

ASU Theoretical Physics Colloquium
Oct 19 2022

Probing physics beyond the Standard Model at low energies

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University of Washington

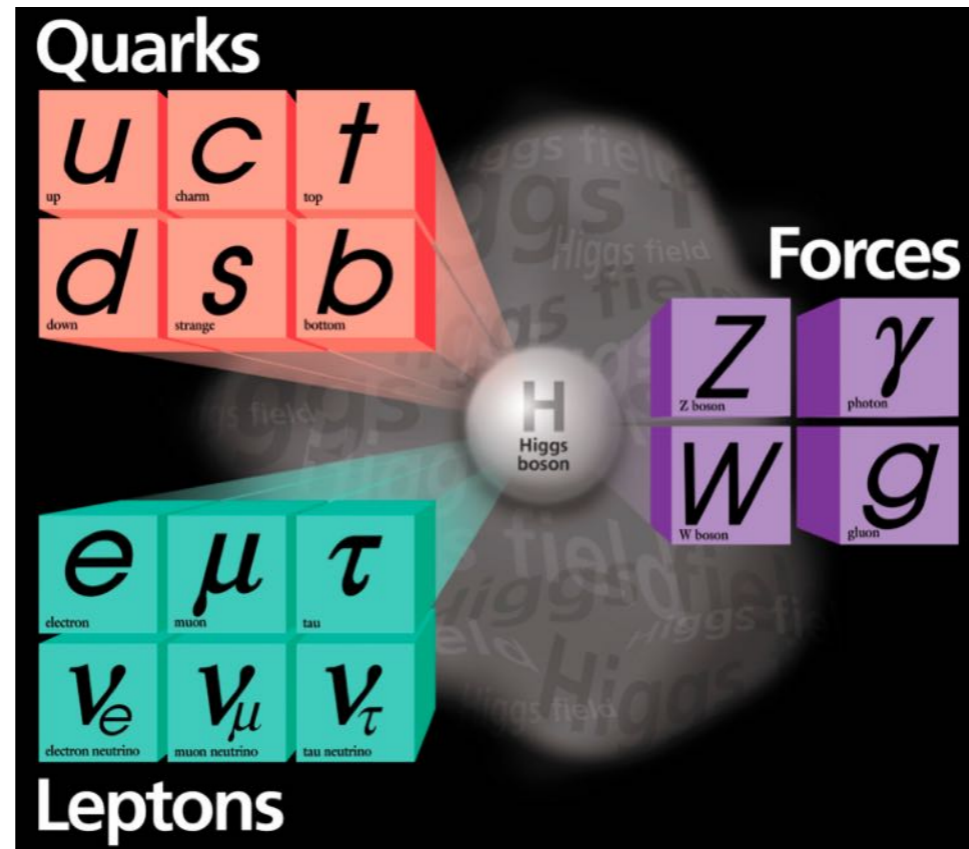


Outline

- The quest for new physics: Energy and Precision Frontiers
- Two exciting Precision Frontier probes
 - β decays as a probe of new physics at the multi-TeV scale
 - Neutrinoless $\beta\beta$ decay and Lepton Number Violation

The quest for new physics

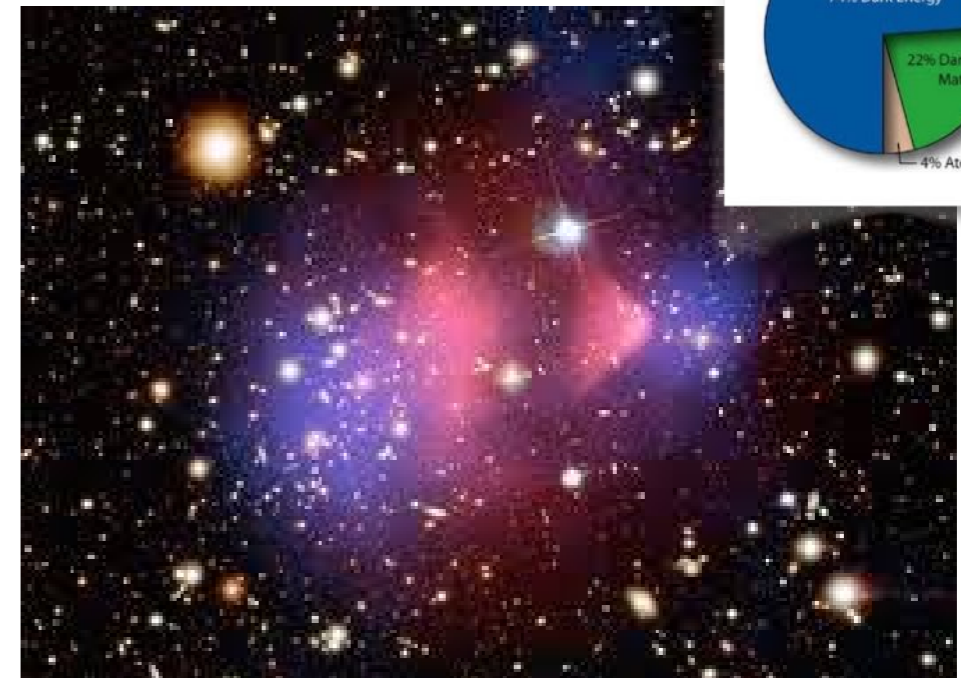
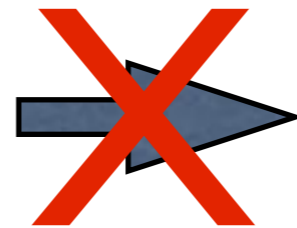
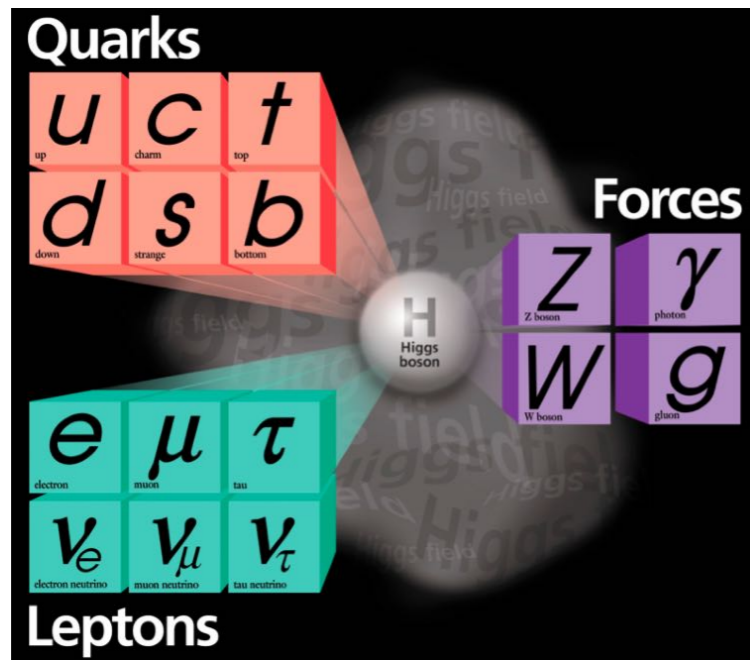
The (known) building blocks of nature



The “Standard Model” is a remarkably successful theory, tested over a wide range of energies (atomic to \sim TeV)

However, the SM is probably not the whole story...

New physics: why?

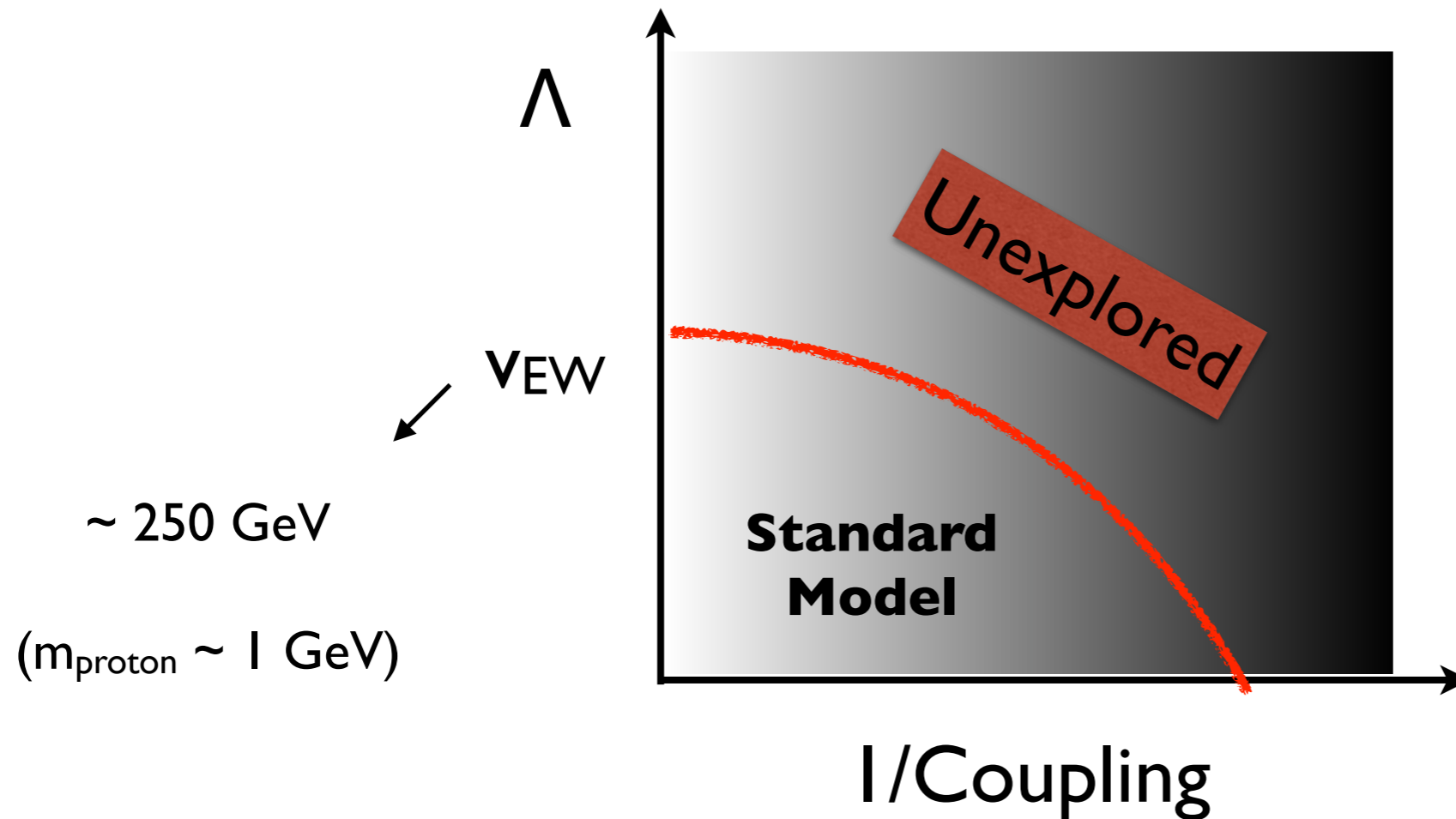


No Baryon Asymmetry, no Dark Matter, no Dark Energy, no Neutrino Mass
Origin of families, Strong CP problem, Higgs naturalness, Unification,...

Addressing these puzzles requires new physics

New physics: where?

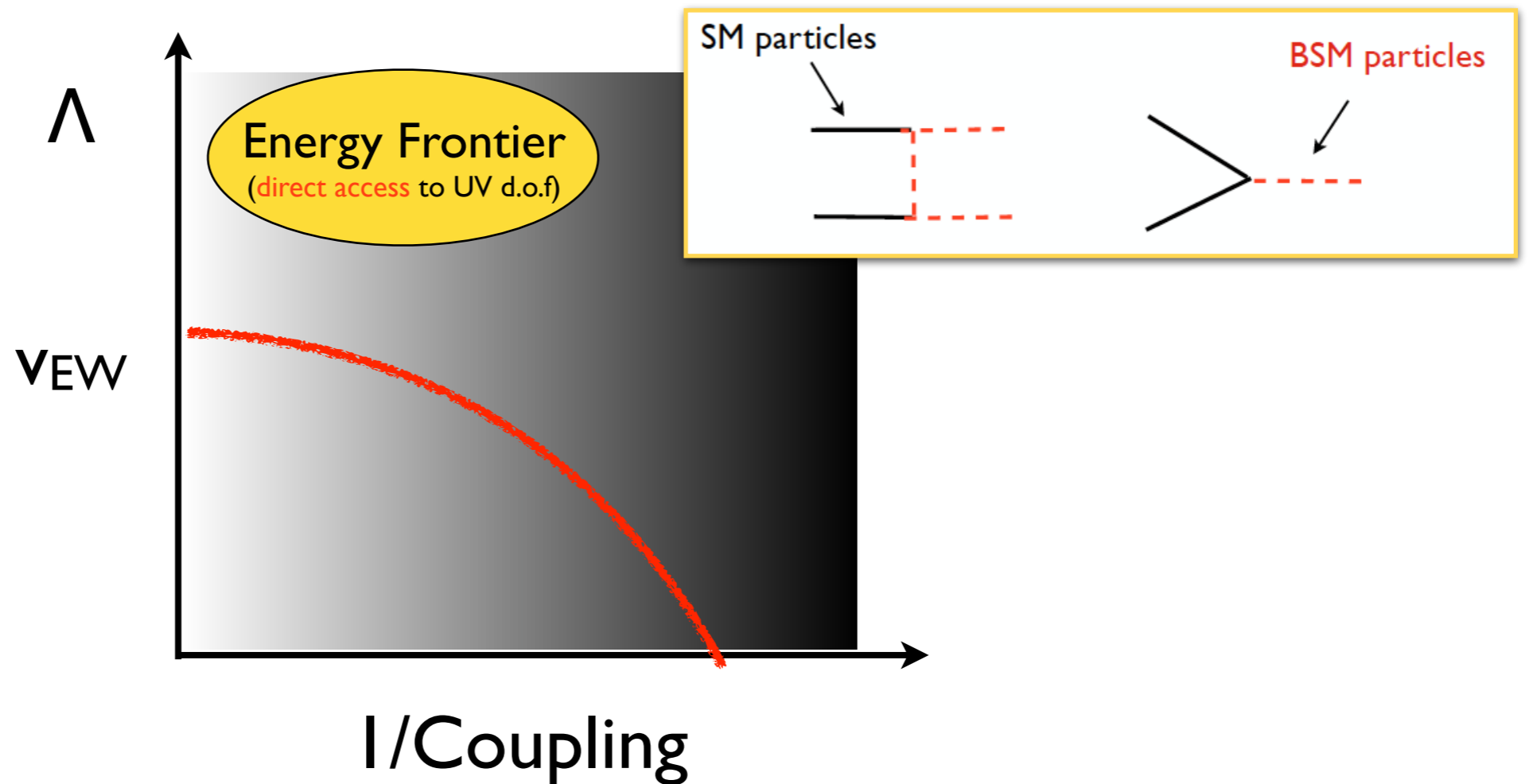
- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?



New physics: how?

- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?

$$E = mc^2$$

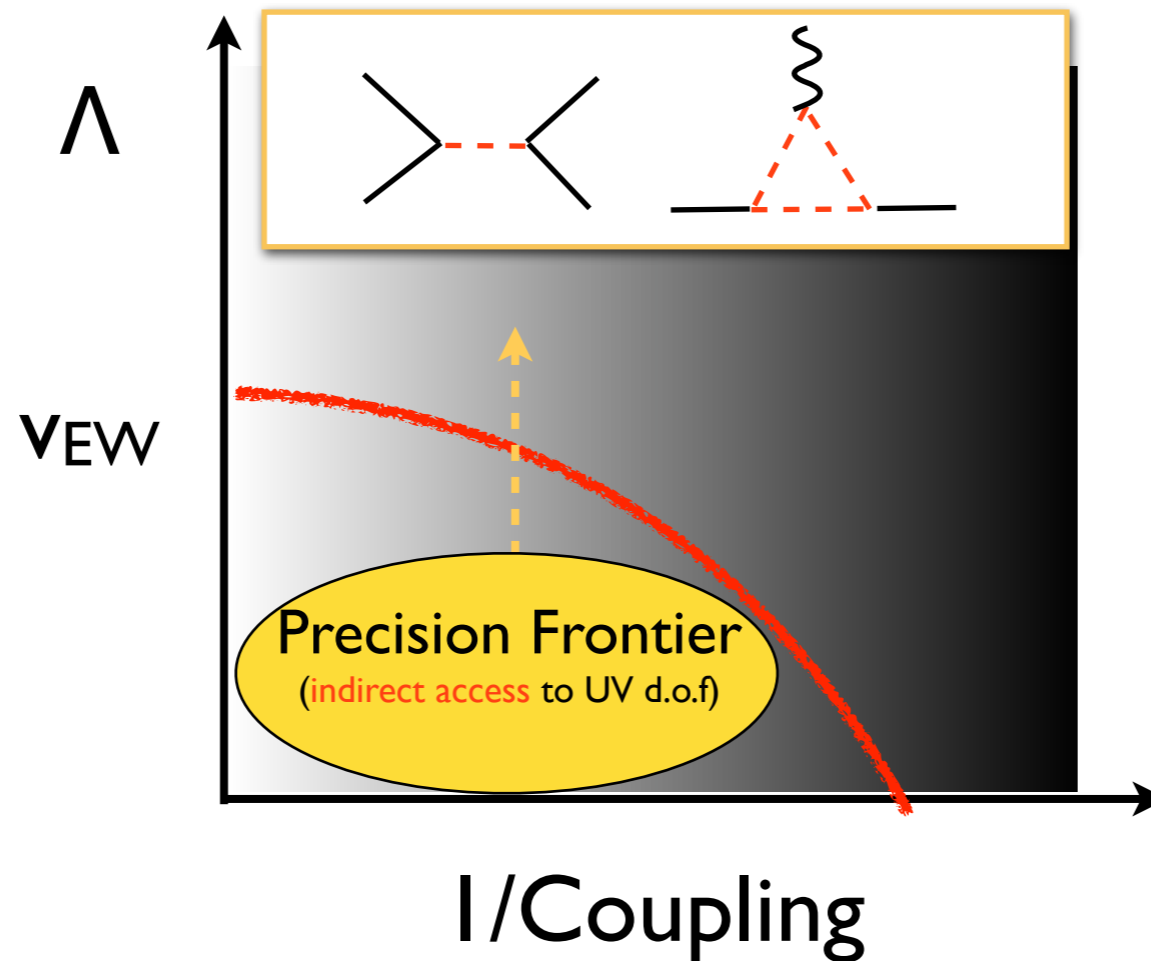


- Two approaches

New physics: how?

- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?

