Photon emission from strongly magnetized QCD plasma

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[X. Wang, I. Shovkovy, L. Yu, M. Huang, Phys. Rev. D 102, 076010 (2020)]

[X. Wang, I. Shovkovy, arXiv:2103.01967]

APS April Meeting
Apr 20, 2021
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APS physics APRIL MEETING 2021

quarks Q2C cosmos

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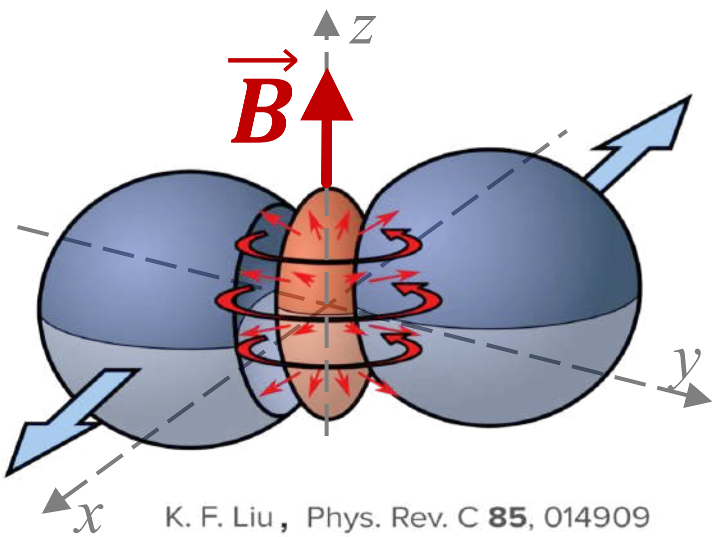
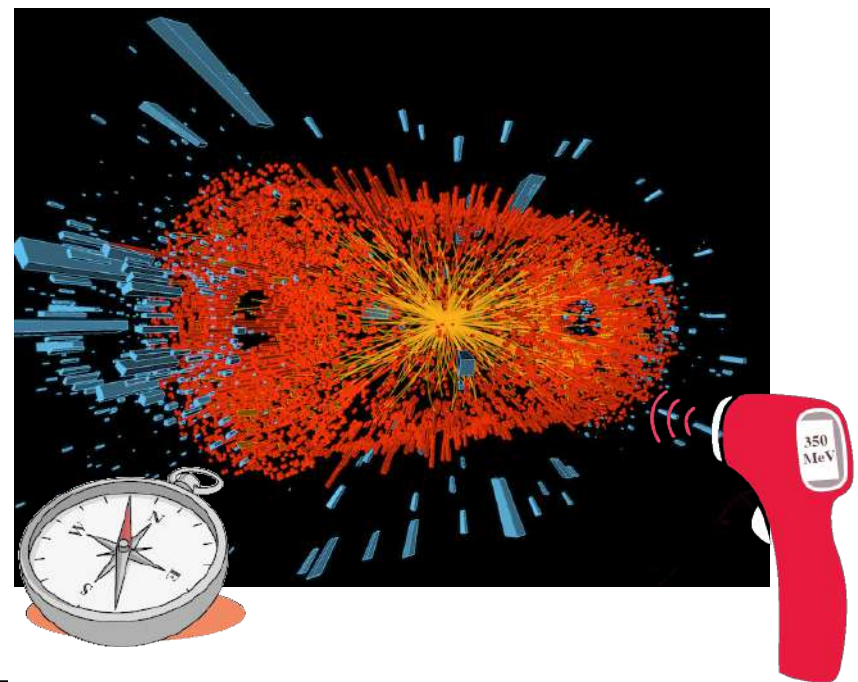
Magnetic field in HICs

- QGP produced at RHIC/LHC is **magnetized**
  - **10^{18} to 10^{19} G** $\sim m_\pi^2 \sim (100\ \mathrm{MeV})^2$



K. F. Liu, Phys. Rev. C **85**, 014909



Main idea

- Photon emission is not only a **thermo**meter but also **magneto**meter of QGP

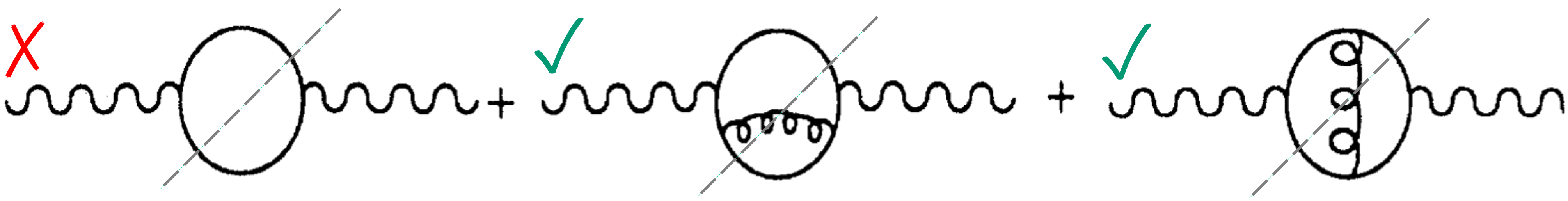
Thermal photons ($B = 0$)

