$, $B_{0,5} = 0$

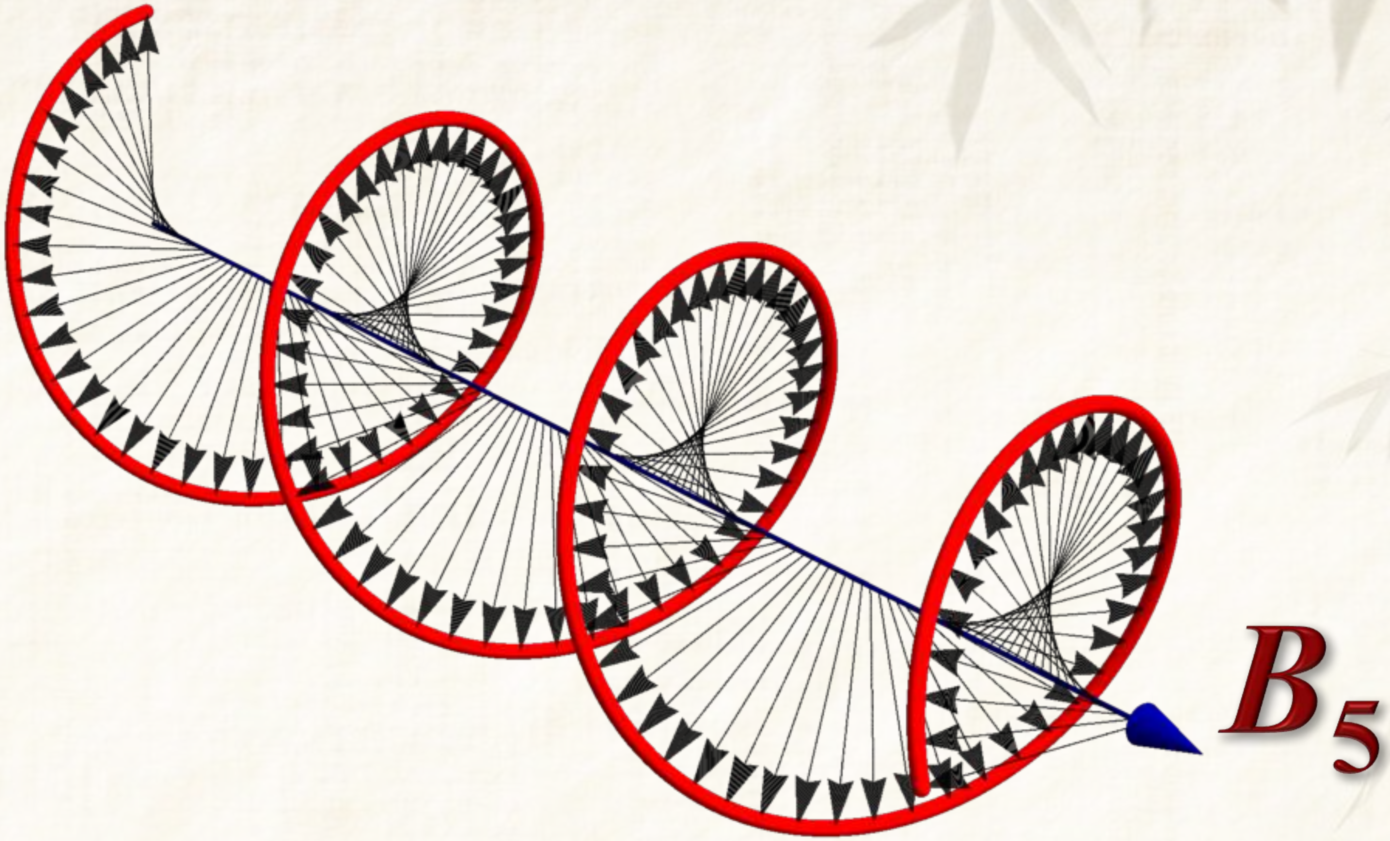


(e) $\mu = 0$, $\mu_5 = \sqrt{3\pi/(4\alpha)}\hbar\Omega_e$, $B_0 = 0$, $b_{\parallel} = 0$