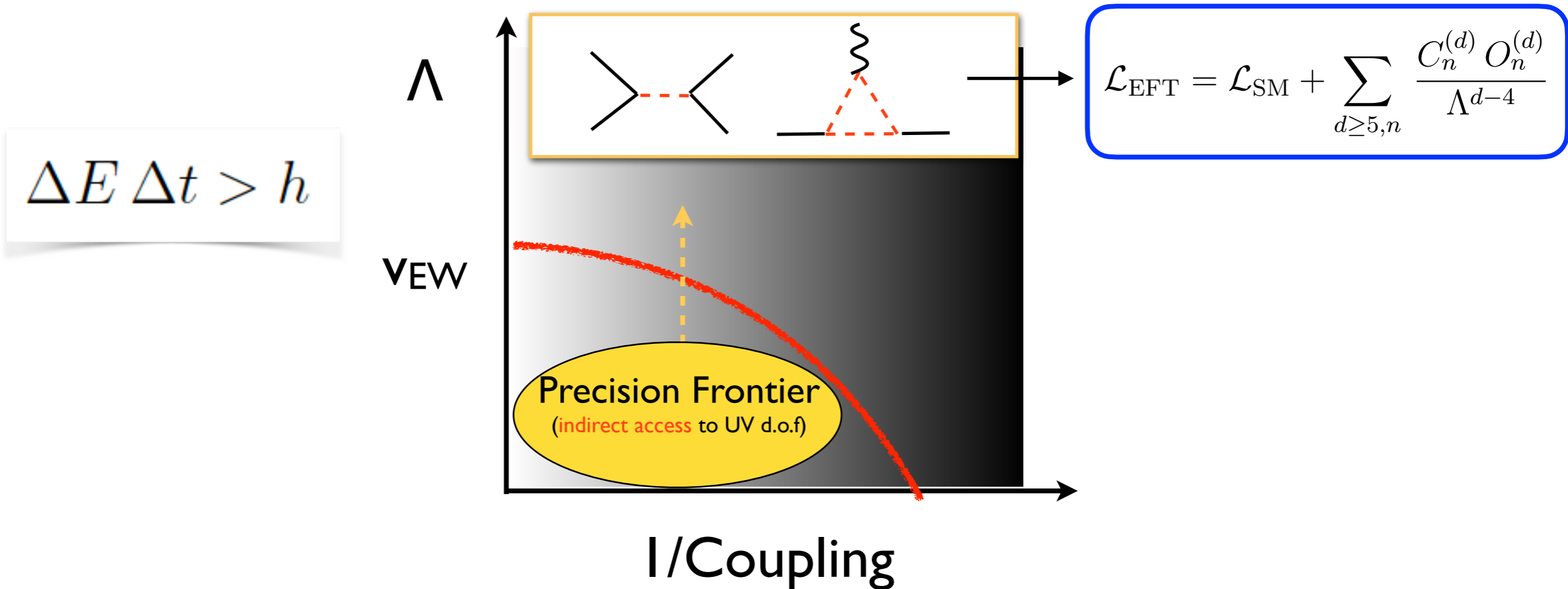
$$\Delta E \Delta t > h$$



- Two approaches

New physics: how?

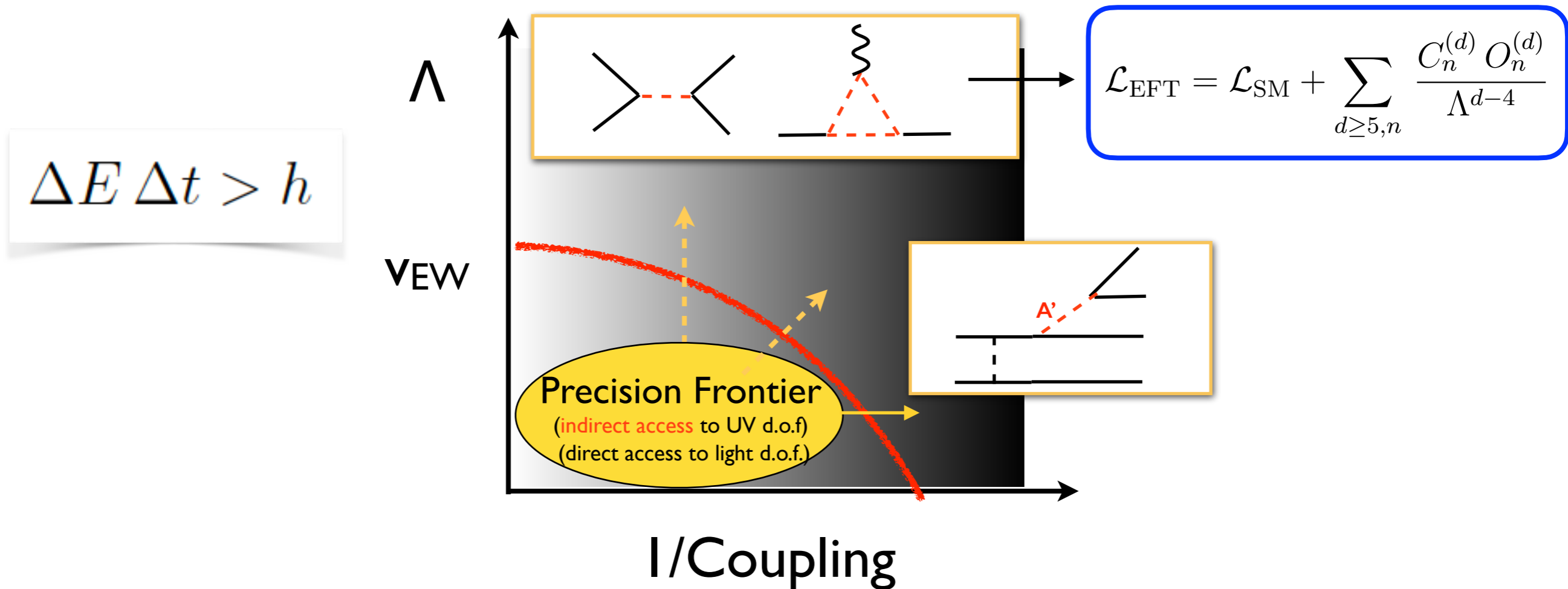
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- Two approaches

New physics: how?

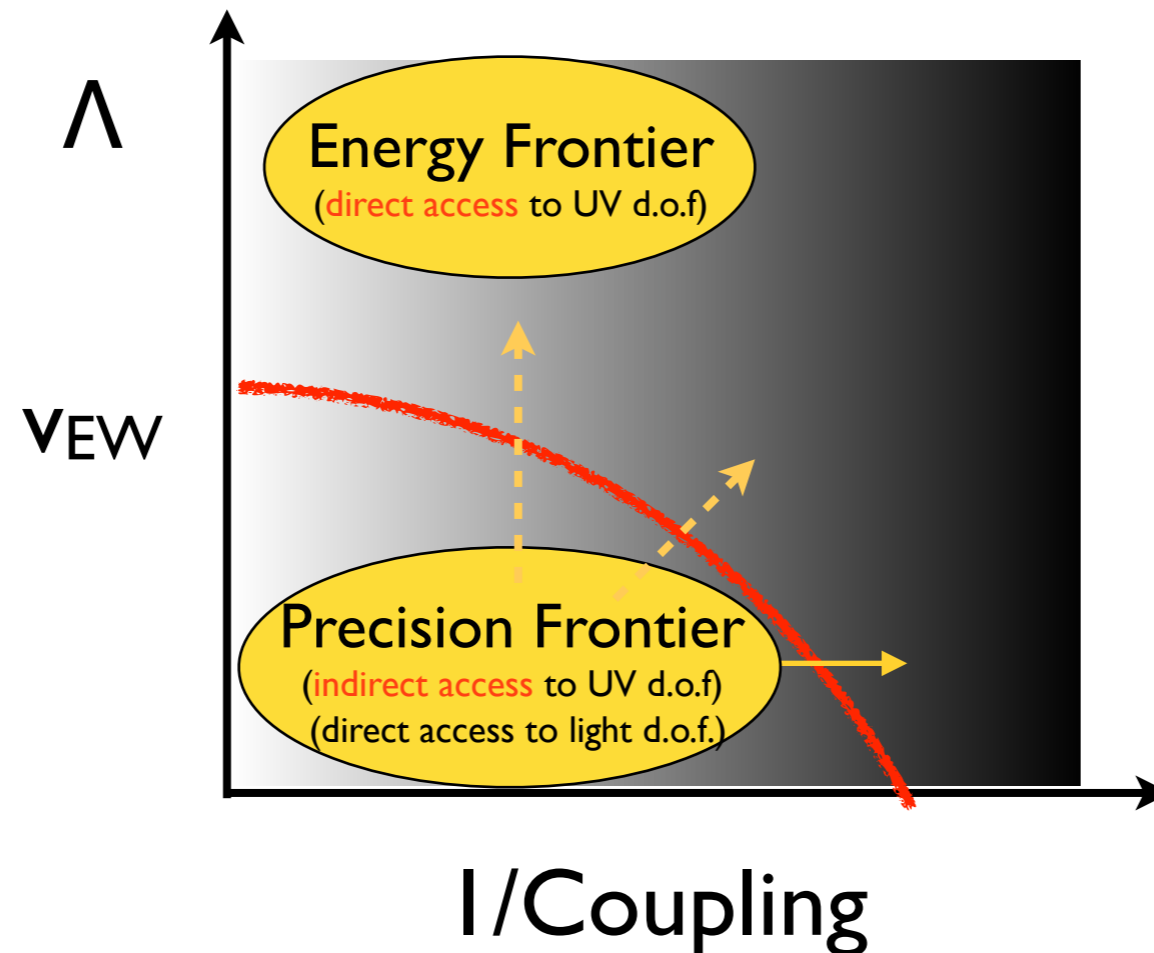
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- Two approaches

New physics: how?

- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?

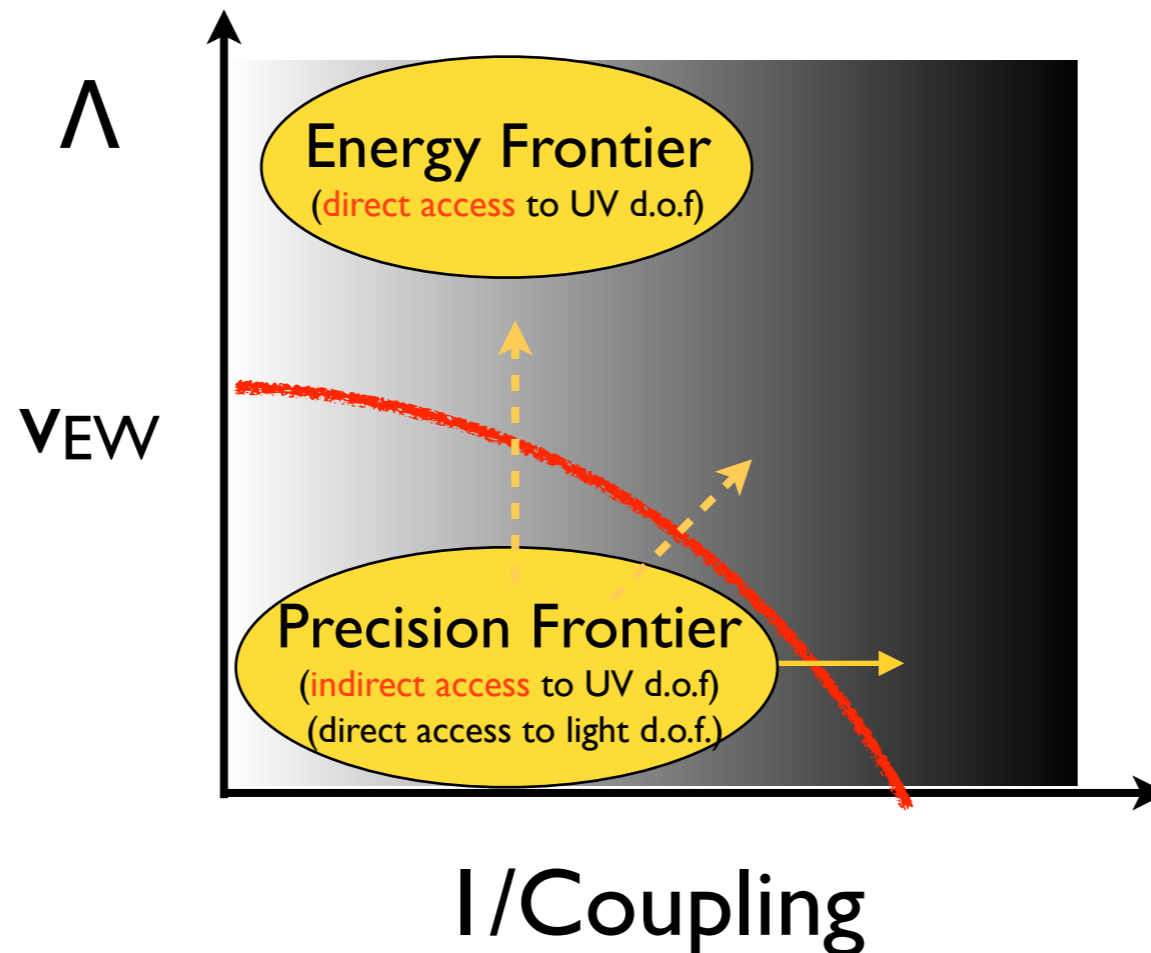


- EWSB mechanism
- Direct access to heavy particles
- ...
- L and B non conservation
- CP violation (w/o flavor)
- Flavor violation: quarks, leptons
- Precision tests sensitive to multi-TeV scale force mediators
- Light and weakly coupled particles: neutrino properties, dark sectors, ...

- Two approaches, both needed to pin down BSM dynamics (content and symmetries of \mathcal{L}_{BSM}) and answer open questions

New physics: how?

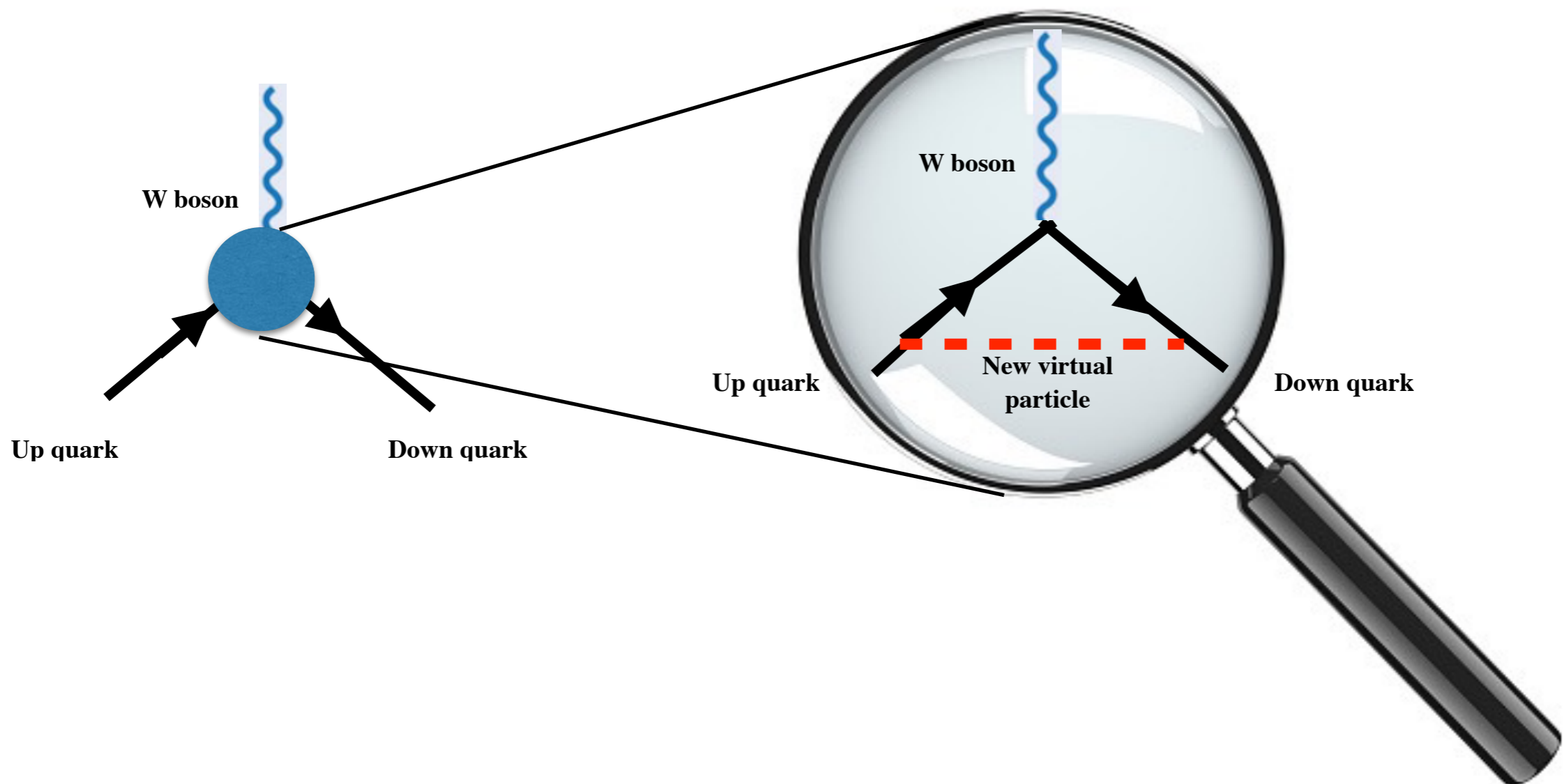
- Where is the new physics? Is it Heavy? Is it Light & weakly coupled?



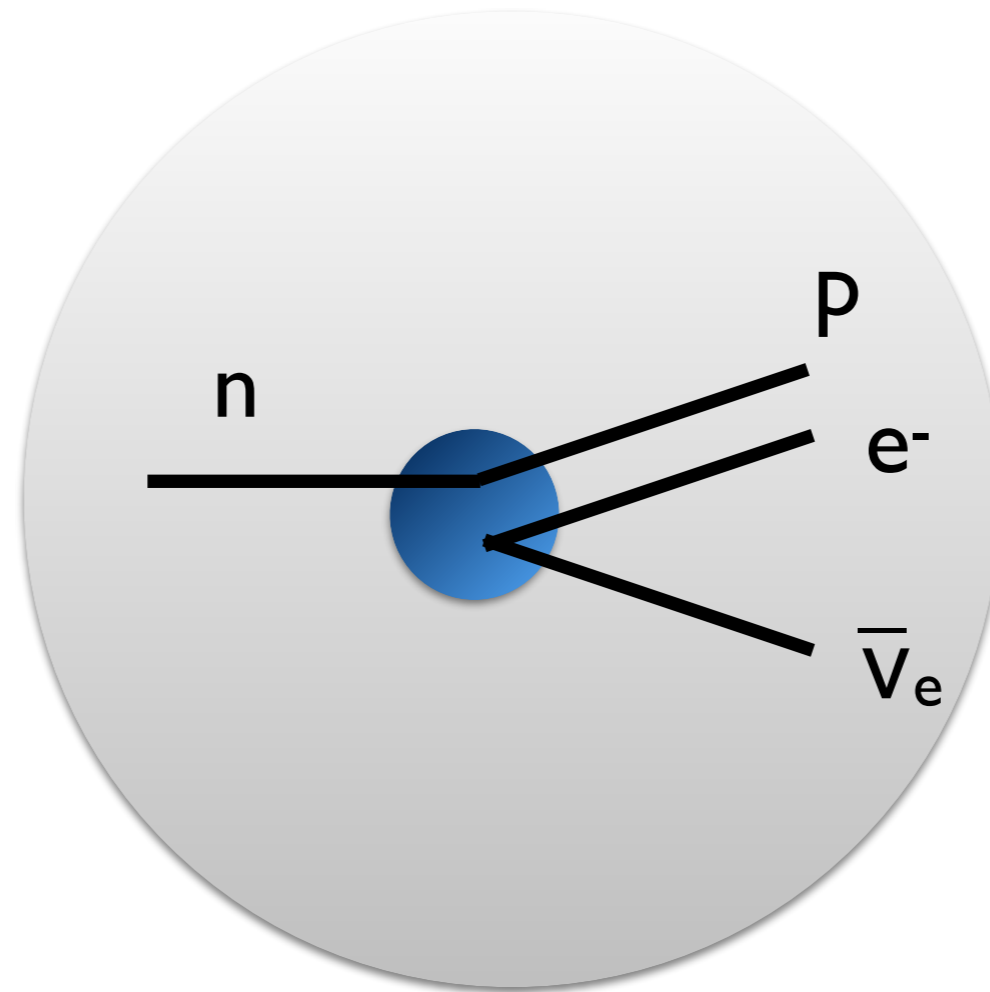
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I will discuss β and $\beta\beta$ decays as probes of BSM weak interactions and $L\neq$ non-conservation, respectively

Beta decays as a probe of new physics



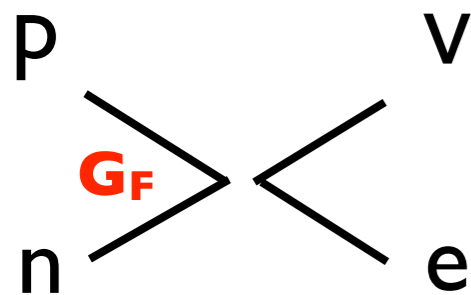
β -decays in the SM and beyond



- Beta decays have played a central role in the development of the SM
- Nowadays: tool to challenge the SM & probe possible new physics

β -decays: historic perspective

Fermi, 1934



Current-current,
parity conserving

Fermi scale:
 $\Lambda = G_F^{-1/2} \sim 250 \text{ GeV}$

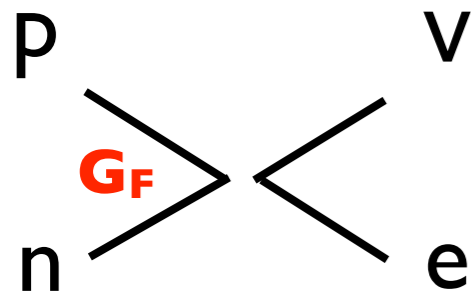
Fermi's theory of beta decays ($n \rightarrow p e \bar{\nu}_e$):

Postulate new local interaction in terms of “light”
degrees of freedom (n, p, e, ν_e): $H \sim G_F \bar{p} \Gamma n \bar{e} \Gamma \nu_e$

Coupling constant $G_F \equiv 1/\Lambda^2$ determined by fitting
the “slow” beta decay rates \Rightarrow
point to mass scale $\Lambda \gg m_n \sim \text{GeV}$

β -decays: historic perspective

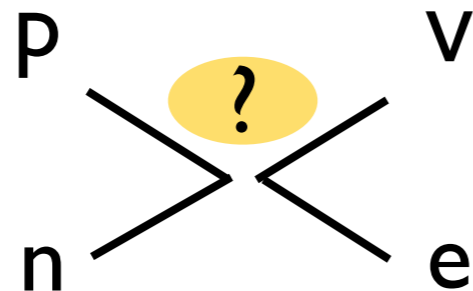
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Lee and Yang, 1956



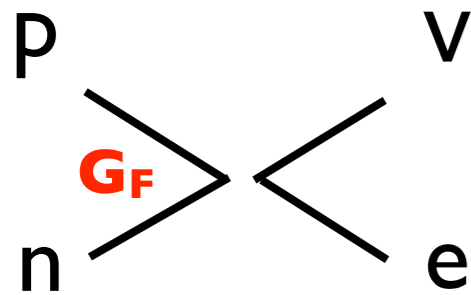
Parity conserving:
VV, AA, SS, TT ...