- The rate of the thermal emission of photons (more precisely, the energy loss rate) is

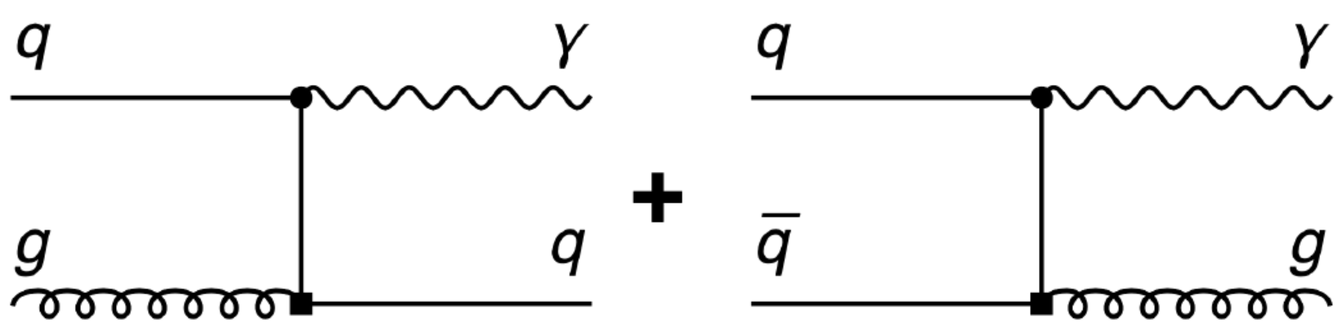
$$k^0 \frac{d^3 R}{dk_x dk_y dk_z} = -\frac{1}{(2\pi)^3} \frac{\mathrm{Im}\left[\Pi^\mu_\mu(k)\right]}{\exp\left(\frac{k_0}{T}\right) - 1}$$

[Kapusta, Lichard, Seibert, Phys. Rev. D 44, 2774 (1991)]
[Baier, Nakkagawa, Niegawa, Redlich, Z. Physik C 53 (1992) 433]

- In the case of hot QCD plasma,



- Processes:



Thermal photons ($B = 0$)

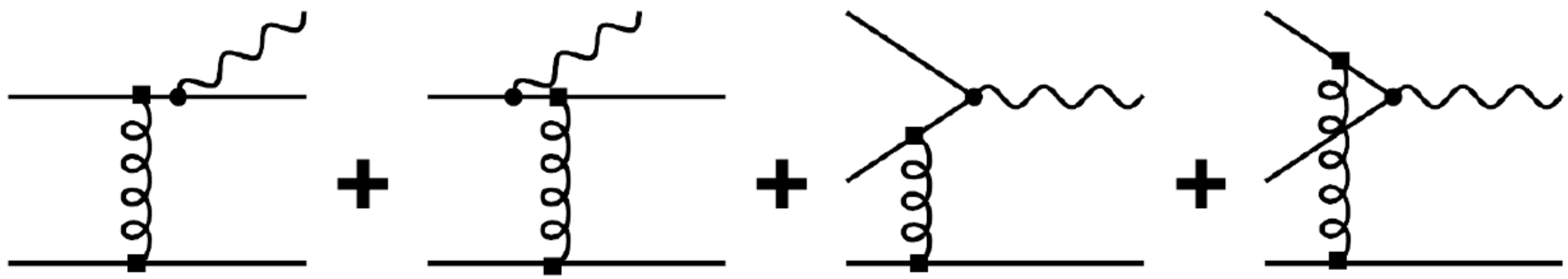
- The approximate result is given by

$\mathcal{R}_{2\to 2}$:
$$E\frac{dR}{d^3p} = \frac{5}{9}\frac{\alpha\alpha_s}{2\pi^2}T^2 e^{-E/T}\ln\left(\frac{2.912}{g^2}\frac{E}{T}\right)$$

[Kapusta, Lichard, Seibert, Phys. Rev. D 44, 2774 (1991)]

- There are important corrections from **bremsstrahlung** and **inelastic pair annihilation**

$\mathcal{R}_{2\to 3}$, $\mathcal{R}_{3\to 2}$:



[Arnold, Moore,Yaffe, JHEP 12 (2001) 009; hep-ph/0111107]

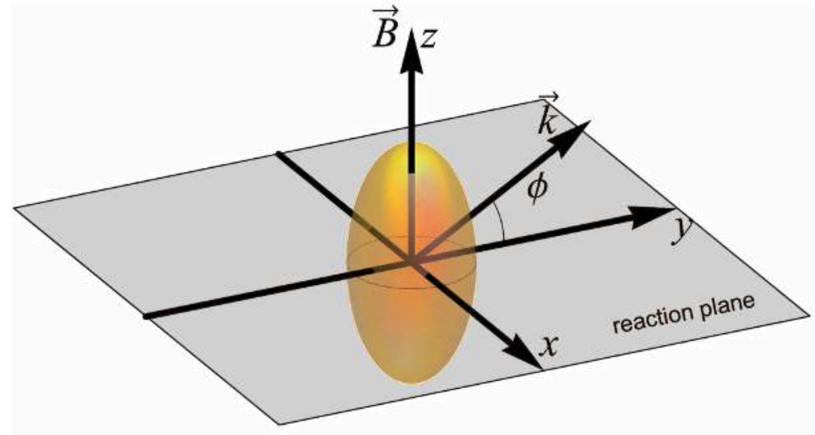
- Next to leading order corrections are ~ 100%