(f) $\mu = 0$, $\mu_5 = \sqrt{3\pi/(4\alpha)}\hbar\Omega_e$, $B_0 = 0$, $B_{0,5} = 0$

[Gorbar, Miransky, Shovkovy, Sukhachov, Phys. Rev. B **95**, 115202 (2017)]



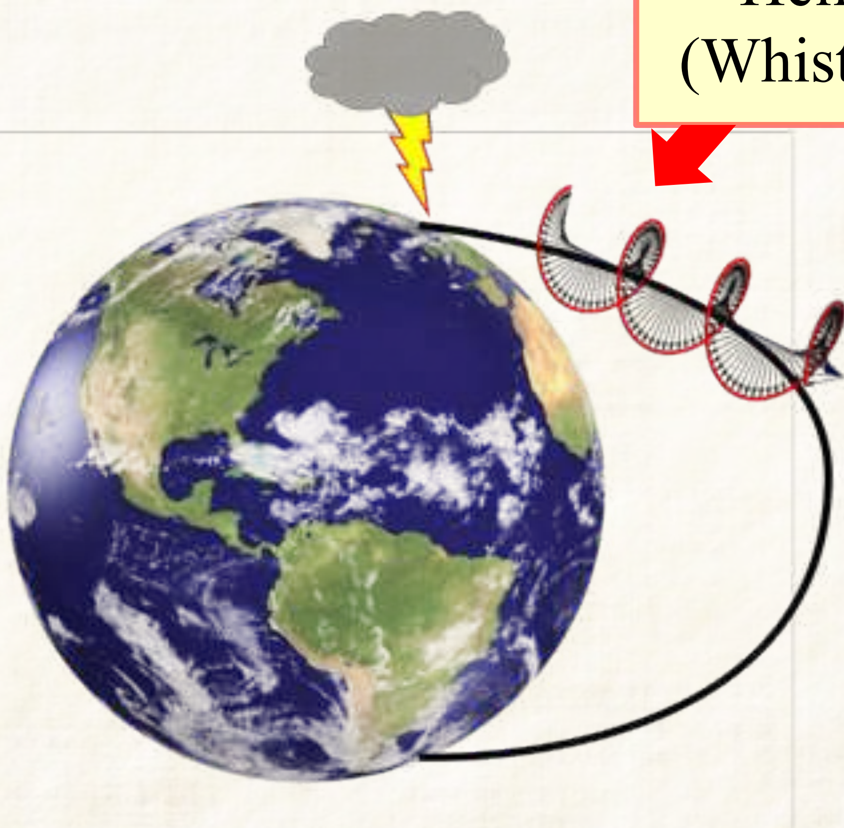
PSEUDOMAGNETIC HELICONS

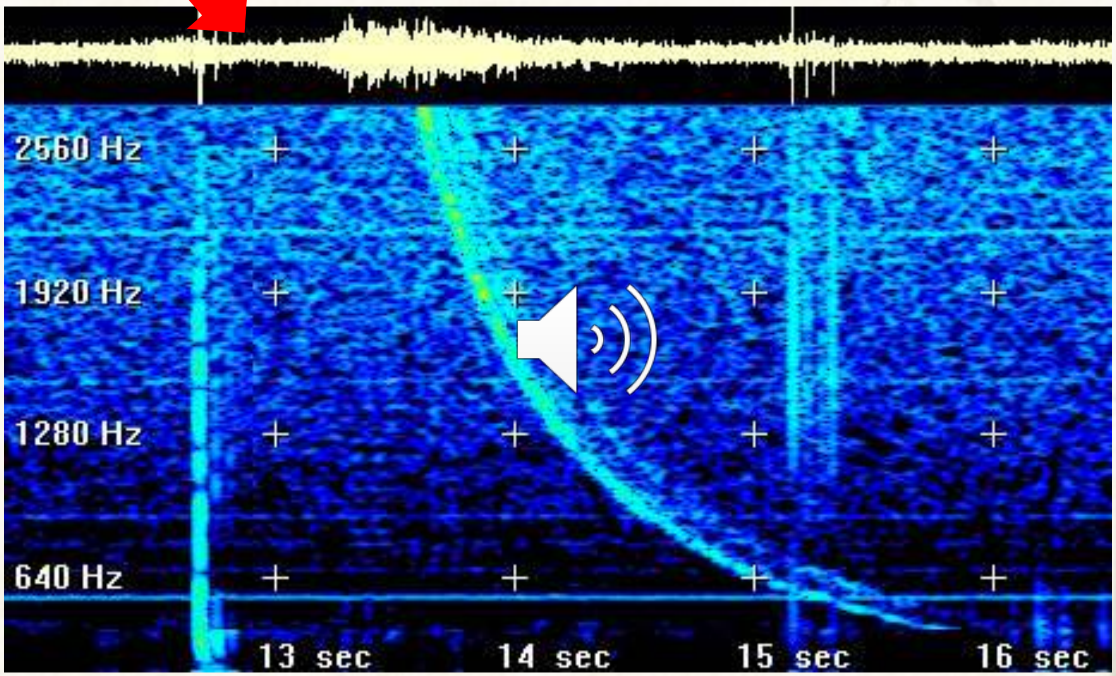
[E.V. Gorbar, V.A. Miransky, I.A. Shovkovy, and P.O. Sukhachov, Phys. Rev. B **95**, 115422 (2017)]

Pseudo-magnetic helicons

- Usual helicons are transverse low-energy gapless excitations propagating along the background magnetic field $\vec{\boldsymbol{B}}_0$ in uncompensated media (e.g., metals with different electron and hole densities)

Helicon in the ionosphere (Whistler) and its spectrogram





(Pseudo-)magnetic helicon

- Helicon dispersion law at $T \to 0$:

$$\omega_h|_{B_{0,5}\to 0,\mu_5\to 0} \overset{b_0\to 0}{=} \frac{eB_0 c^3 \hbar^2 \pi v_F^2 k^2}{\pi \hbar^2 c^2 \Omega_e^2 \mu + 2B_0 e^4 v_F^2 b_\parallel} + O(k^3)$$