Parity violating: VA, SP, ...

Lee and Yang:
use most general Lorentz-
invariant interaction

β -decays: historic perspective

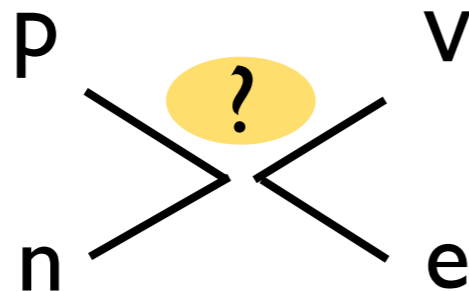
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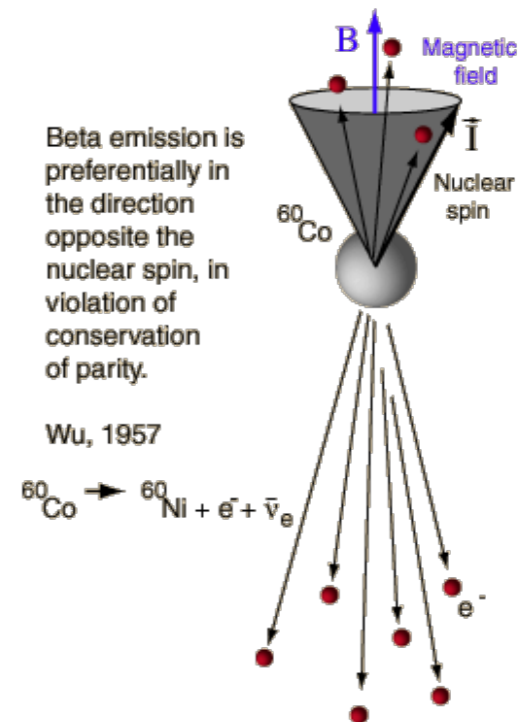
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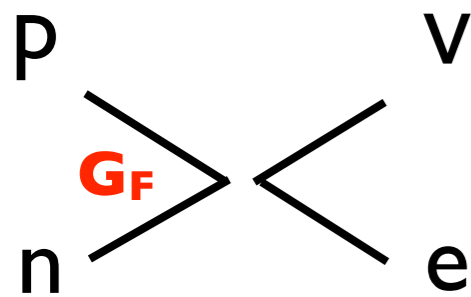


C-S Wu

Experiment: parity is violated!
(but could be VA, SP, ...)

β -decays: historic perspective

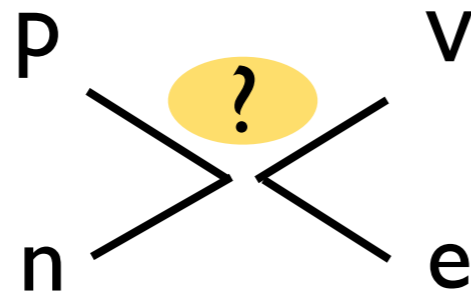
Fermi, 1934



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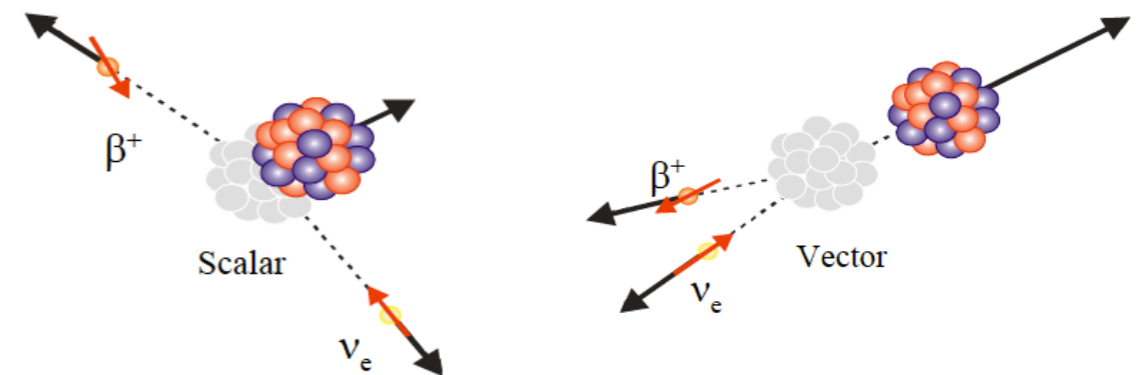
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Parity conserving:
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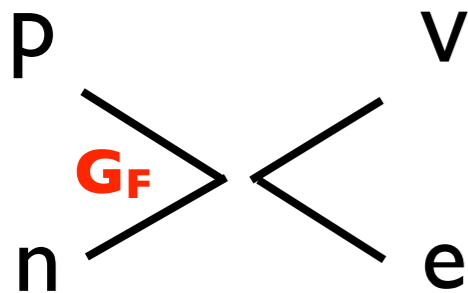
Differential decay distributions
depend on structure of currents



Model diagnosing!

β -decays: historic perspective

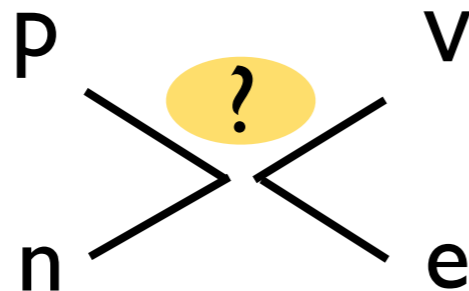
Fermi, 1934



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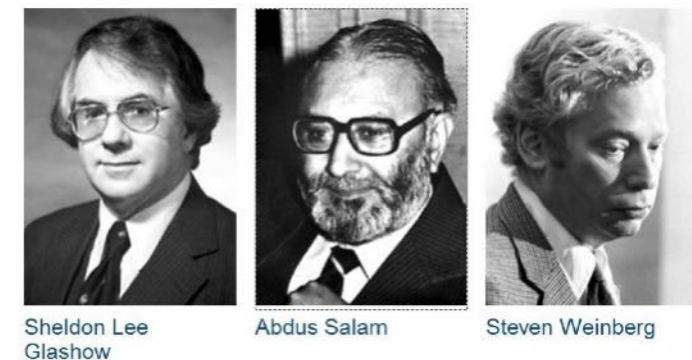


Parity conserving:
VV, AA, SS, TT ...
Parity violating: VA, SP, ...

Marshak & Sudarshan,
Feynman & Gell-Mann 1958

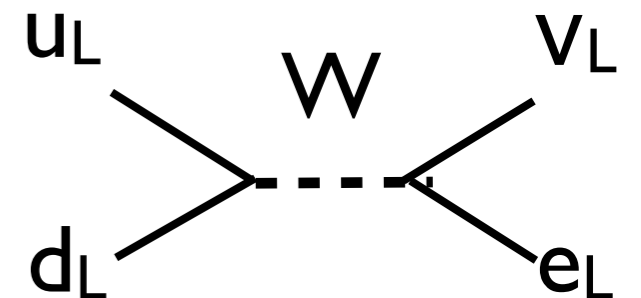


Glashow,
Salam,
Weinberg



It's $(V-A)*(V-A)$!!

"V-A was the key"
S. Weinberg



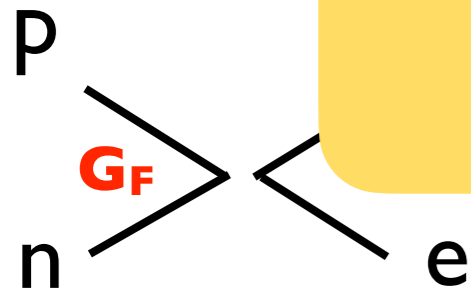
Embed in **non-abelian**
chiral gauge theory,
predict neutral currents

β -decays: historic perspective

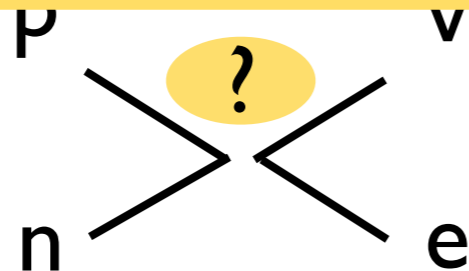
Fermi, 1934



- Lessons: nuclear beta decays were able to
- “Detect” physics originating at $\Lambda = G_F^{-1/2} \sim 250 \text{ GeV}$
 - Point to key features of the underlying interactions, that led to the formulation of the Standard Model



Current-current,
parity conserving



Parity conserving:
VV, AA, SS, TT ...
Parity violating: VA, SP, ...

Fermi scale:
 $\Lambda = G_F^{-1/2} \sim 250 \text{ GeV}$



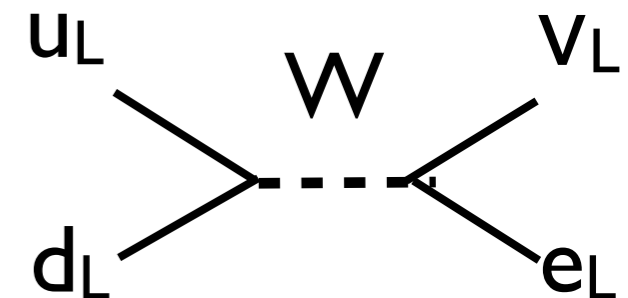
Sheldon Lee
Glashow

Abdus Salam

Steven Weinberg

It's $(V-A)*(V-A) !!$

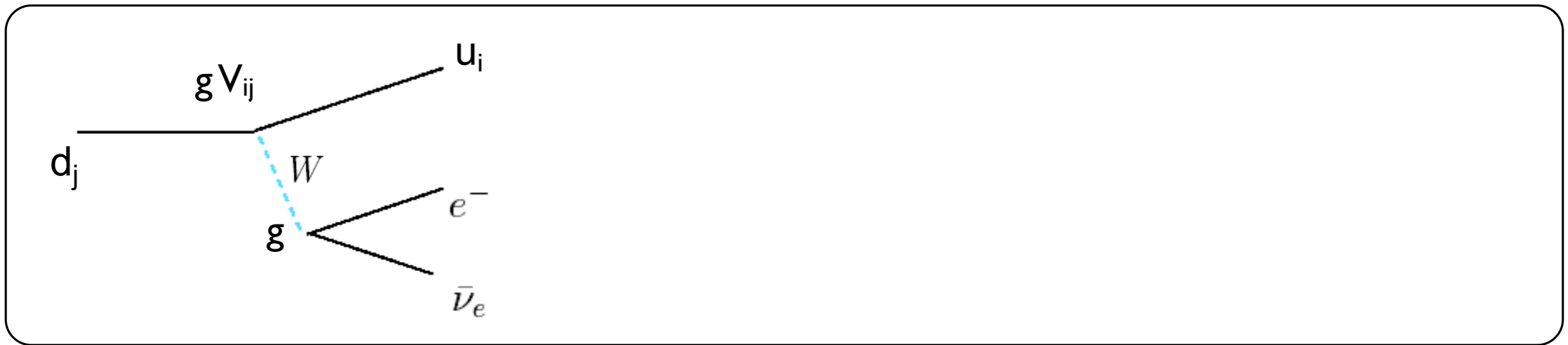
“V-A was the key”
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Embed in **non-abelian
chiral gauge theory**,
predict neutral currents

β -decays: current perspective

- In the SM, W exchange \Rightarrow universality relations



$$G_F^{(\beta)} \sim g^2 V_{ij} / M_W^2 \sim G_F^{(\mu)} V_{ij} \sim 1/v^2 V_{ij}$$

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Cabibbo-Kobayashi-Maskawa

Cabibbo Universality

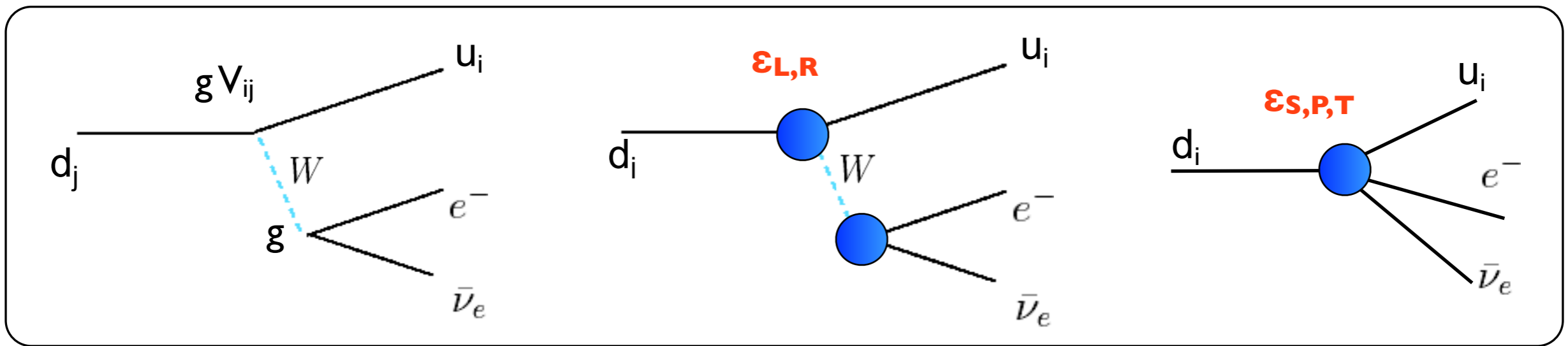
$$|V_{ud}|^2 + |V_{us}|^2 + \cancel{|V_{ub}|^2} = 1$$

$$[G_F]_e / [G_F]_\mu = 1$$

Lepton Flavor Universality (LFU)

β -decays: current perspective

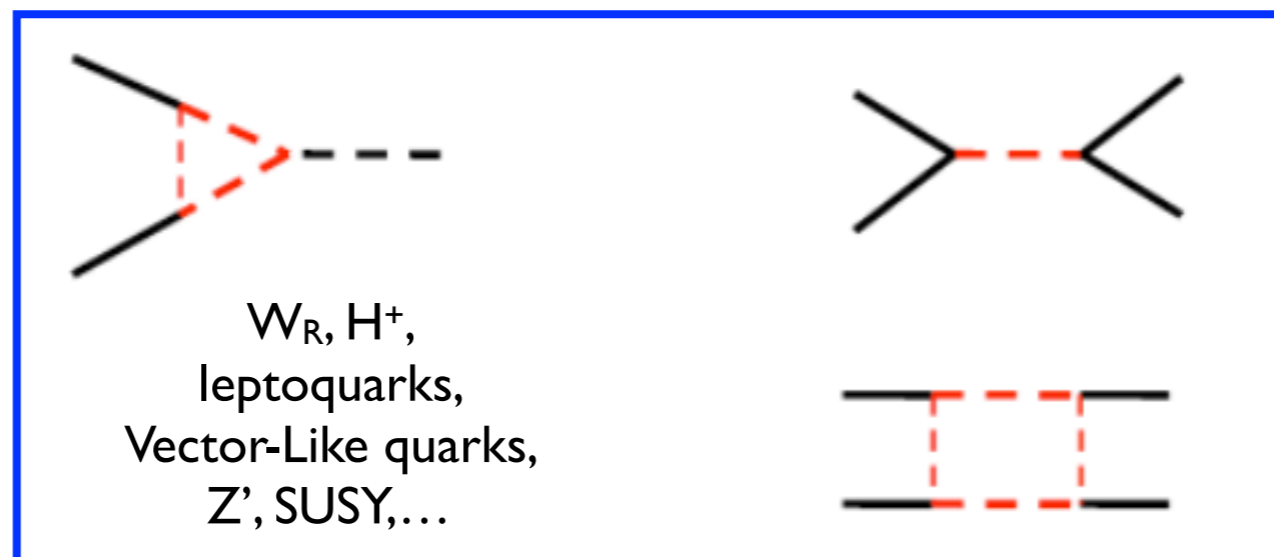
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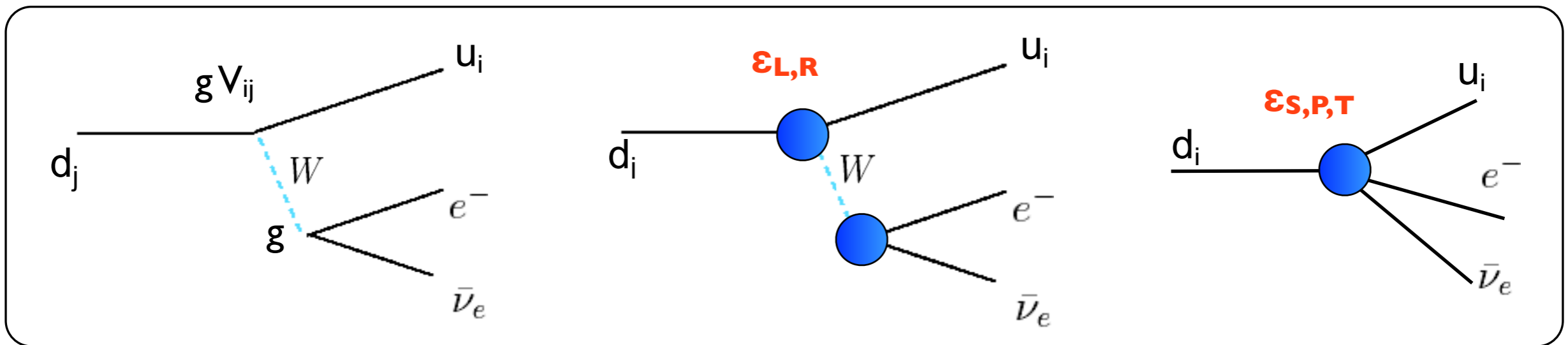
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$$E \ll \Lambda \quad \downarrow \quad \epsilon_\Gamma \sim \tilde{\epsilon}_\Gamma \sim (v/\Lambda)^2$$