[Arnold, Moore,Yaffe, JHEP 12 (2001) 009; hep-ph/0111107]
[Ghiglieri et al., JHEP 05 (2013) 010; arXiv:1302.5970]

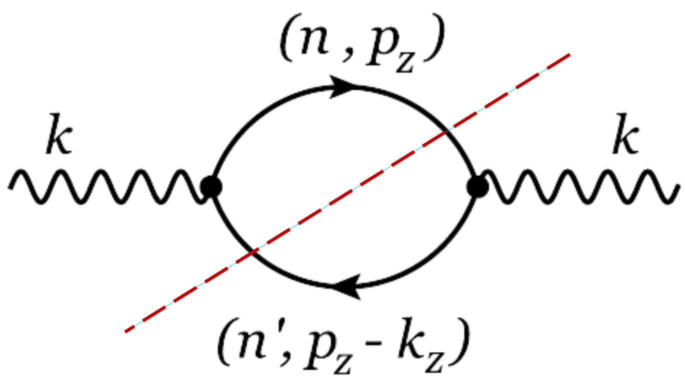
Photon emission rate

- The expression for the rate is

$$k^0 \frac{d^3 R}{dk_x dk_y dk_z} = -\frac{1}{(2\pi)^3} \frac{\text{Im}\left[\Pi^{\mu}_{\mu}(k)\right]}{\exp\left(\frac{k_0}{T}\right) - 1}$$



- At $\vec{B} \neq 0$, the imaginary part of the polarization tensor

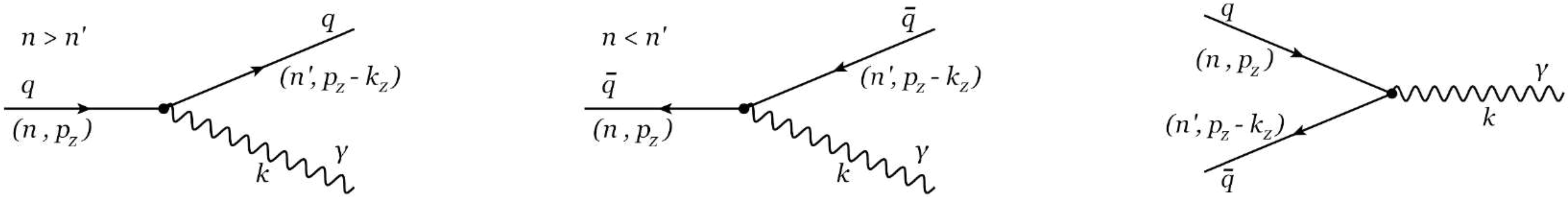


is nonzero at leading order in α_s!

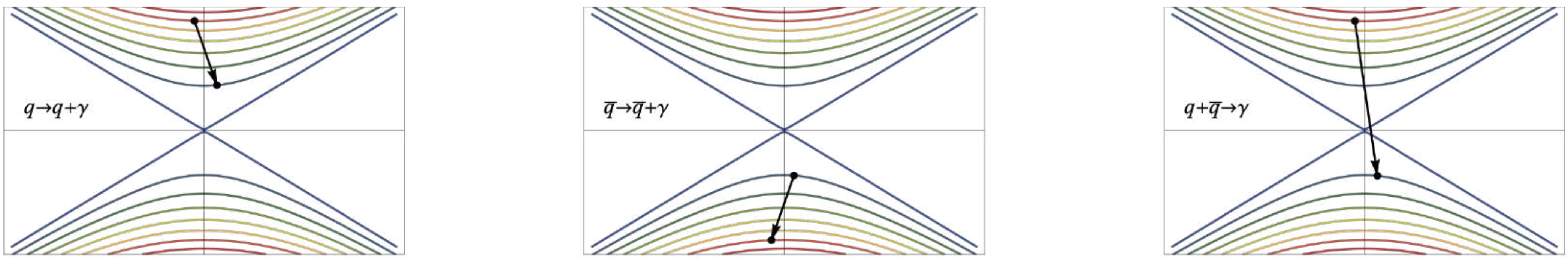
[Wang, Shovkovy, Yu, Huang, Phys. Rev. D 102, 076010 (2020), arXiv:2006.16254]

Physics processes

- Relevant physics processes (0^{th} order in α_s):



(a) (b) (c)



The energy momentum conservation

$$E_{n,p_z,f} - \lambda E_{n',p_z-k_z,f} + \eta\Omega = 0$$

is satisfied for these 1 → 2 and 2 → 1 processes

[Wang, Shovkovy, Yu, Huang, Phys. Rev. D 102, 076010 (2020), arXiv:2006.16254]