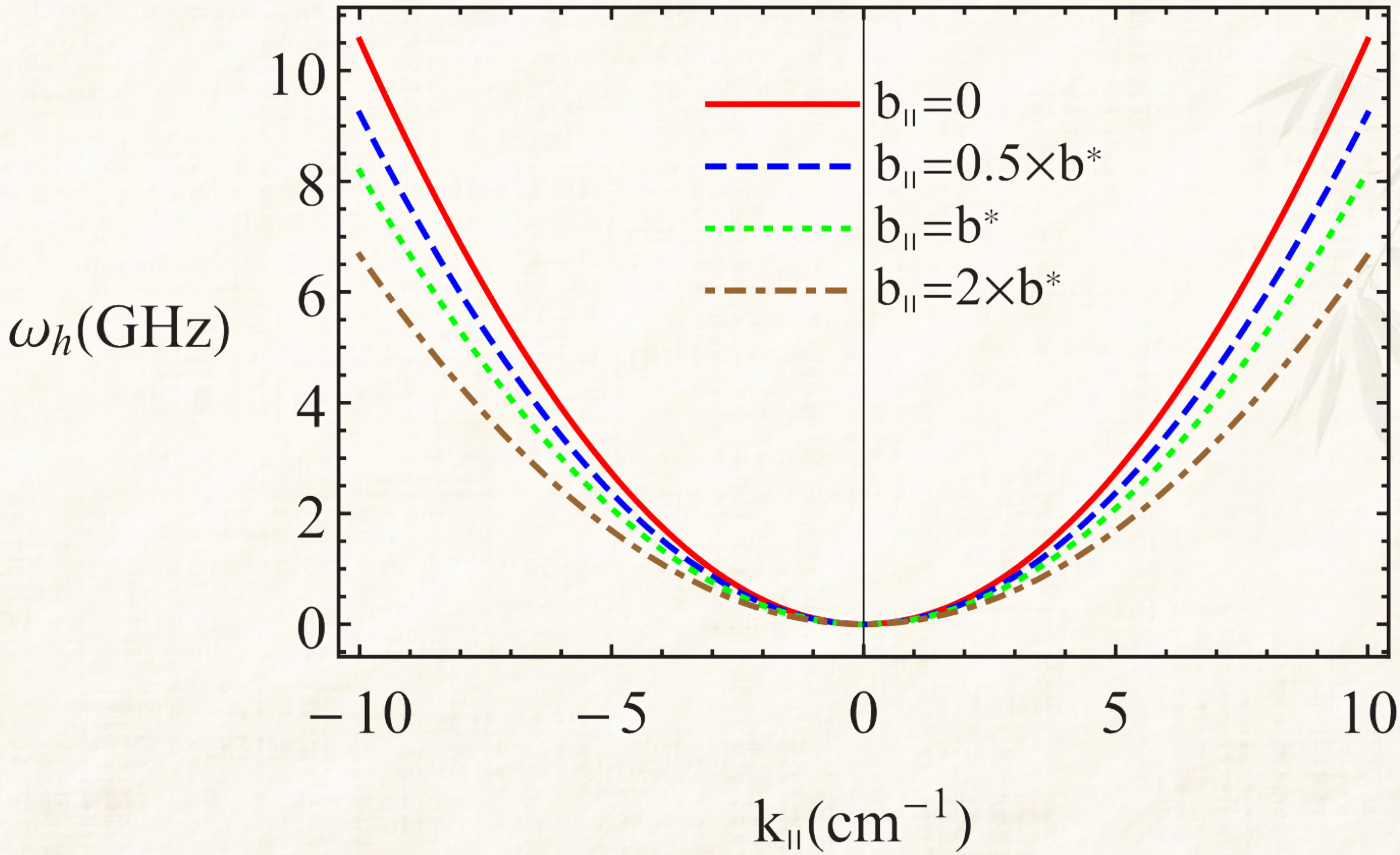
$$\omega_h|_{B_0\to 0,\mu\to 0} \overset{b_0\to -\mu_5/e}{=} \frac{eB_{0,5} c^3 \hbar^2 \pi v_F^2 k^2}{\pi \hbar^2 c^2 \Omega_e^2 \mu_5 + 2B_{0,5} e^4 v_F^2 b_\parallel} + O(k^3)$$

- Properties:
  - Gapless electromagnetic wave propagates in metals **without magnetic field**!
  - Chiral shift modifies effective helicon mass
  - In the equilibrium regime $eb_0 = -\mu_5$, the linear in the wave vector term is **absent**

[E.V. Gorbar, V.A. Miransky, I.A. Shovkovy, and P.O. Sukhachov, Phys. Rev. B **95**, 115422 (2017)]