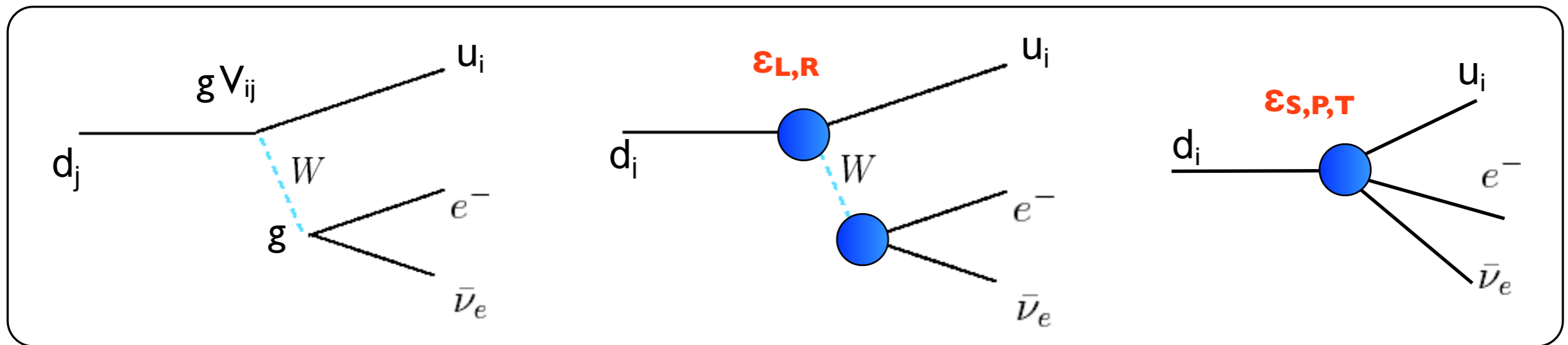
$$\mathcal{L}_{\text{SM}} = \frac{G_F V_{ud}}{\sqrt{2}} \sum_{\Gamma} \left[\epsilon_\Gamma \bar{\ell} \Gamma \nu_L \cdot \bar{u} \Gamma d + \tilde{\epsilon}_\Gamma \bar{\ell} \Gamma \nu_R \cdot \bar{u} \Gamma d \right]$$

Ten effective couplings

$$\Gamma = L, R, S, P, T$$

β -decays: current perspective

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Ten effective couplings

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- BSM effects can spoil universality. Precision of 0.1-0.01% probes $\Lambda > 10 \text{ TeV}$

Cabibbo universality tests

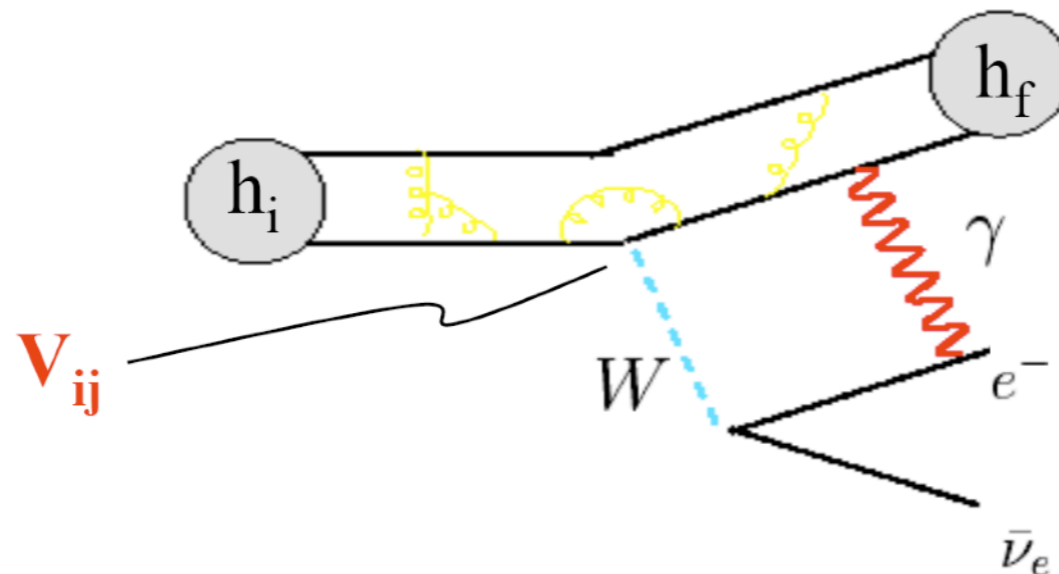
Extract $V_{ud} = \cos\theta_C$ and $V_{us} = \sin\theta_C$ from various decays

$$\Gamma_k = (G_F^{(\mu)})^2 \times |\bar{V}_{ij}|^2 \times |M_{\text{had}}|^2 \times (1 + \delta_{RC}) \times F_{\text{kin}}$$

Channel-dependent
effective CKM element

Hadronic matrix
element

Radiative corrections:
 $(\alpha/\pi) \sim 2. \times 10^{-3}$



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$$|\bar{V}_{ij}|^2 = |V_{ij}|^2 \times \left(1 + \sum_{\alpha} c_k^{\alpha} \epsilon_{\alpha} \right)$$

Calculable coefficients

BSM effective couplings

$$|\bar{V}_{ud}|^2 + |\bar{V}_{us}|^2 + |\bar{V}_{ub}|^2 = 1 + \Delta_{\text{CKM}}(\epsilon_i)$$

Paths to V_{ud} and V_{us}

V_{ud}	$0^+ \rightarrow 0^+$ ($\pi^\pm \rightarrow \pi^0 e \nu$)	$n \rightarrow p e \bar{\nu}$	$\pi \rightarrow \mu \nu$
V_{us}	$K \rightarrow \pi l \nu$	($\Lambda \rightarrow p e \bar{\nu}, \dots$)	$K \rightarrow \mu \nu$

(Hadronic
 τ decays)

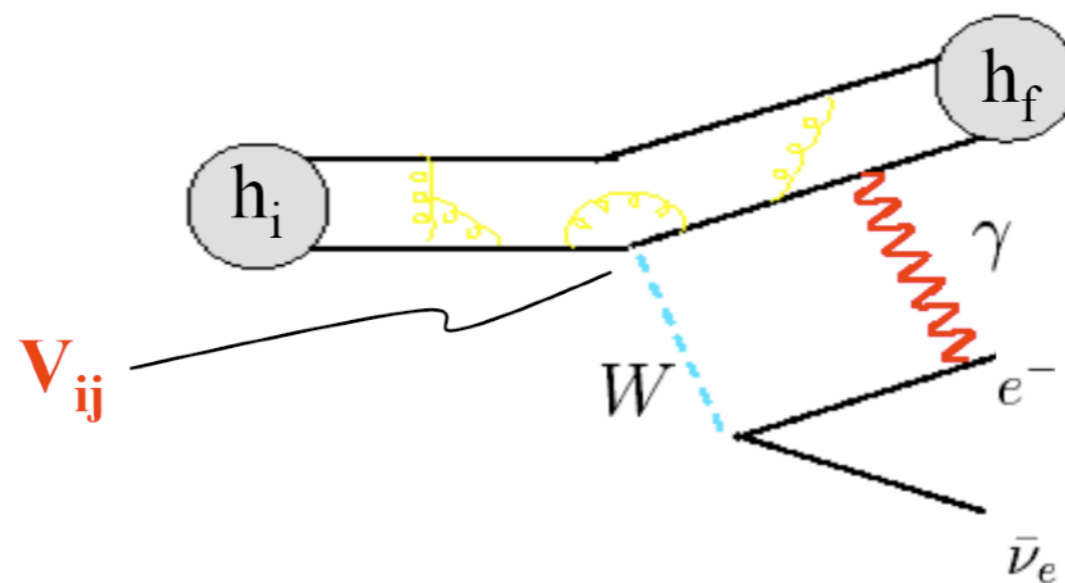
Quark current
mediating the decay



V

V, A

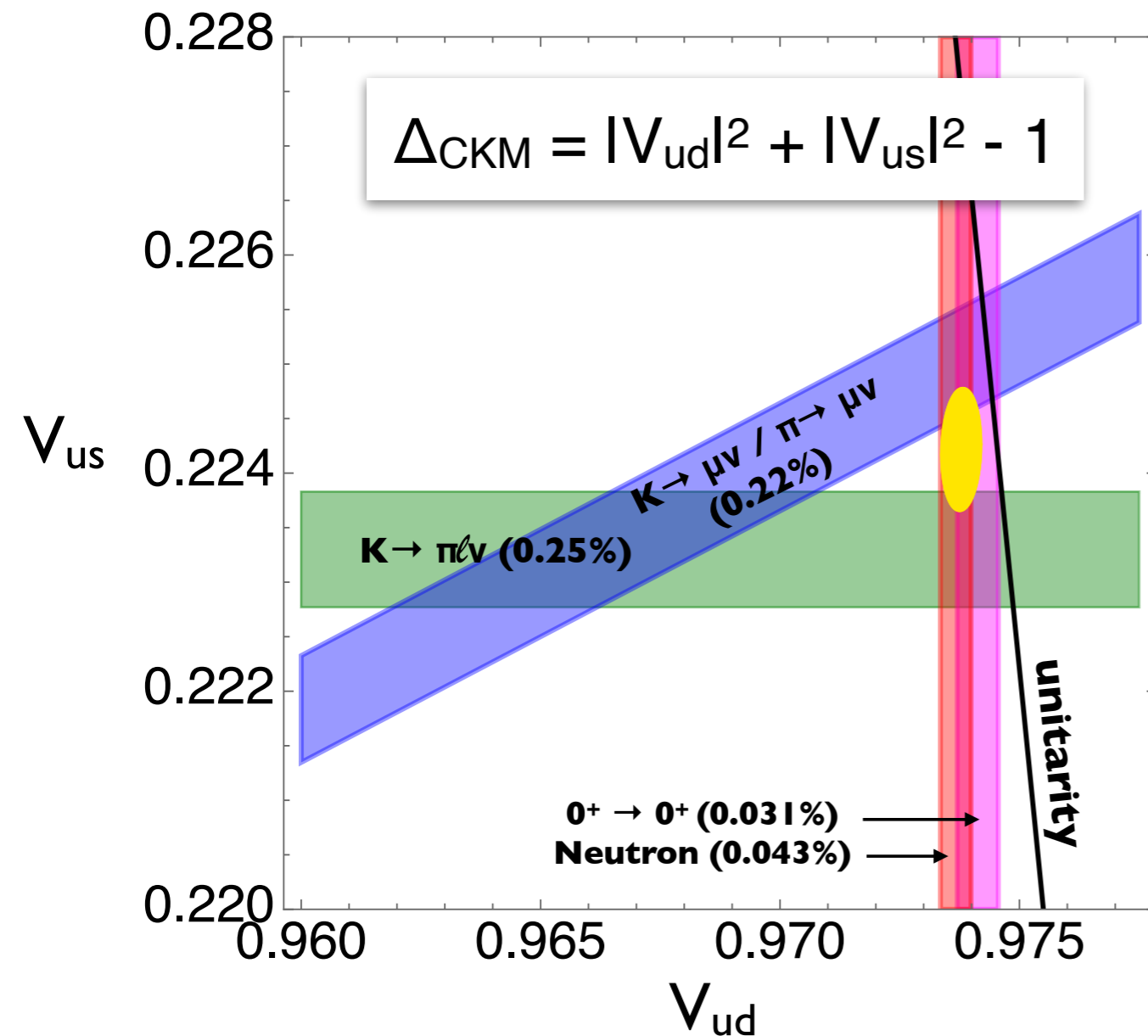
A



Input from *many* experiments and theory papers

The Cabibbo angle “anomaly”

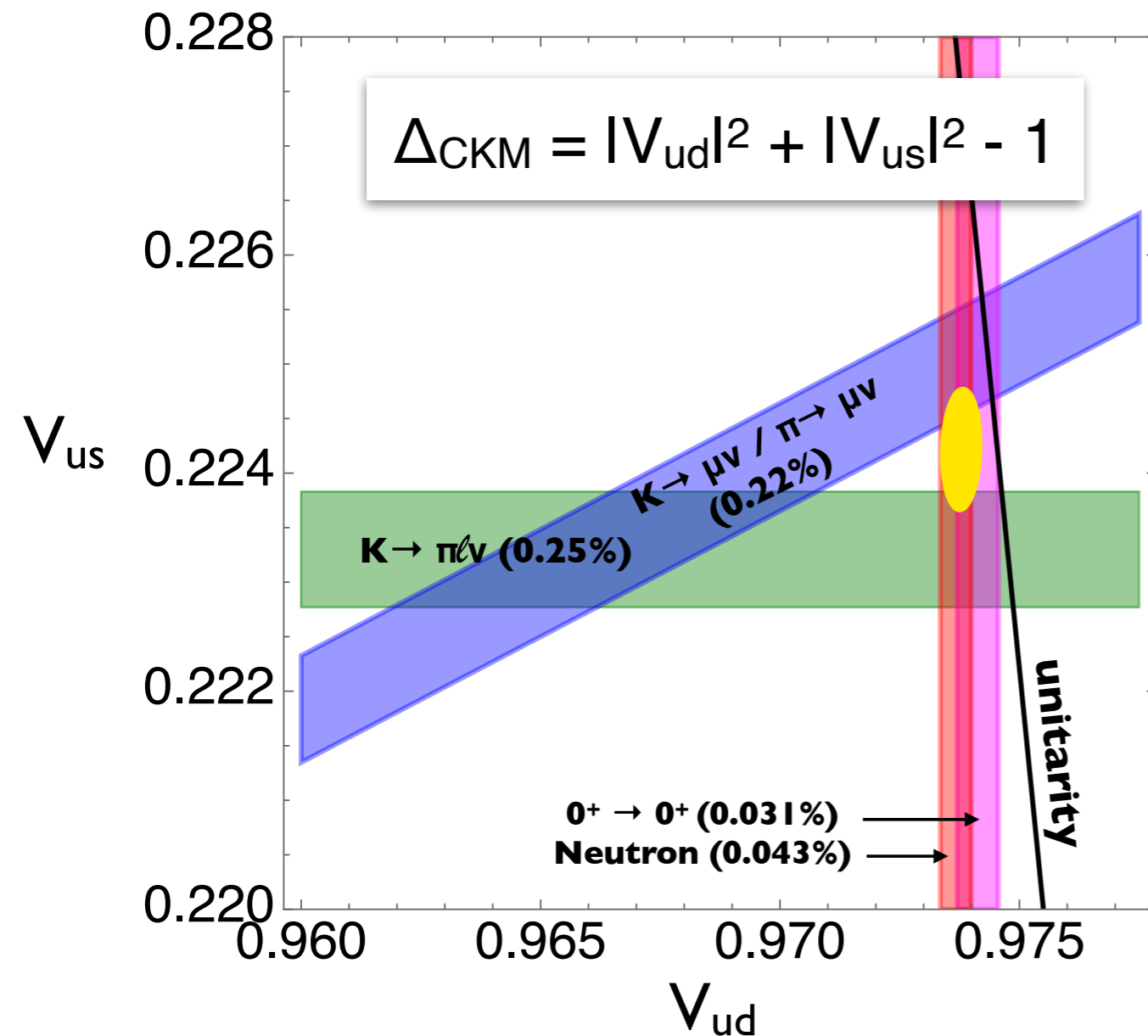
VC-Crivellini-Hoferichter-Moulson 2208.11707



- Two ‘anomalies’:
 - $\sim 3\sigma$ effect in global fit ($\Delta_{CKM} = -1.48(53) \times 10^{-3}$)
 - $\sim 3\sigma$ problem in meson sector (K12 vs K13)

The Cabibbo angle “anomaly”

VC-Crivellin-Hoferichter-Moulson 2208.11707



- Two ‘anomalies’:
 - $\sim 3\sigma$ effect in global fit ($\Delta_{CKM} = -1.48(53) \times 10^{-3}$)
 - $\sim 3\sigma$ problem in meson sector (Kl2 vs Kl3)

- Three versions of Δ_{CKM} :

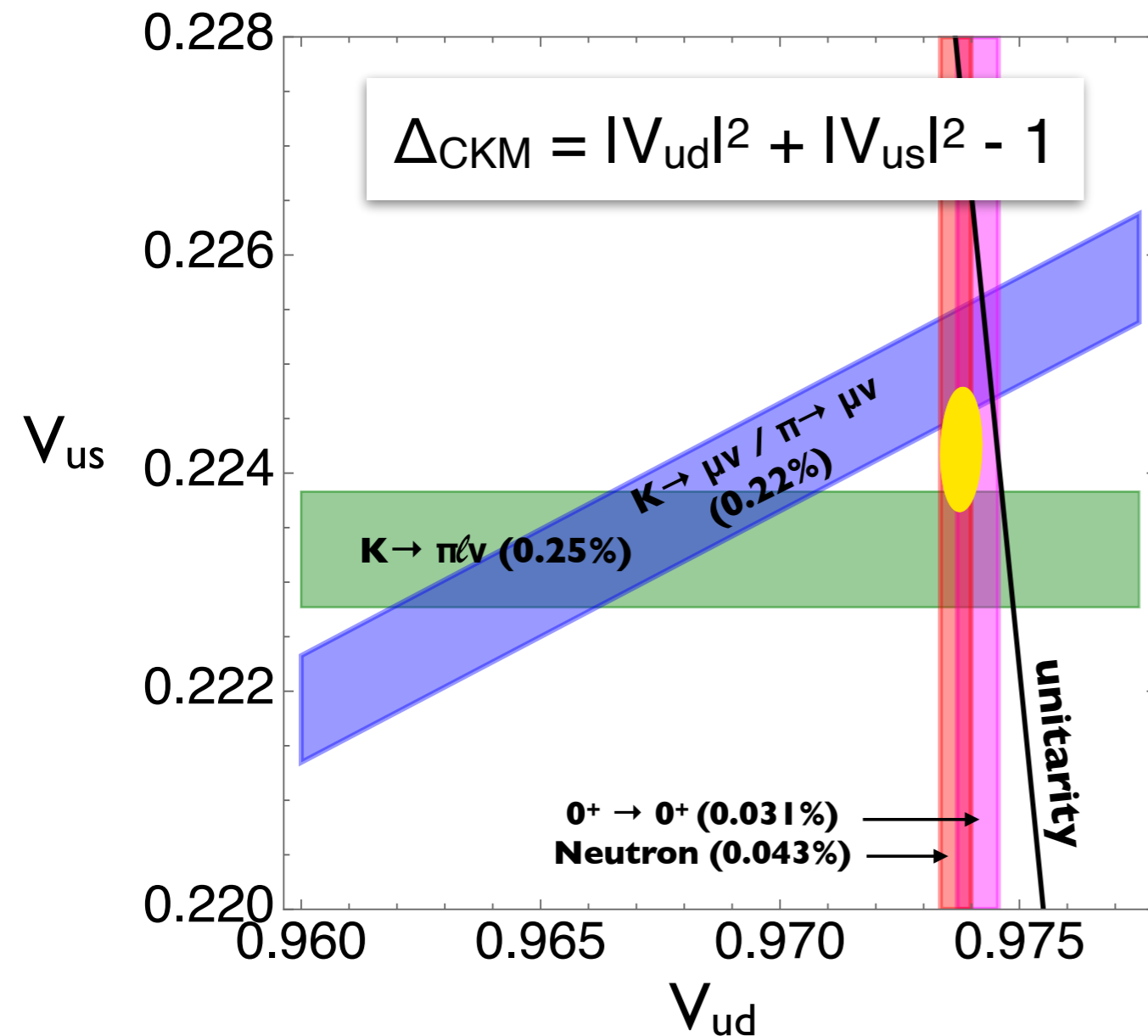
$$\begin{aligned} \Delta_{CKM}^{(1)} &= |V_{ud}^\beta|^2 + |V_{us}^{K\ell 3}|^2 - 1 \\ &= -1.76(56) \times 10^{-3} \end{aligned}$$

$$\begin{aligned} \Delta_{CKM}^{(2)} &= |V_{ud}^\beta|^2 + |V_{us}^{K\ell 2/\pi\ell 2, \beta}|^2 - 1 \\ &= -0.98(58) \times 10^{-3} \end{aligned}$$

$$\begin{aligned} \Delta_{CKM}^{(3)} &= |V_{ud}^{K\ell 2/\pi\ell 2, K\ell 3}|^2 + |V_{us}^{K\ell 3}|^2 - 1 \\ &= -1.64(63) \times 10^{-2} \end{aligned}$$