Photon emission rate

The explicit expression for $\mathrm{Im}\left[\Pi^{\mu}_{R,\mu}(\Omega,\mathbf{k})\right]$ is

$$\mathrm{Im}\left[\Pi^{\mu}_{R,\mu}\right] = \sum_{f=u,d}\frac{N_c\alpha_f}{2\pi l_f^4}\sum_{n>n'}^{\infty}\frac{g(n,n')\left[\theta\left(k_-^f-|k_y|\right)-\theta\left(|k_y|-k_+^f\right)\right]}{\sqrt{[(k_-^f)^2-k_y^2][(k_+^f)^2-k_y^2]}}\left(\mathcal{F}_1^f+\mathcal{F}_4^f\right)$$

$$-\sum_{f=u,d}\frac{N_c\alpha_f}{4\pi l_f^4}\sum_{n=0}^{\infty}\frac{g_0(n)\theta\left(|k_y|-k_+^f\right)}{\sqrt{k_y^2[k_y^2-(k_+^f)^2]}}\left(\mathcal{F}_1^f+\mathcal{F}_4^f\right),$$

where $g(n,n')$ and $g_0(n)$ are combinations of the Fermi-Dirac distribution functions.

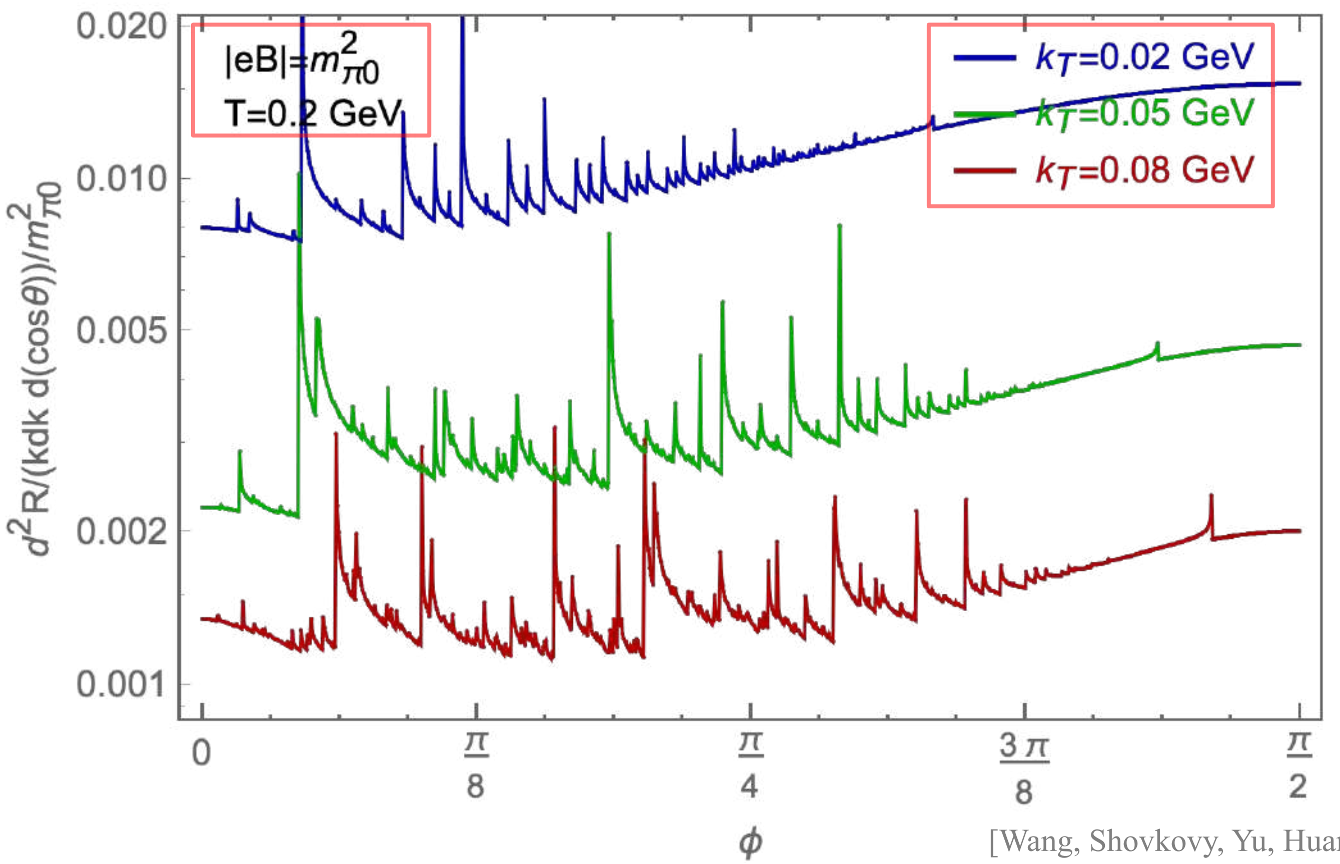
The momentum *thresholds* are determined by