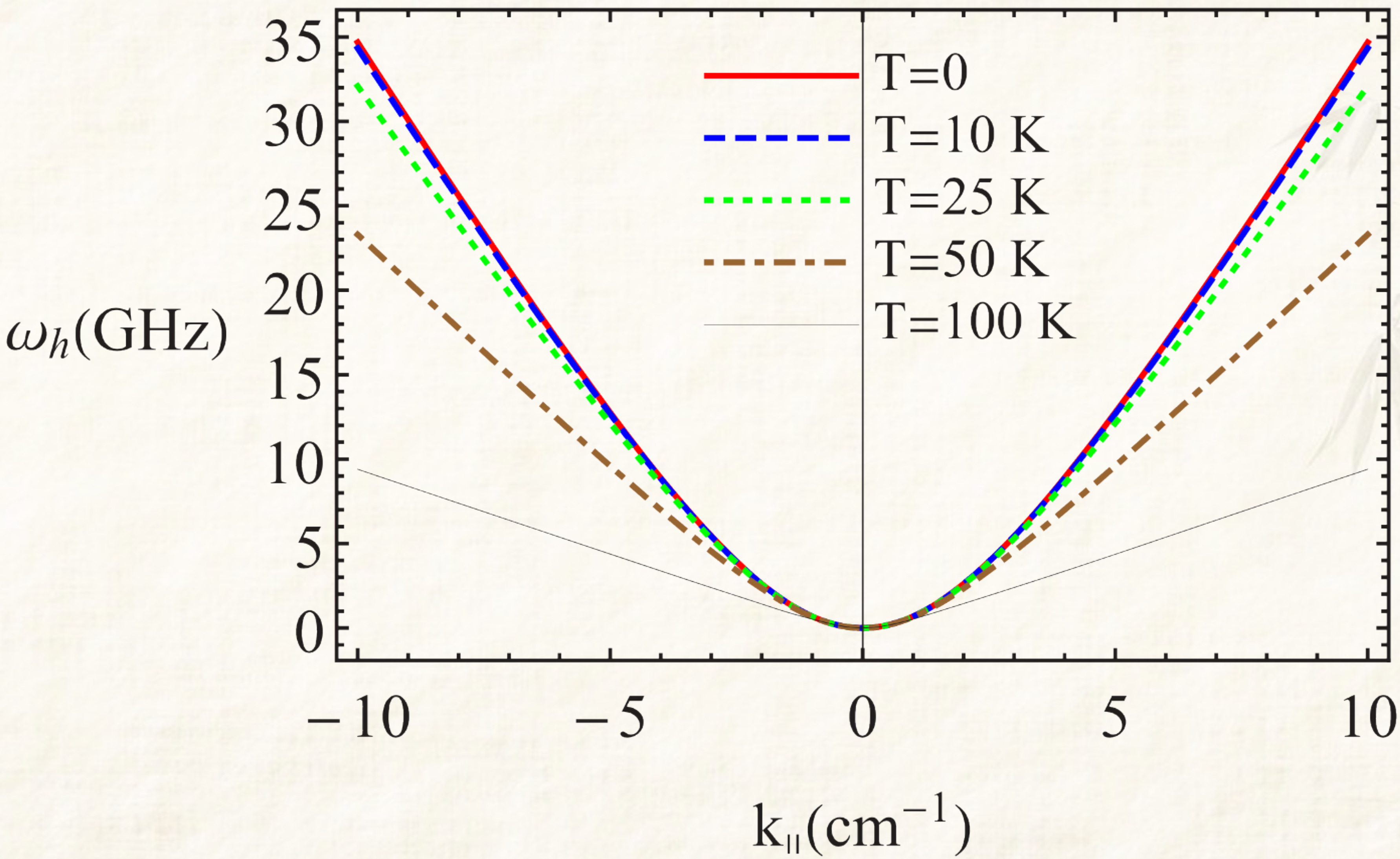
Helicons at different $b_{\parallel}$

$eb_0 = -\mu_5$, $B_{0,5} = 10^{-2}\mathrm{T}$, $\mu_5 = 5$ meV, $\mu = 0$, $eb^* = 0.3\pi\hbar v_F/c_3$



Helicons at different T

$eb_0 = -\mu_5$, $B_{0,5} = 10^{-2}\mathrm{T}$, $b_{\parallel} = 0.5b^*$, $\mu_5 = 5$ meV, $\mu = 0$



Summary

- Consistent chiral kinetic theory is needed
- Chiral magnetic plasmons (χMPs) are sensitive to local charge (non-)conservation
- Properties of χMPs carry information about b_0 and $\vec{b}$
- χMPs are not only due to the oscillation of *electric* charge, but also *chiral* charge
- New types of collective modes, pseudomagnetic helicons, may exist in Weyl materials