The Cabibbo angle “anomaly”

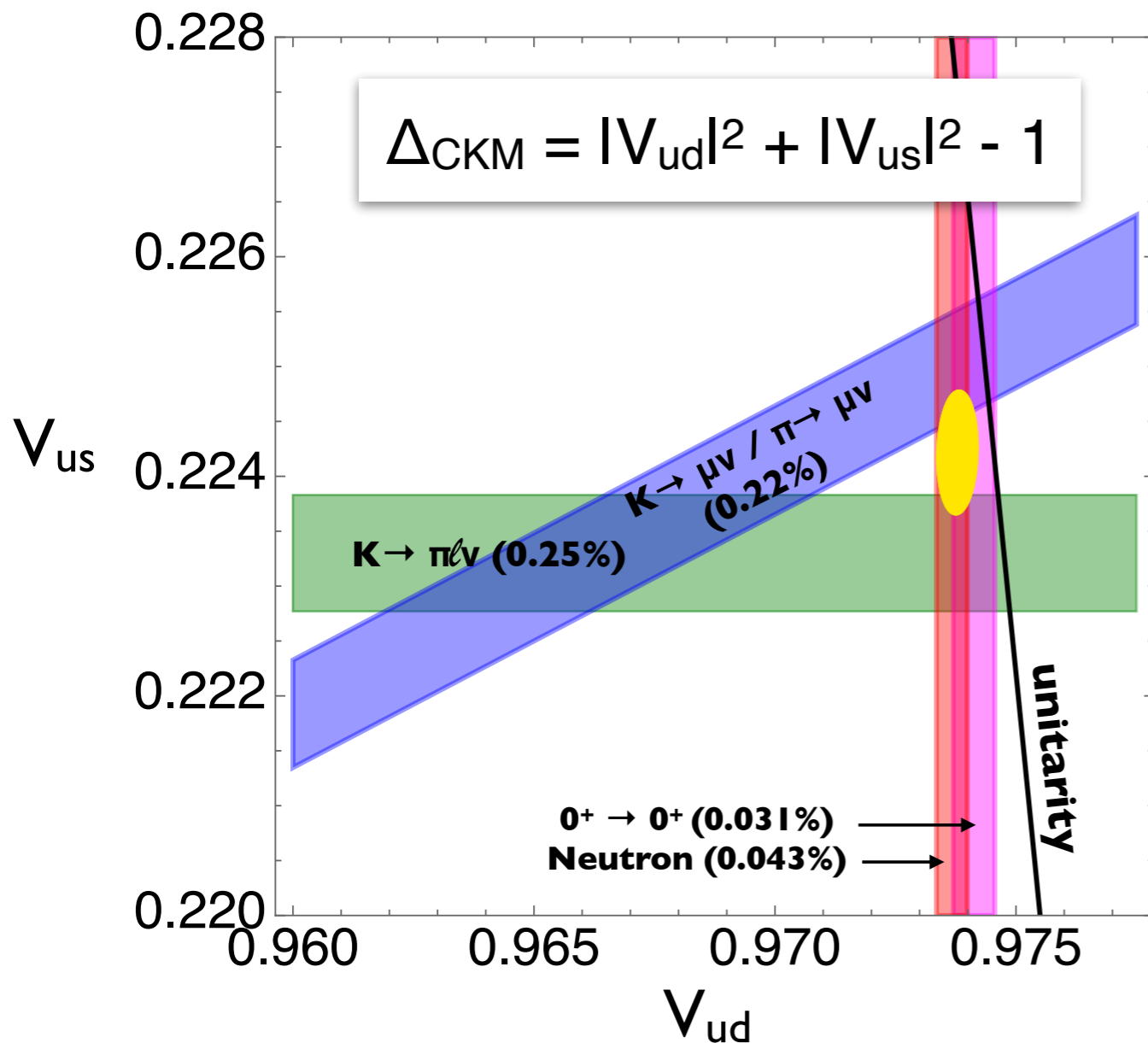
VC-Crivellin-Hoferichter-Moulson 2208.11707



- **Expected experimental improvements:**
 - neutron decay (will match nominal nuclear uncertainty)
 - pion decay (3x to 10x at PIONEER phases II, III)
 - possibly new K BR measurements at NA62
- **Expected theoretical improvements:**
 - radiative corrections in QCD+QED with lattice gauge theory for $Kl3$ and neutron;
 - EFT-based first-principles nuclear structure corrections

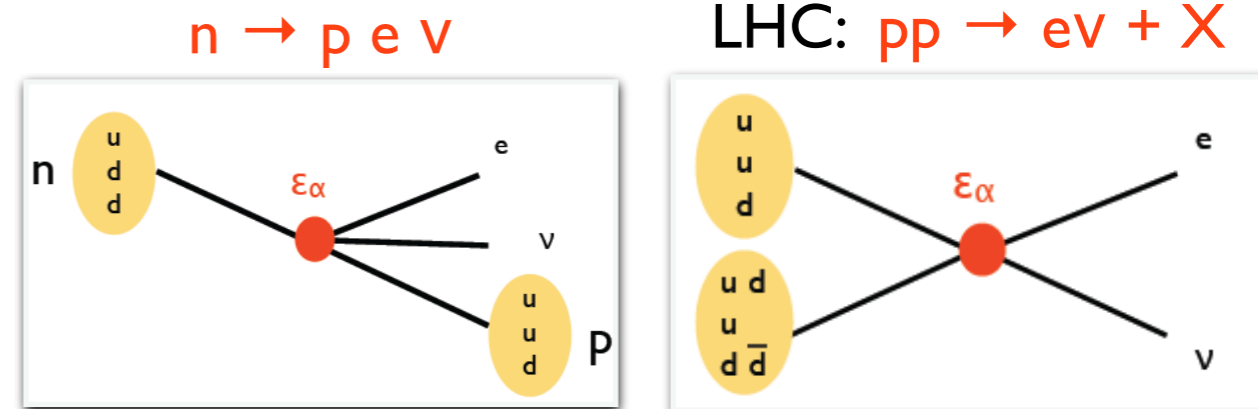
The Cabibbo angle “anomaly”

VC-Crivellin-Hoferichter-Moulson 2208.11707



- Possible BSM explanations:

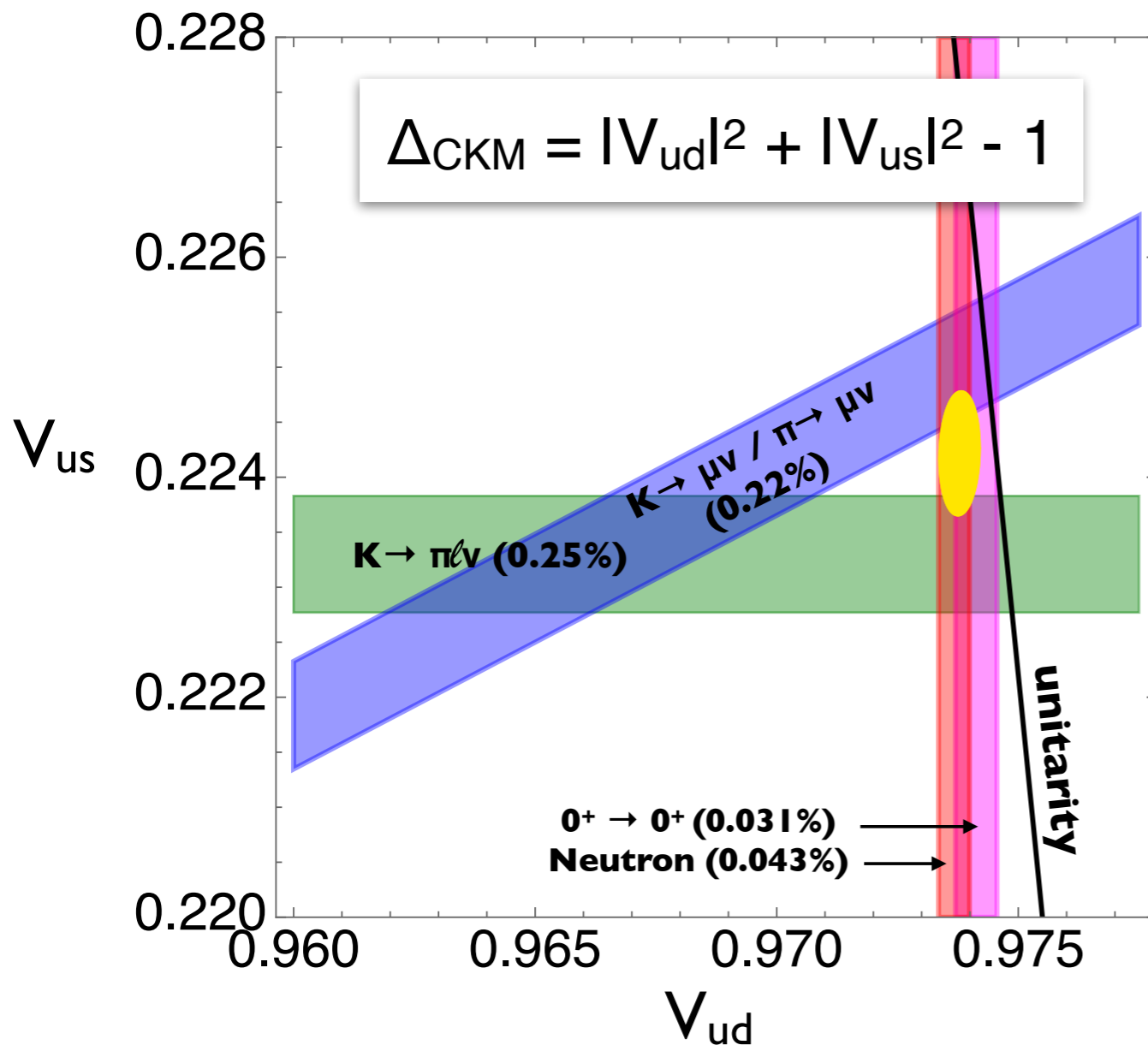
- 4-fermion operators strongly constrained by LHC measurements



VC, Gonzalez-Alonso, Graesser 1210.4553

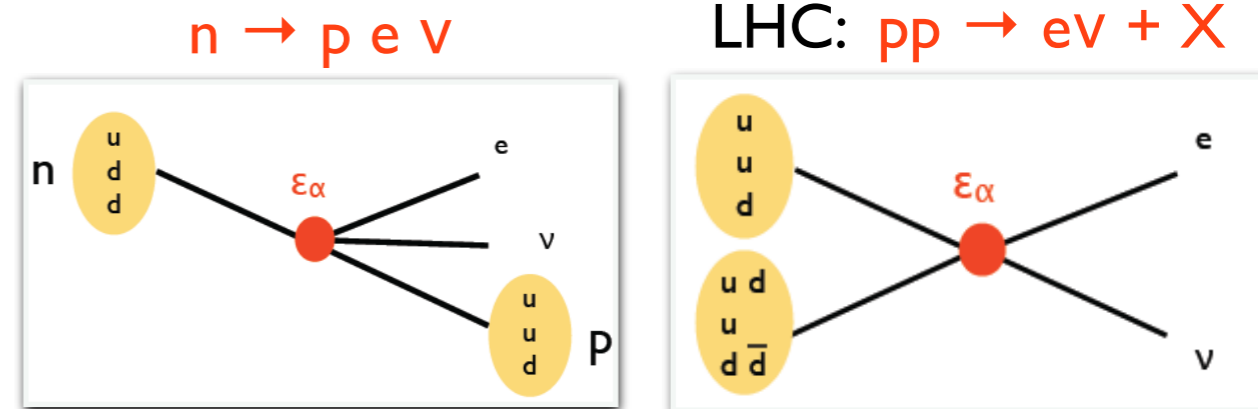
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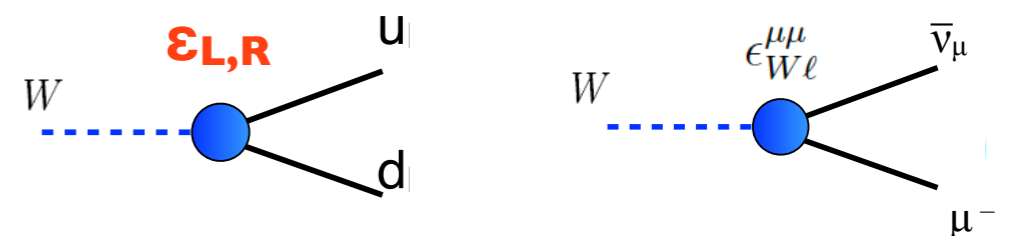
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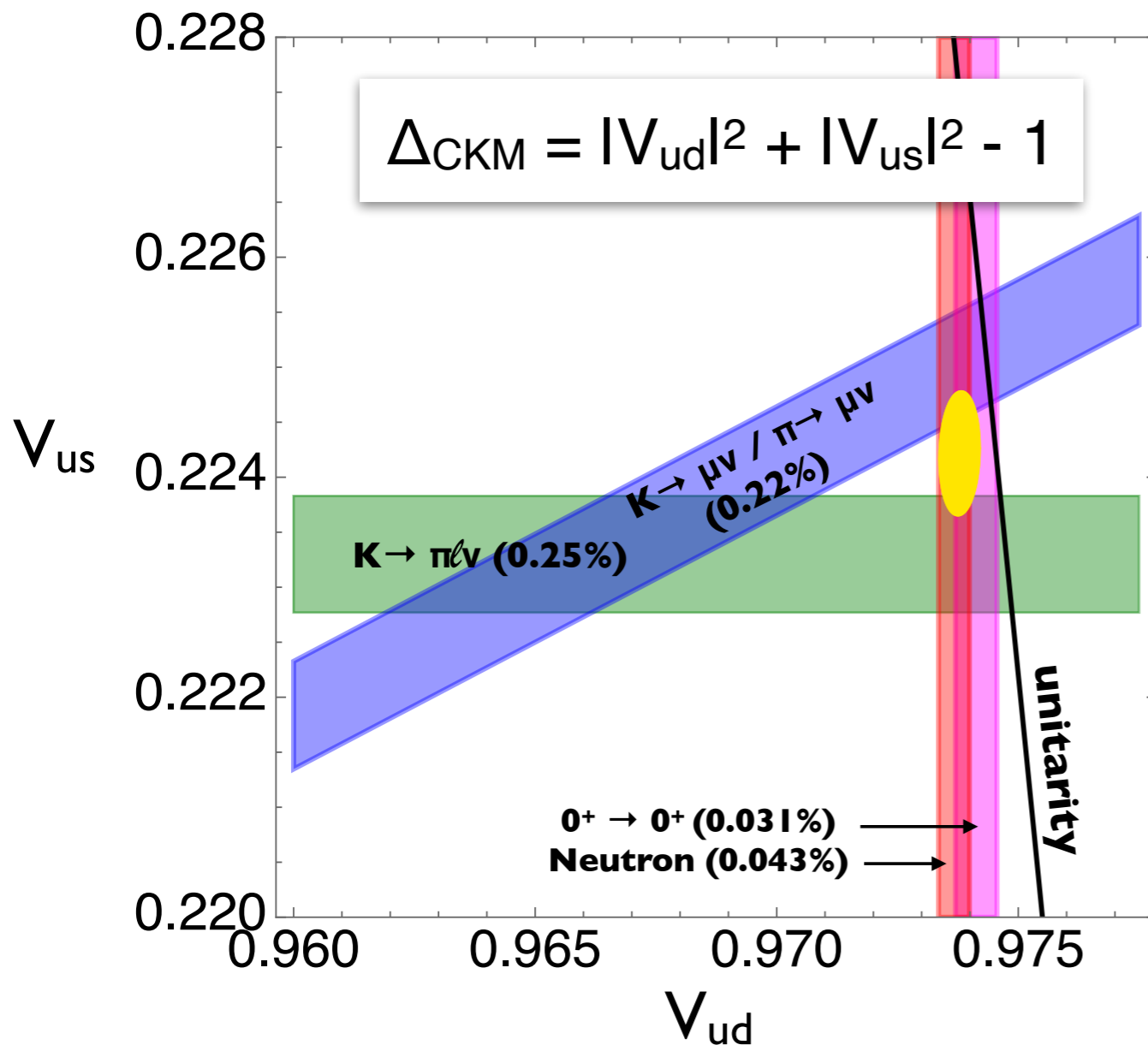
- Vertex corrections (quark and leptons) remain viable candidates



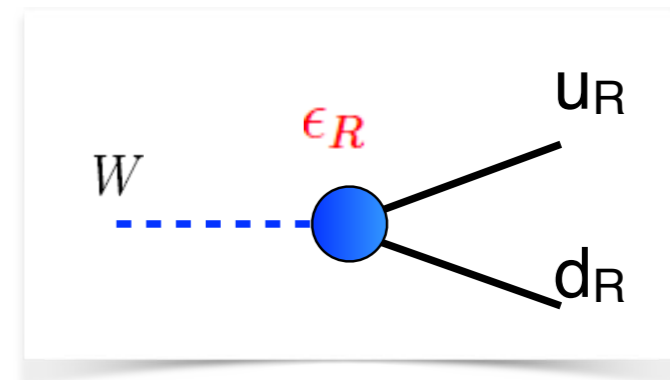
For mini-review see A. Crivellin 2207.02507

The Cabibbo angle “anomaly”

VC-Crivellini-Hoferichter-Moulson 2208.11707



- All discrepancies can be solved by R-handed quark currents (both ‘ud’ and ‘us’): ϵ_R , $\epsilon_R^{(s)}$

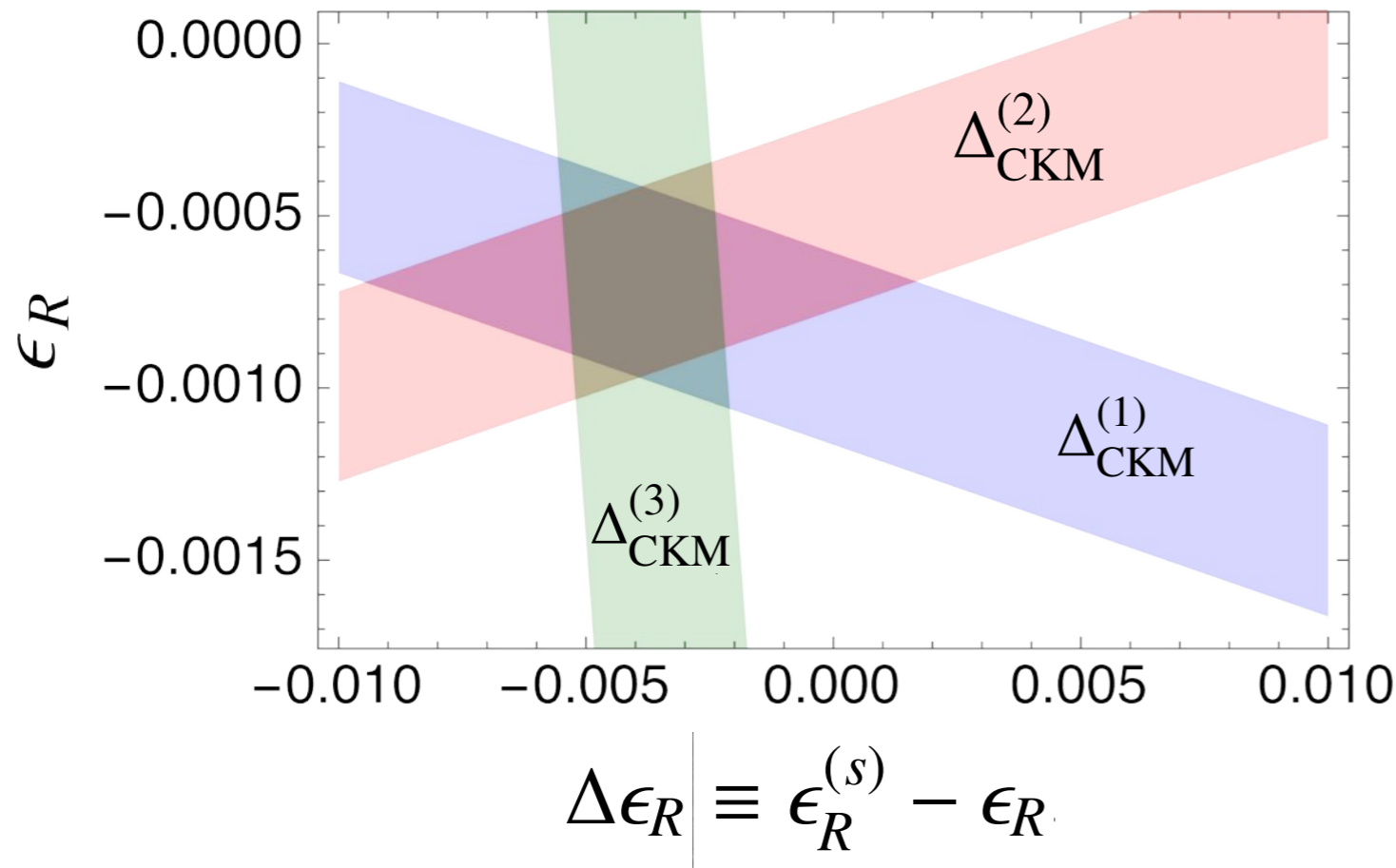


- CKM elements extracted from vector (axial) channels are shifted by $1 + \epsilon_R$ ($1 - \epsilon_R$)

Alioli et al 1703.04751
 Grossman-Passemar-Schacht 1911.07821
 VC, Diaz-Calderon, et al, 2112.02087
 Belfatto-Berezhiani 2103.05549

Unveiling R-handed quark currents?