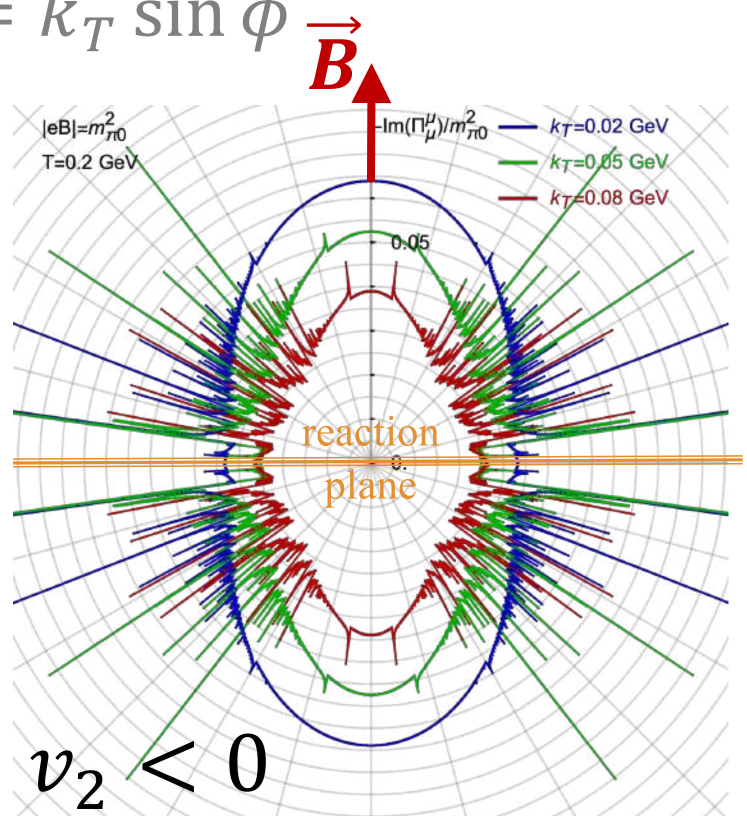
$$k_\pm^f = \left|\sqrt{m^2+2n|e_fB|}\pm\sqrt{m^2+2n'|e_fB|}\right|$$

[Wang, Shovkovy, Yu, Huang, Phys. Rev. D 102, 076010 (2020), arXiv:2006.16254]

Angular dependence (1)

- At small k_T, the emission rate is maximal at $\phi = \frac{\pi}{2}$ (i.e., **perpendicular** to the reaction plane)
- Effectively, this gives photon "flow" with $v_2 < 0$
- Note: $k_x = 0$, $k_y = k_T \cos\phi$ and $k_z = k_T \sin\phi$