VC-Crivellin-Hoferichter-Moulson 2208.11707



$$\begin{aligned}\Delta_{\text{CKM}}^{(1)} &= 2\epsilon_R + 2\Delta\epsilon_R V_{us}^2, \\ \Delta_{\text{CKM}}^{(2)} &= 2\epsilon_R - 2\Delta\epsilon_R V_{us}^2, \\ \Delta_{\text{CKM}}^{(3)} &= 2\epsilon_R + 2\Delta\epsilon_R(2 - V_{us}^2)\end{aligned}$$

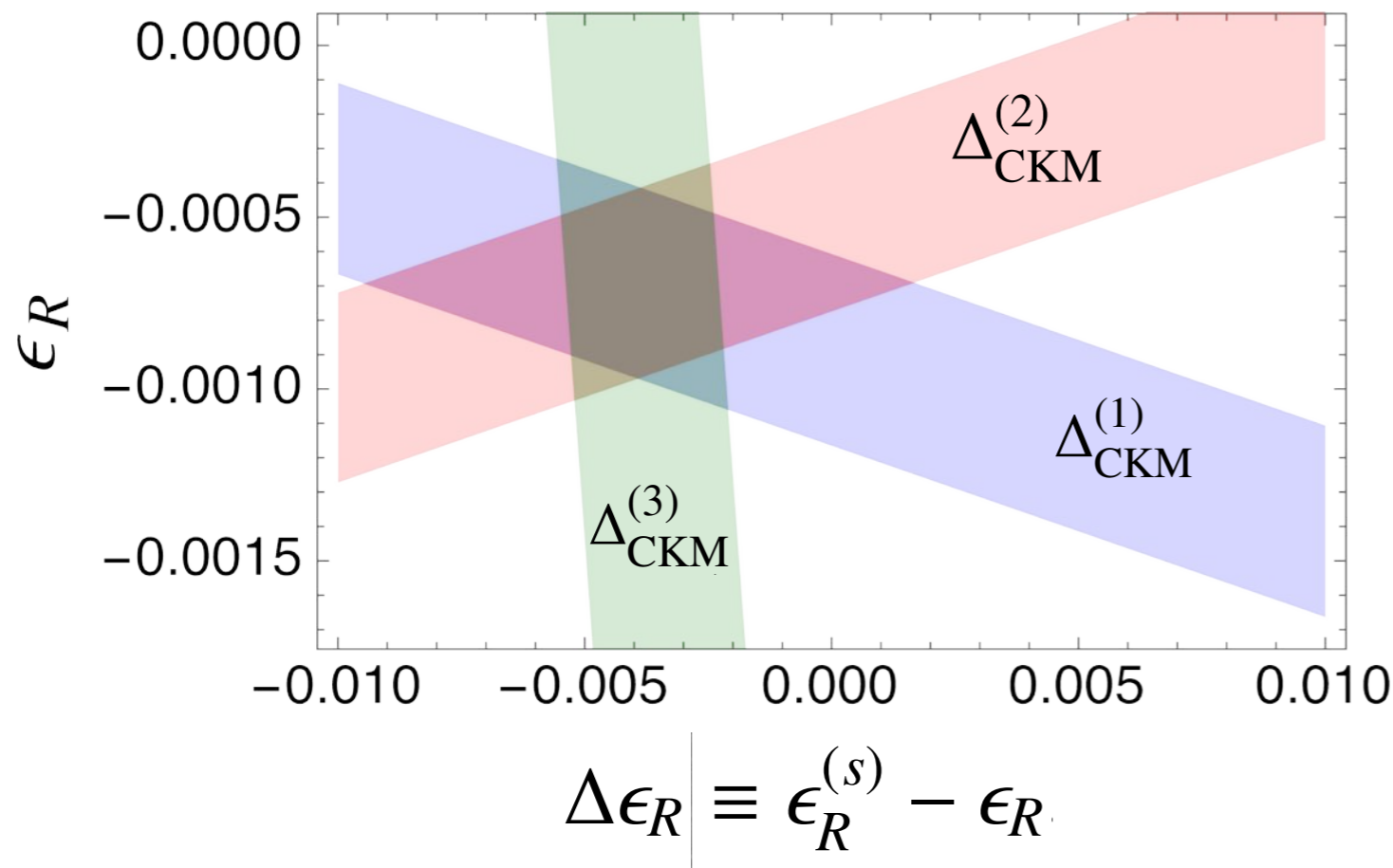


$$\begin{aligned}\epsilon_R &= -0.69(27) \times 10^{-3} \\ \Delta\epsilon_R &= -3.9(1.6) \times 10^{-3}\end{aligned}$$

$\Lambda_R \sim 5-10 \text{ TeV}$ 2.5σ effect

Unveiling R-handed quark currents?

VC-Crivellin-Hoferichter-Moulson 2208.11707



- Preferred ranges are not in conflict with other data from β decays or the LHC, such as $pp \rightarrow e\nu + X$ and $pp \rightarrow W h + X$ (projected constraints at few % level)

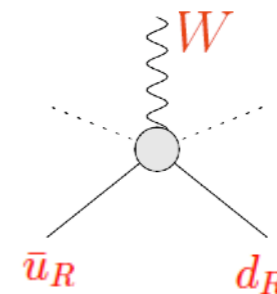
$$\begin{aligned}\Delta_{\text{CKM}}^{(1)} &= 2\epsilon_R + 2\Delta\epsilon_R V_{us}^2, \\ \Delta_{\text{CKM}}^{(2)} &= 2\epsilon_R - 2\Delta\epsilon_R V_{us}^2, \\ \Delta_{\text{CKM}}^{(3)} &= 2\epsilon_R + 2\Delta\epsilon_R(2 - V_{us}^2)\end{aligned}$$



$$\begin{aligned}\epsilon_R &= -0.69(27) \times 10^{-3} \\ \Delta\epsilon_R &= -3.9(1.6) \times 10^{-3}\end{aligned}$$

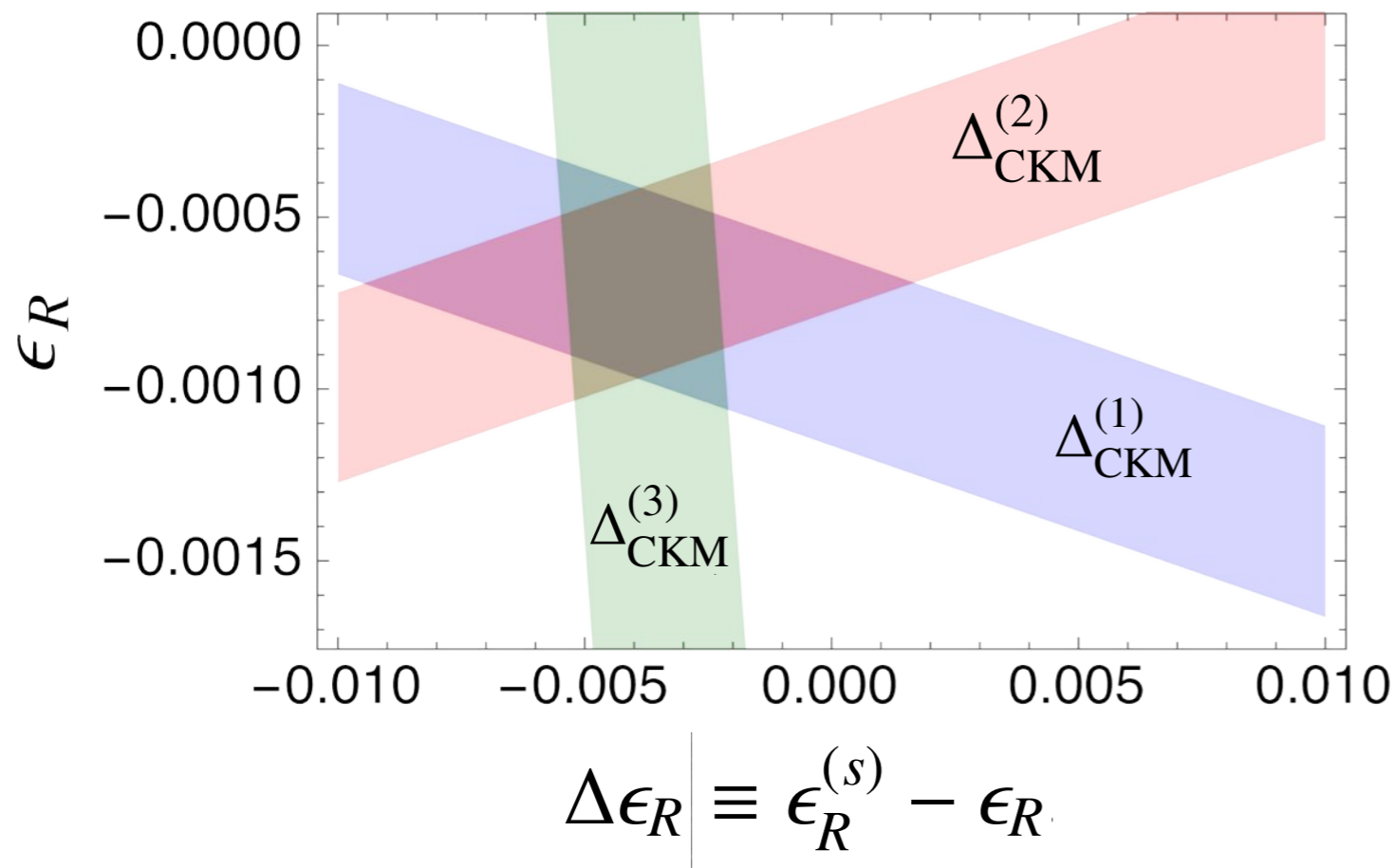
$\Lambda_R \sim 5-10 \text{ TeV}$ 2.5σ effect

$$O_{\varphi\varphi} = i(\varphi^T \epsilon D_\mu \varphi)(\bar{u}\gamma^\mu d)$$



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$$\begin{aligned}\Delta_{\text{CKM}}^{(1)} &= 2\epsilon_R + 2\Delta\epsilon_R V_{us}^2, \\ \Delta_{\text{CKM}}^{(2)} &= 2\epsilon_R - 2\Delta\epsilon_R V_{us}^2, \\ \Delta_{\text{CKM}}^{(3)} &= 2\epsilon_R + 2\Delta\epsilon_R(2 - V_{us}^2)\end{aligned}$$



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$$\Lambda_R \sim 5-10 \text{ TeV} \quad 2.5\sigma \text{ effect}$$

- A measurement of $\text{BR}(K \rightarrow \pi\mu\nu)/\text{BR}(K \rightarrow \mu\nu)$ at 0.2% level (possible at NA62, CERN) will corroborate or rule out the presence of R-handed effects
- CKM unitarity test is a very competitive and compelling probe of multi-TeV scale physics well in the LHC era!

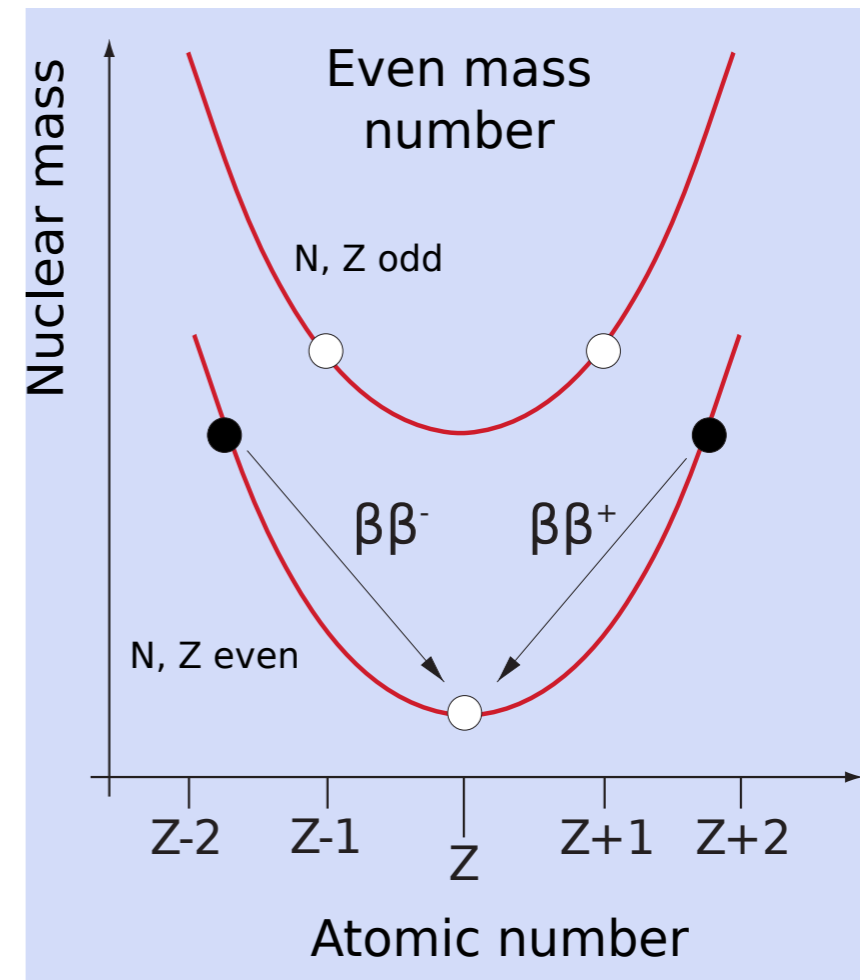
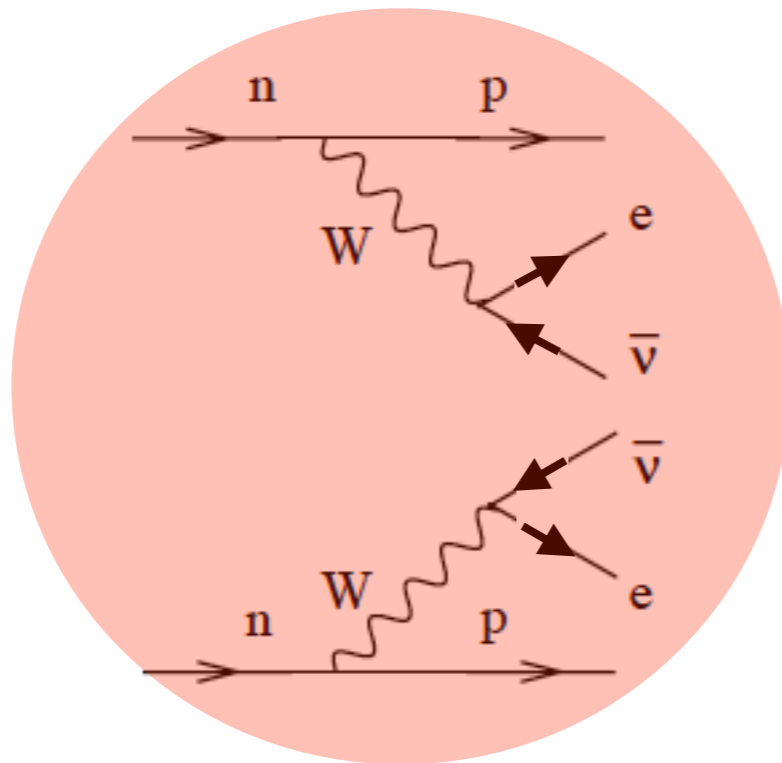
Neutrinoless double beta decay

Double beta decay

- For certain even-even nuclei (^{48}Ca , ^{76}Ge , ^{136}Xe , ...), single β decay is energetically forbidden \rightarrow $\beta\beta$ decay!



M. Goppert Mayer, 1935



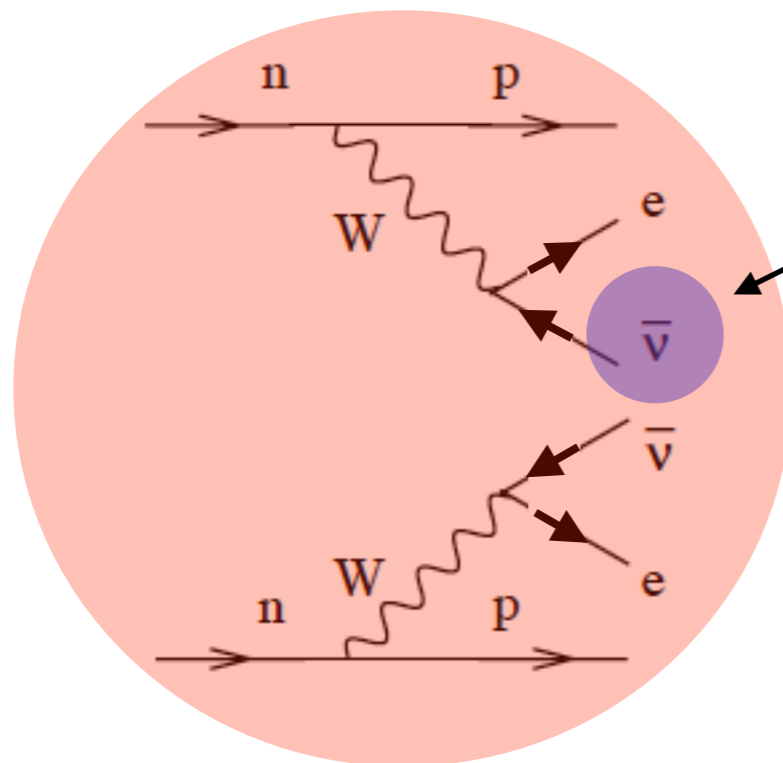
- $2\nu\beta\beta$ is the rarest process ever observed, with $T_{1/2} \sim 10^{21}$ years (first observation in 1987)

Neutrinoless double beta decay?