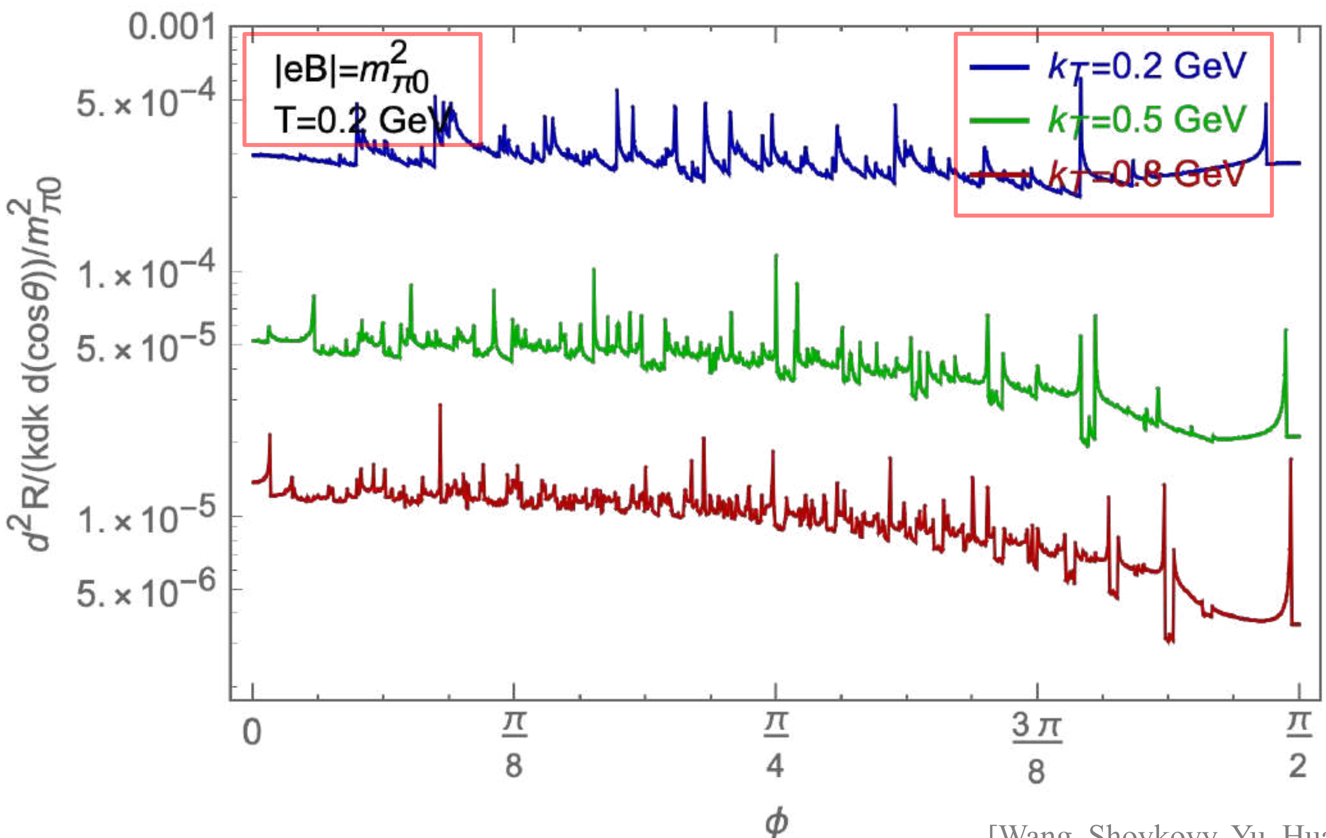


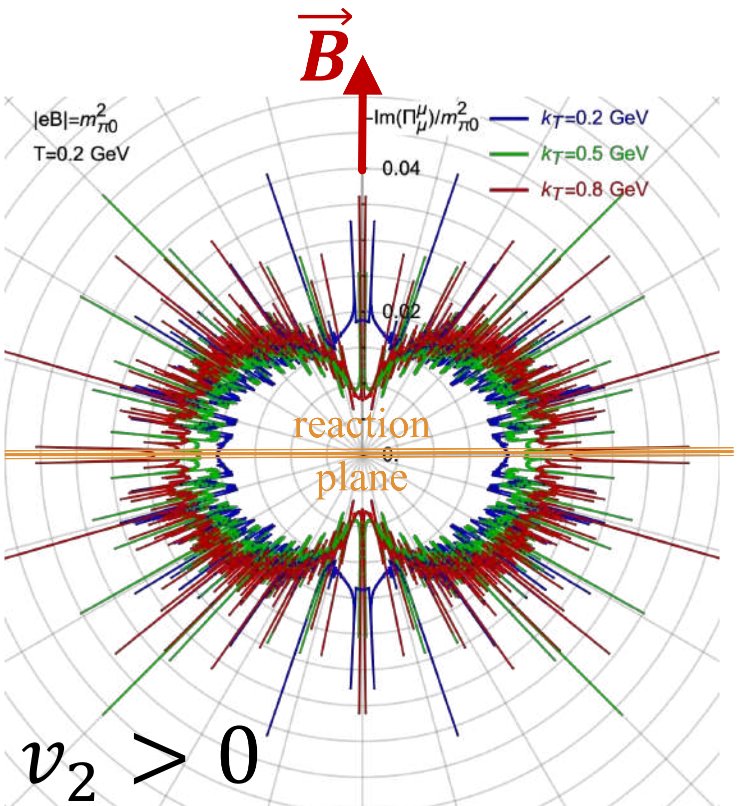


[Wang, Shovkovy, Yu, Huang, Phys. Rev. D 102, 076010 (2020), arXiv:2006.16254]

Angular dependence (2)

- At large k_T, the emission rate is maximal at $\phi = 0$ (i.e., **parallel** to the reaction plane)
- Effectively, this gives photon "flow" with $v_2 > 0$

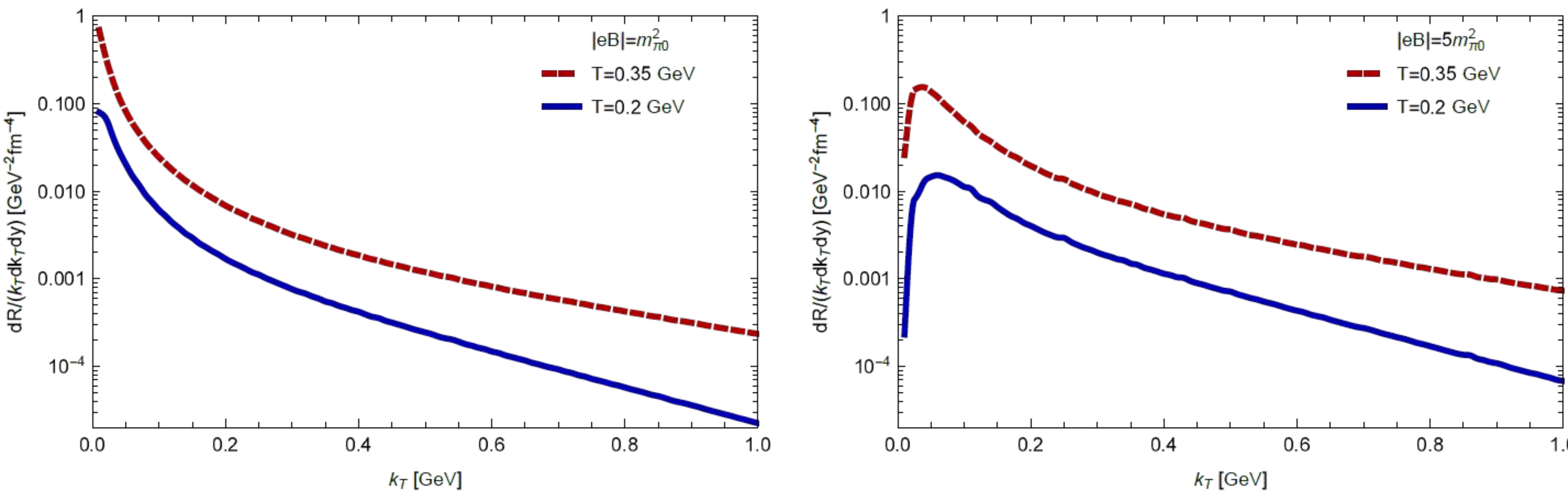




[Wang, Shovkovy, Yu, Huang, Phys. Rev. D 102, 076010 (2020), arXiv:2006.16254]