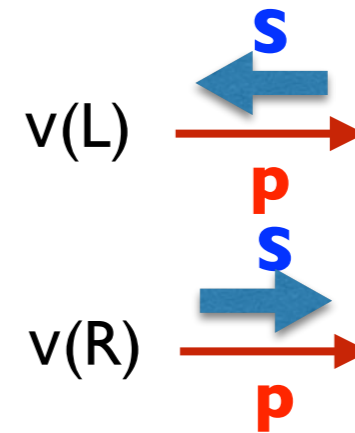
- Yes, if neutrinos are massive Majorana particles (i.e. their own antiparticles)



W. H. Furry, 1939



This is just $\nu(R)$, which mixes with $\nu(L)$ via mass insertion



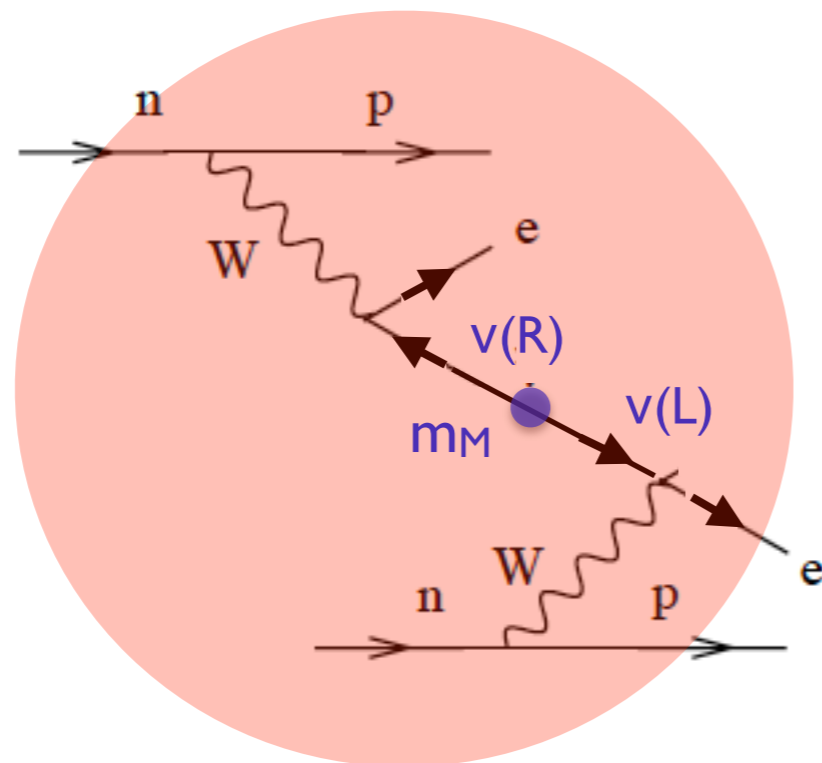
E. Majorana, 1937

Neutrinoless double beta decay?

- Yes, if neutrinos are massive Majorana particles (i.e. their own antiparticles)



W. H. Furry, 1939



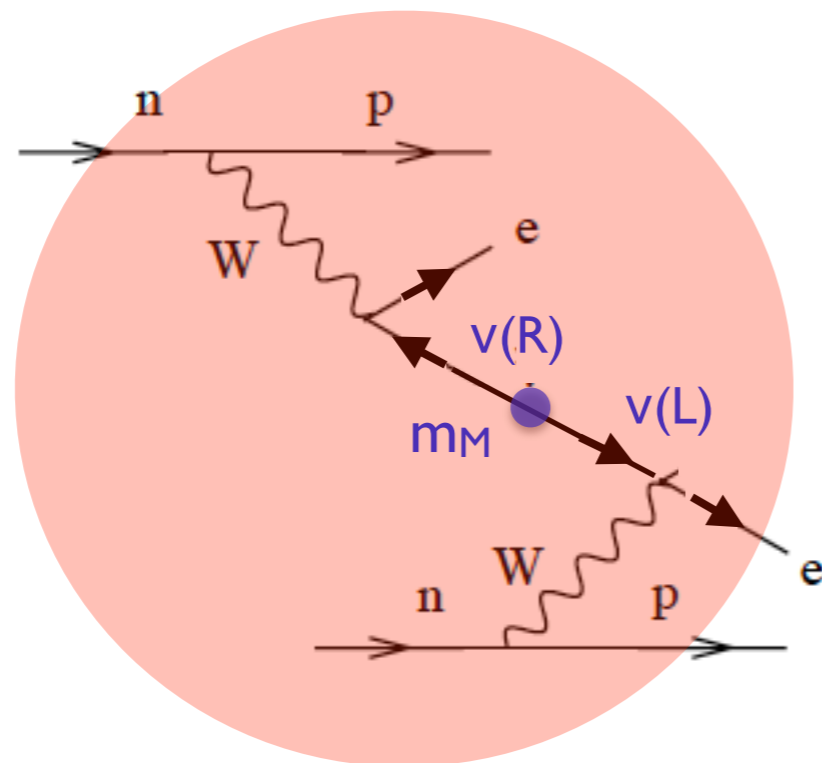
“Subject to the usual limitations on the meaning of such language, one can say that a (virtual) neutrino is emitted together with one of the electrons and reabsorbed when the other electron is emitted.”

Neutrinoless double beta decay?

- Yes, if neutrinos are massive Majorana particles (i.e. their own antiparticles)



W. H. Furry, 1939



“Subject to the usual limitations on the meaning of such language, one can say that a (virtual) neutrino is emitted together with one of the electrons and reabsorbed when the other electron is emitted.”

- Key point: in $0\nu\beta\beta$ Lepton Number changes by two units. Majorana ν exchange is just one possible mechanism. Furry understood this:

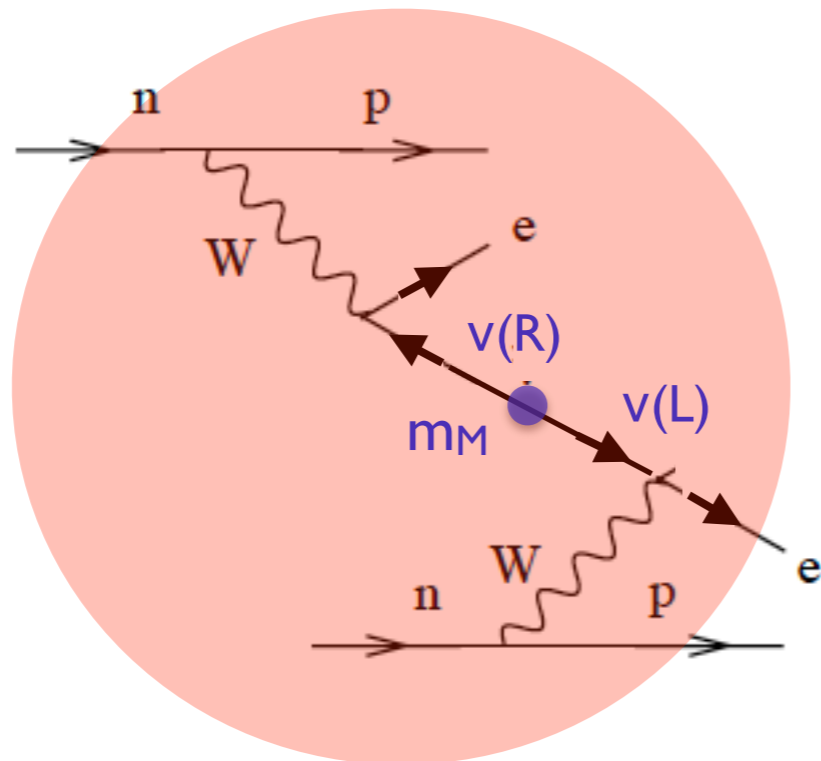
“The Majorana form of the theory is not the only one that permits this new form of disintegration [...]. The Majorana theory provides, so to speak, a canonical form.”

Neutrinoless double beta decay?

- Yes, if neutrinos are massive Majorana particles (i.e. their own antiparticles)

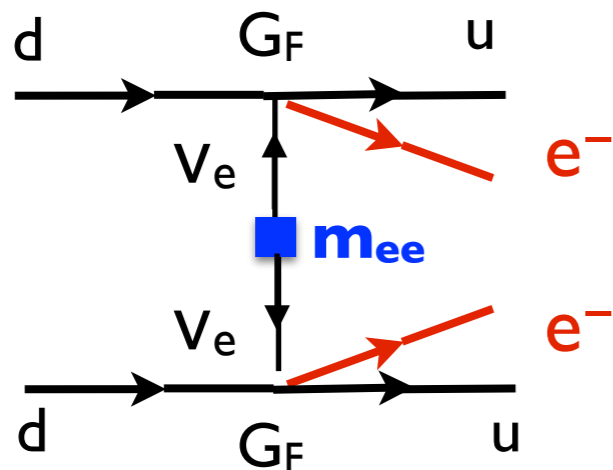


W. H. Furry, 1939

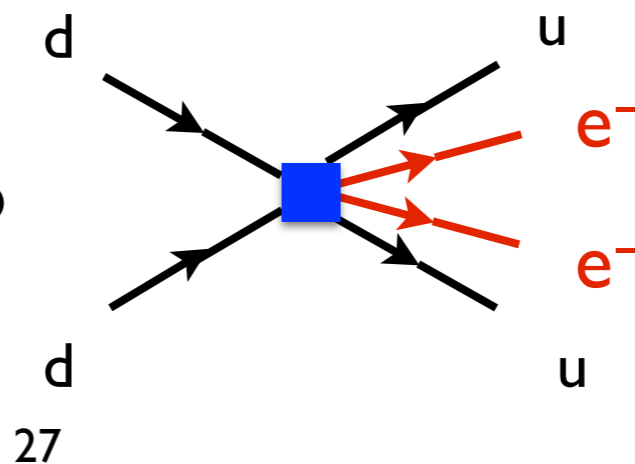


“Subject to the usual limitations on the meaning of such language, one can say that a (virtual) neutrino is emitted together with one of the electrons and reabsorbed when the other electron is emitted.”

- Modern viewpoint on Lepton Number Violation:



but also

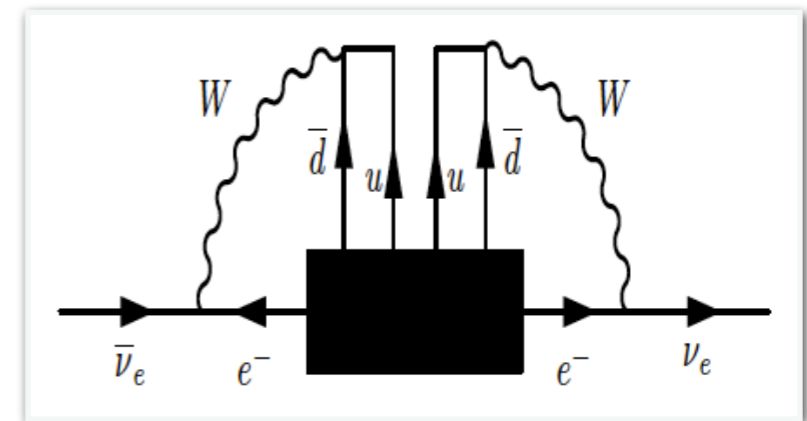


Exchange of heavier neutrinos or other Majorana particles. At low-energy induce six-fermion operator $\sim 1/\Lambda^5$

Significance of $0\nu\beta\beta$

- B-L conserved in SM $\rightarrow 0\nu\beta\beta =$ new physics, with far-reaching implications

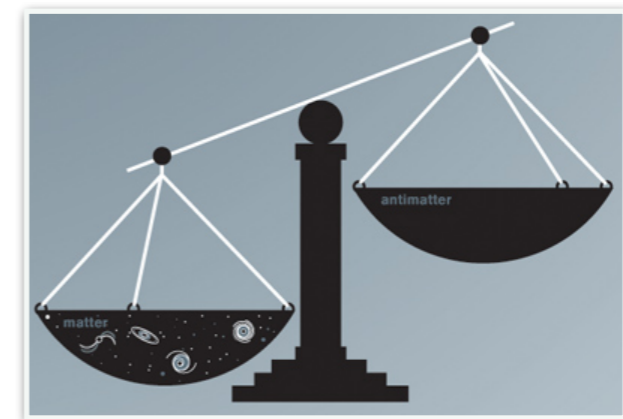
- Demonstrate that neutrinos are their own antiparticles (Majorana fermions)



Shechter-Valle 1982

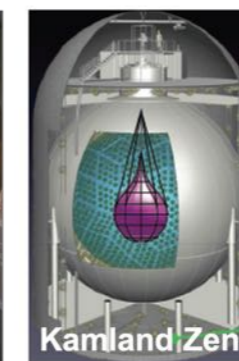
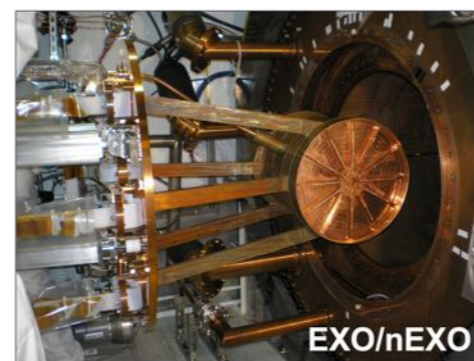
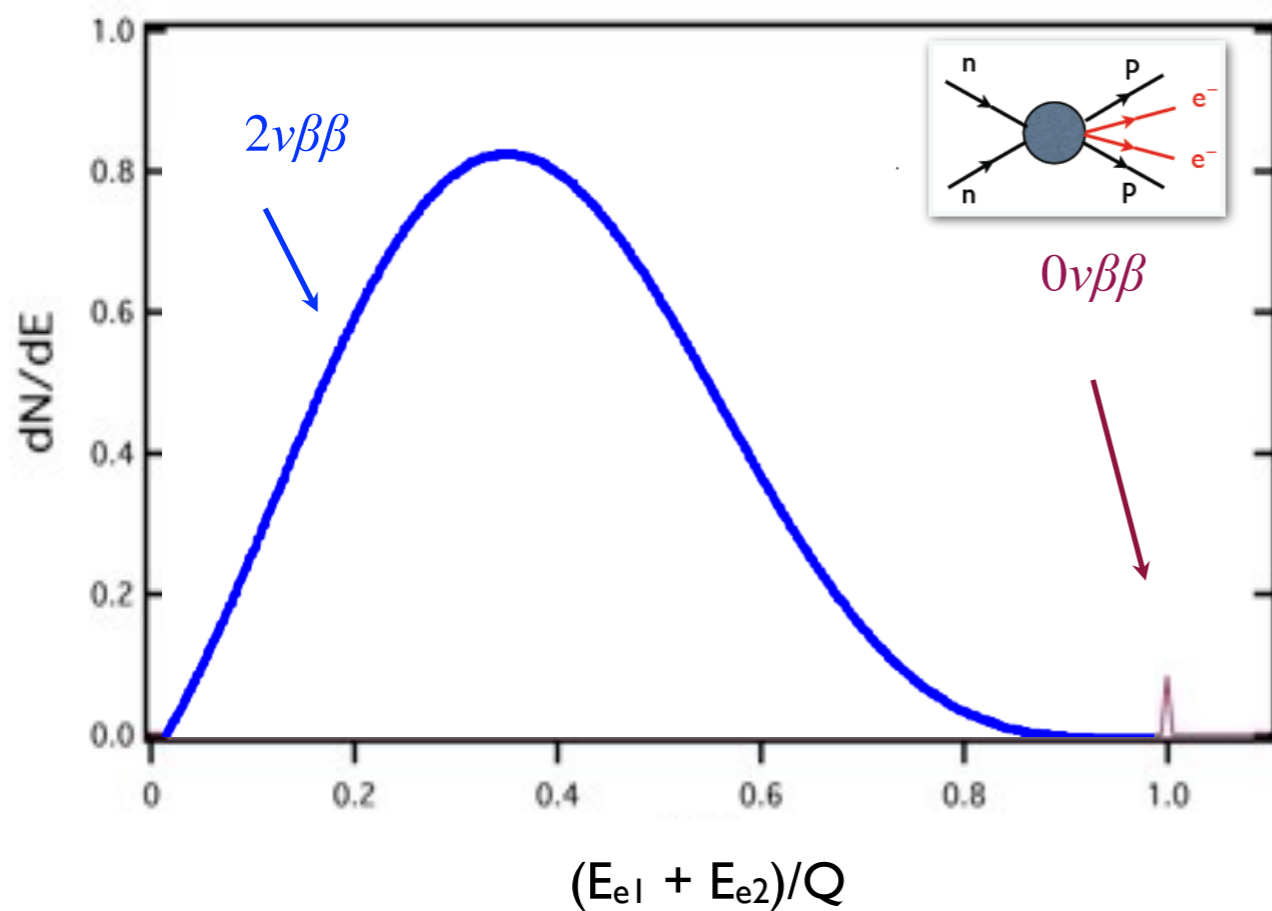
- Establish L non-conservation, key ingredient to generate the baryon asymmetry via leptogenesis

Fukujita-Yanagida 1987



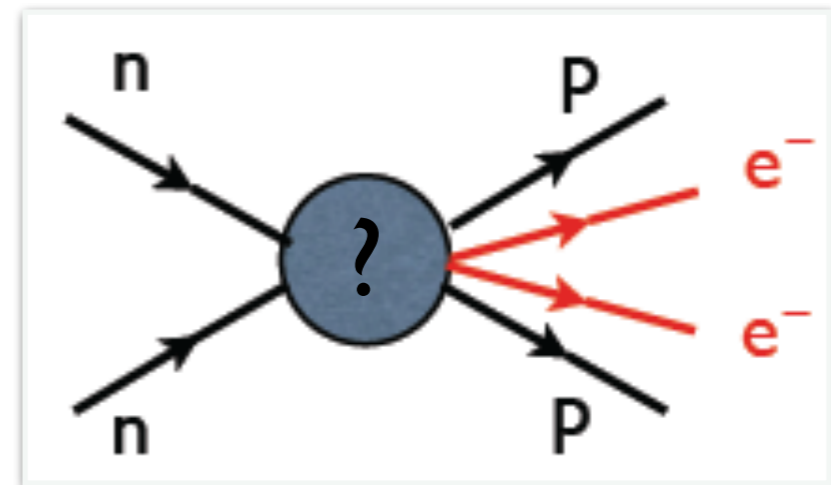
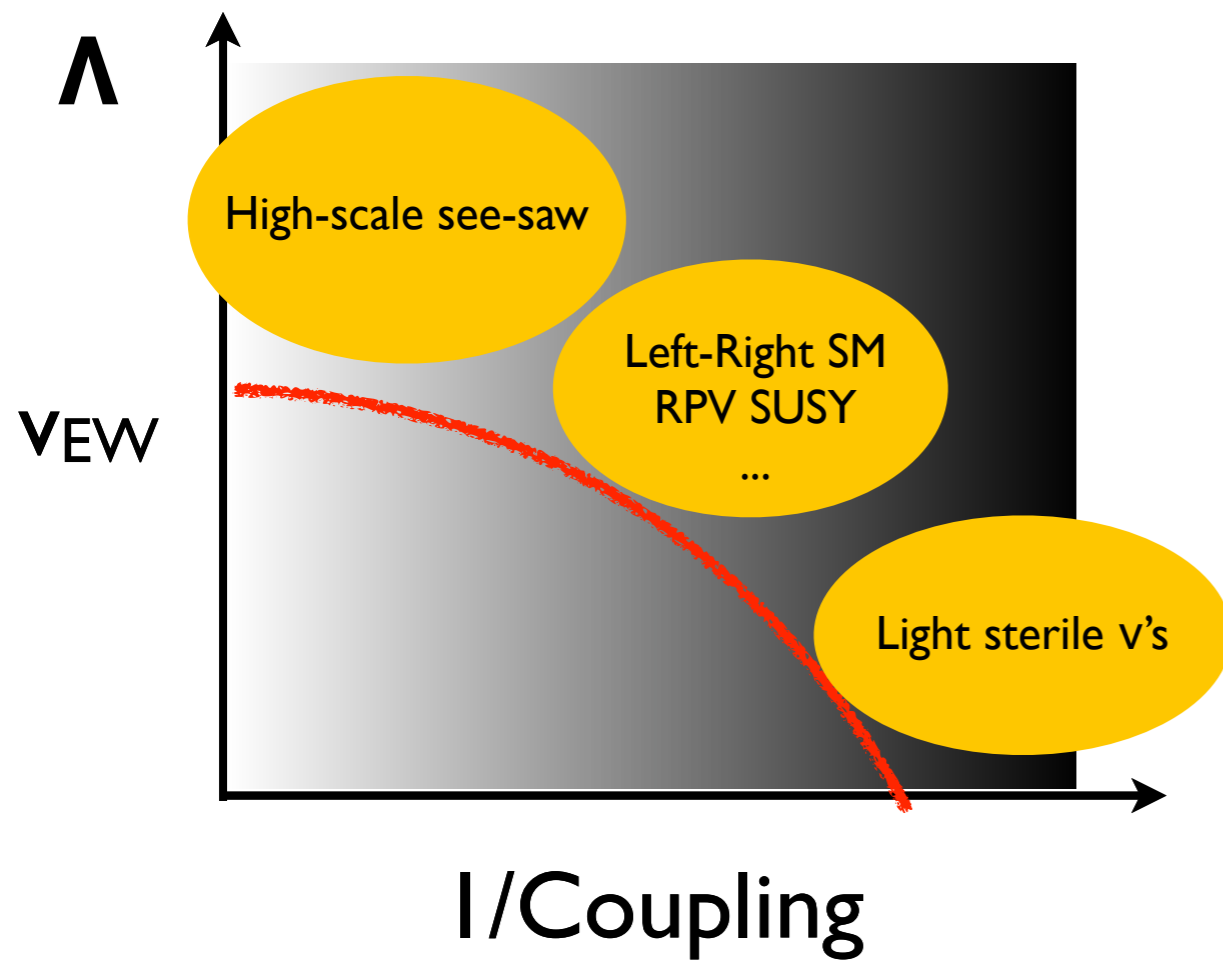
The quest is on...