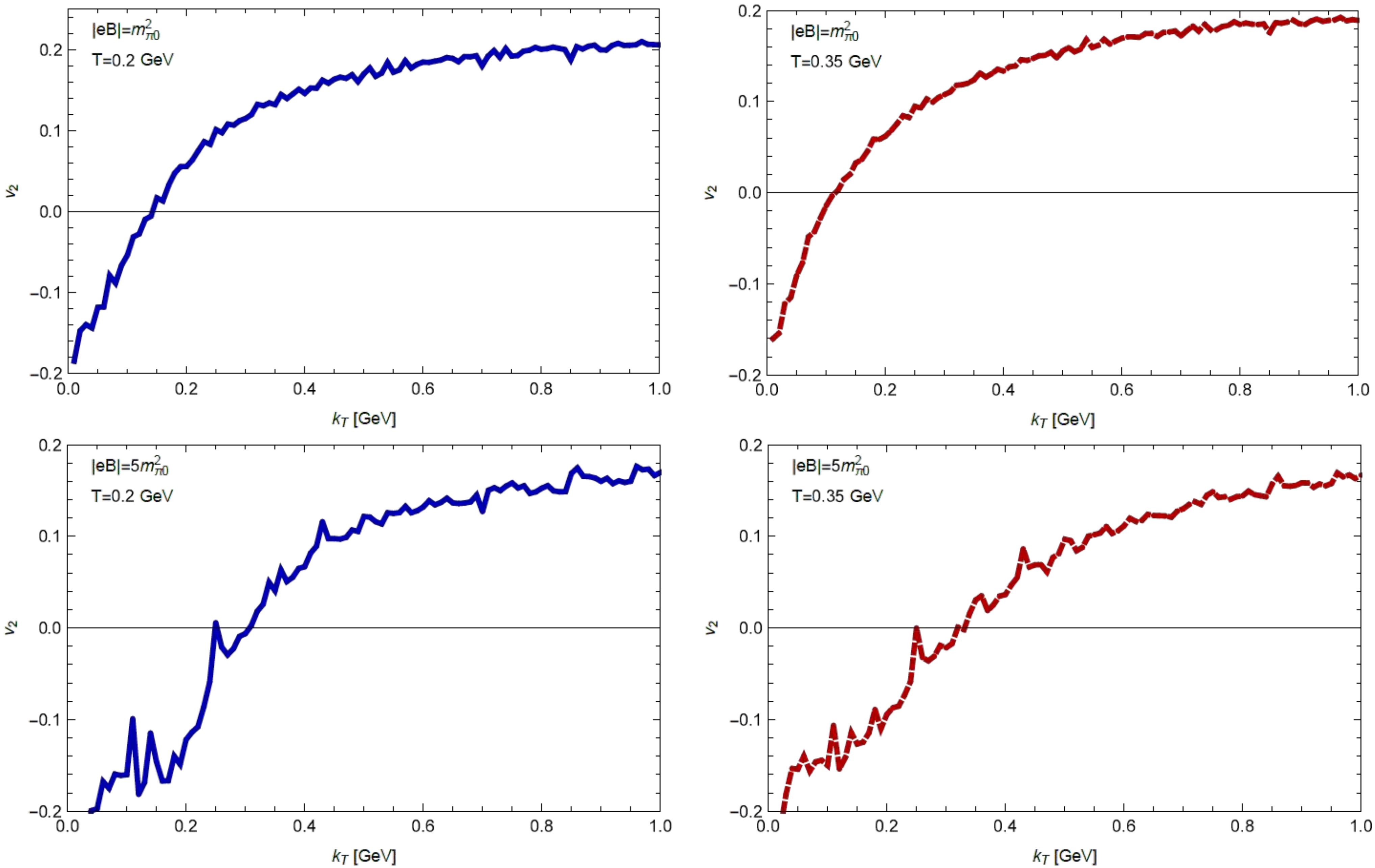
Thermal rate at $\vec{B} \neq 0$

- The photon production rate
  - decreases with energy (k_T) at large k_T
  - increases with temperature
  - goes to zero when $k_T \to 0$ (quantization effects)
  - has a peak at a small nonzero k_T



[Wang, Shovkovy, Yu, Huang, Phys. Rev. D 102, 076010 (2020), arXiv:2006.16254]

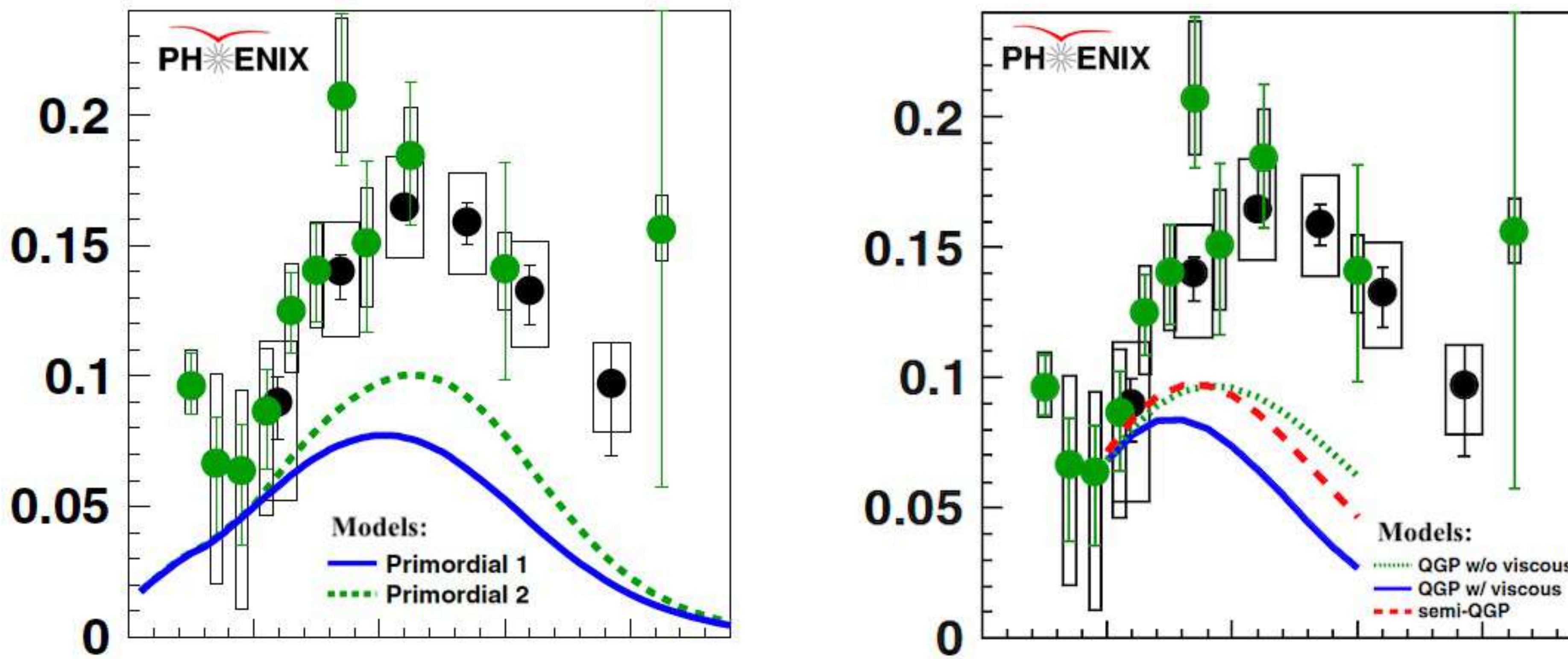
Nonzero elliptic "flow" (v_2)



[Wang, Shovkovy, Yu, Huang, Phys. Rev. D 102, 076010 (2020), arXiv:2006.16254]

Photon v_2 puzzle

- Most photons are produced early (before flow develops)
- Thus, v_2 for photons should be very small



[Adare et al (PHENIX Collaboration), Phys. Rev. C **94**, 064901 (2016)]

Magnetic enhancement of v_2

- Estimate of v_2 in a hot magnetized QGP

$$\mathcal{R}_{\substack{2\to 1\\ 1\to 2}}: \quad v_2 \sim 20\%$$

- Noting that

$$\mathcal{R}_{\substack{2\to 1\\ 1\to 2}} \sim \mathcal{R}_{2\to 2} \sim \mathcal{R}_{\substack{2\to 3\\ 3\to 2}}$$

- Naïve estimate at $p_T \sim 1$ GeV gives

$$6.7\% \lesssim v_2 \lesssim 20\%$$

- A more realistic estimate should consider non-isotropic expansion & non-thermal processes

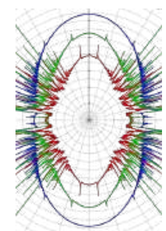
Summary

- $\vec{B} \neq 0$: photons are produced at 0^{th} order in α_s

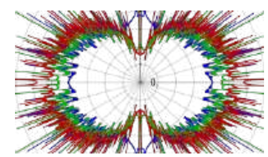
  (i) $q \to q + \gamma$, (ii) $\bar{q} \to \bar{q} + \gamma$, (iii) $q + \bar{q} \to \gamma$

- Photon emission has pronounced ellipticity

  – $v_2 < 0$ at small k_T ($k_T \lesssim \sqrt{|eB|}$)



  – $v_2 > 0$ at large k_T ($k_T \gtrsim \sqrt{|eB|}$)



- A nonzero ellipticity of thermal emission is a "measure" of the magnetic field in collisions