- Several international “**ton-scale**” experiments with different isotopes and technologies under way, with sensitivity up to $T_{1/2} \sim 10^{28}$ yr



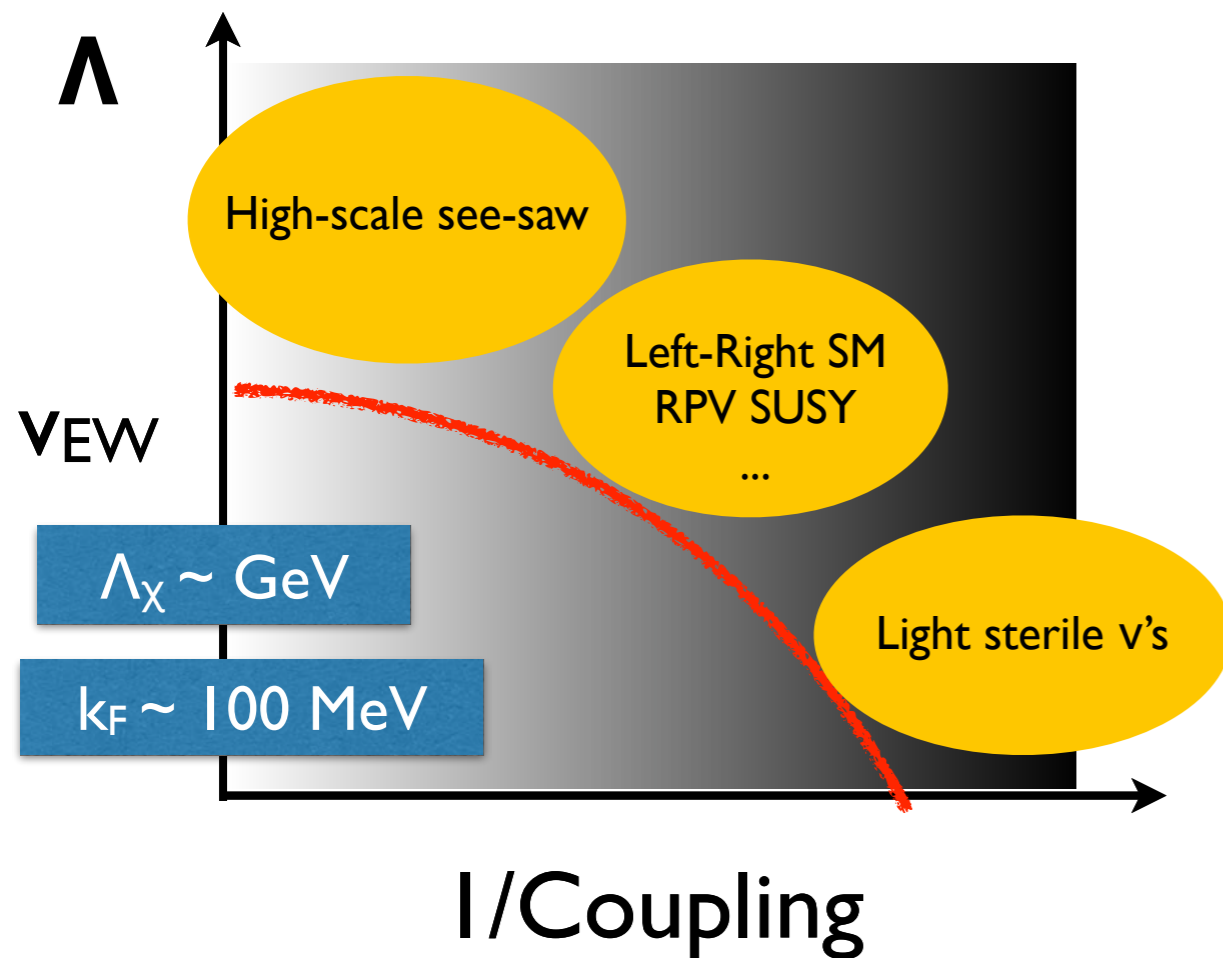
$0\nu\beta\beta$ physics reach

- $0\nu\beta\beta$ searches @ $T_{1/2} > 10^{27-28}$ yr will have broad sensitivity to LNV mechanisms



$0\nu\beta\beta$ physics reach

- $0\nu\beta\beta$ searches @ $T_{1/2} > 10^{27-28}$ yr will have broad sensitivity to LNV mechanisms



- Multi-scale problem best tackled through 'end-to-end' EFT: only chance to achieve controllable uncertainty
- Synergy of **EFT**, **Lattice QCD**, and first-principles **nuclear structure**

SMEFT

LEFT

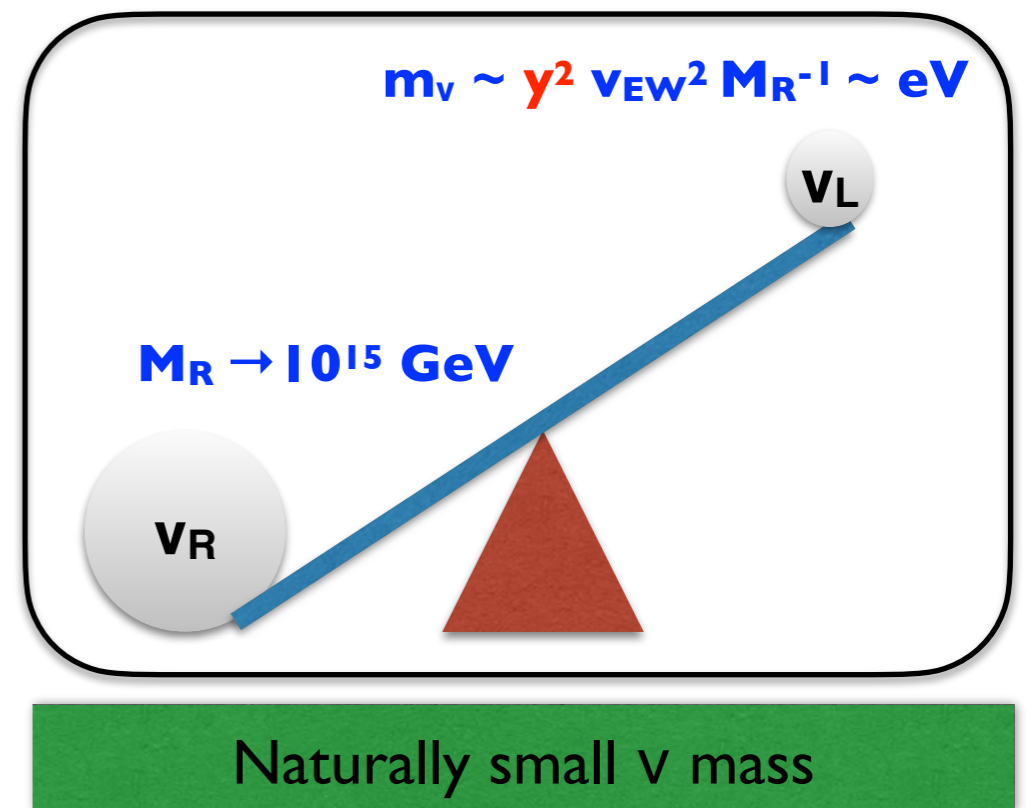
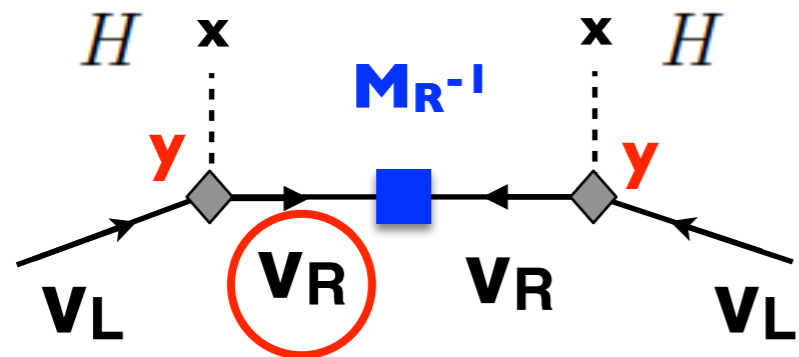
Chiral EFT

$$T_{1/2} \propto (m_W/\Lambda)^A (\Lambda_\chi/m_W)^B (k_F/\Lambda_\chi)^C$$

Snowmass white paper 2203.21169 and refs therein

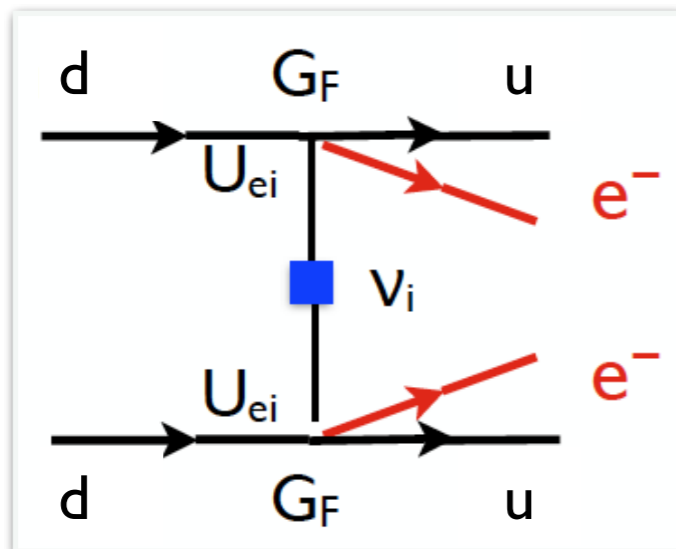
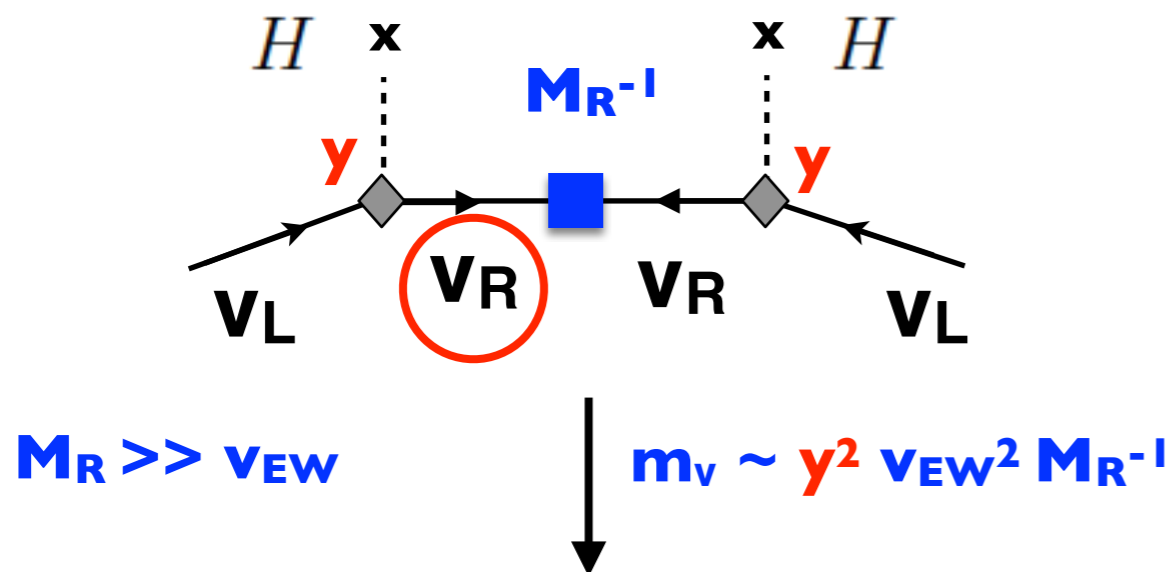
High-scale seesaw

- Majorana mass generated by exchange of heavy particles, such as heavy neutrinos that are neutral under all SM charges (=sterile)



High-scale seesaw

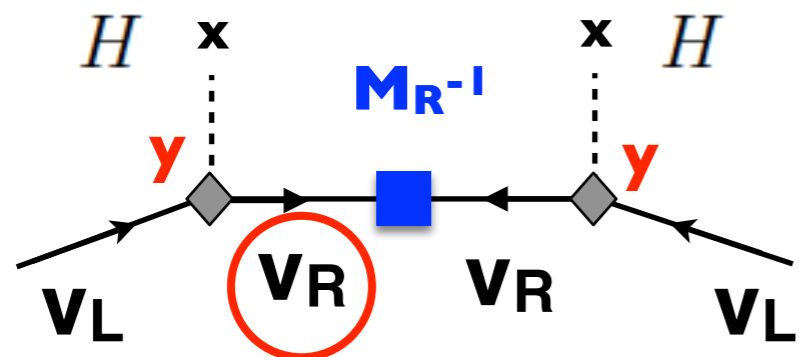
- Majorana mass generated by exchange of heavy particles, such as heavy neutrinos that are neutral under all SM charges (=sterile)



$0\nu\beta\beta$ mediated by light neutrinos

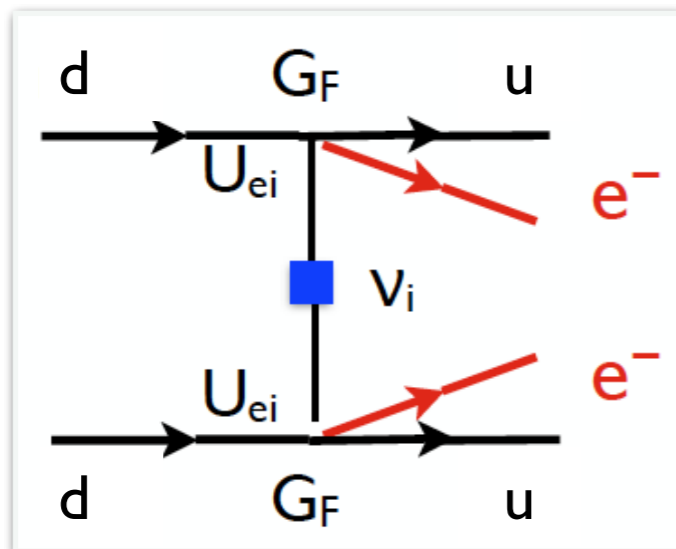
High-scale seesaw

- Majorana mass generated by exchange of heavy particles, such as heavy neutrinos that are neutral under all SM charges (=sterile)



$$M_R \gg v_{EW}$$

$$m_\nu \sim y^2 v_{EW}^2 M_R^{-1}$$



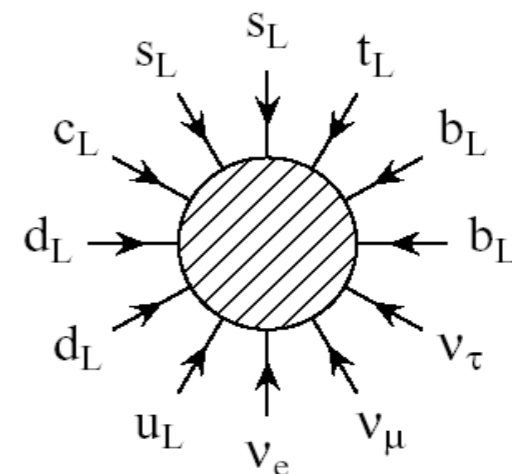
$0\nu\beta\beta$ mediated by light neutrinos

Baryogenesis via Leptogenesis

- CP- and L- violating out-of-equilibrium decays of heavy $\nu_{Ri} \Rightarrow n_L$

$$\Gamma(\nu_R \rightarrow H^* \ell) \neq \Gamma(\nu_R \rightarrow H \bar{\ell})$$

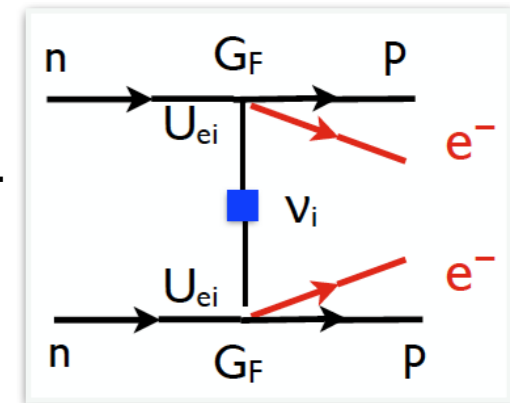
- EW sphalerons $\Rightarrow n_B = \# n_L$



Discovery potential / target

- $0\nu\beta\beta$ can be predicted in terms of ν mass parameters: $\Gamma \propto |M_{0\nu}|^2 (m_{\beta\beta})^2$

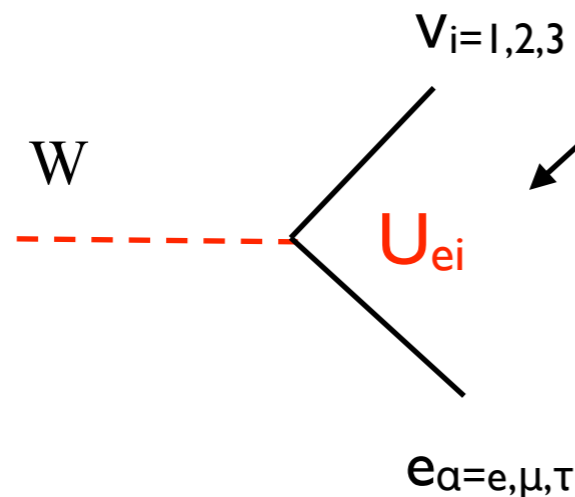
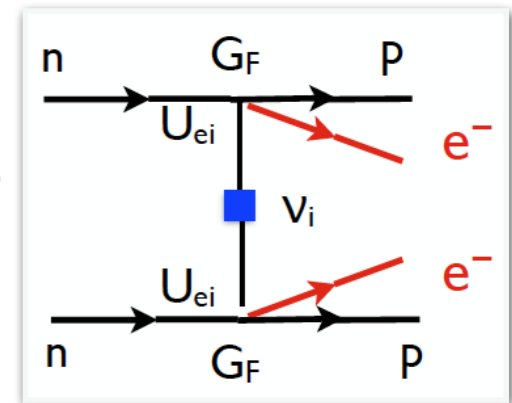
$$\langle m_{\beta\beta} \rangle^2 = \left| \sum_i U_{ei}^2 m_{\nu i} \right|^2$$



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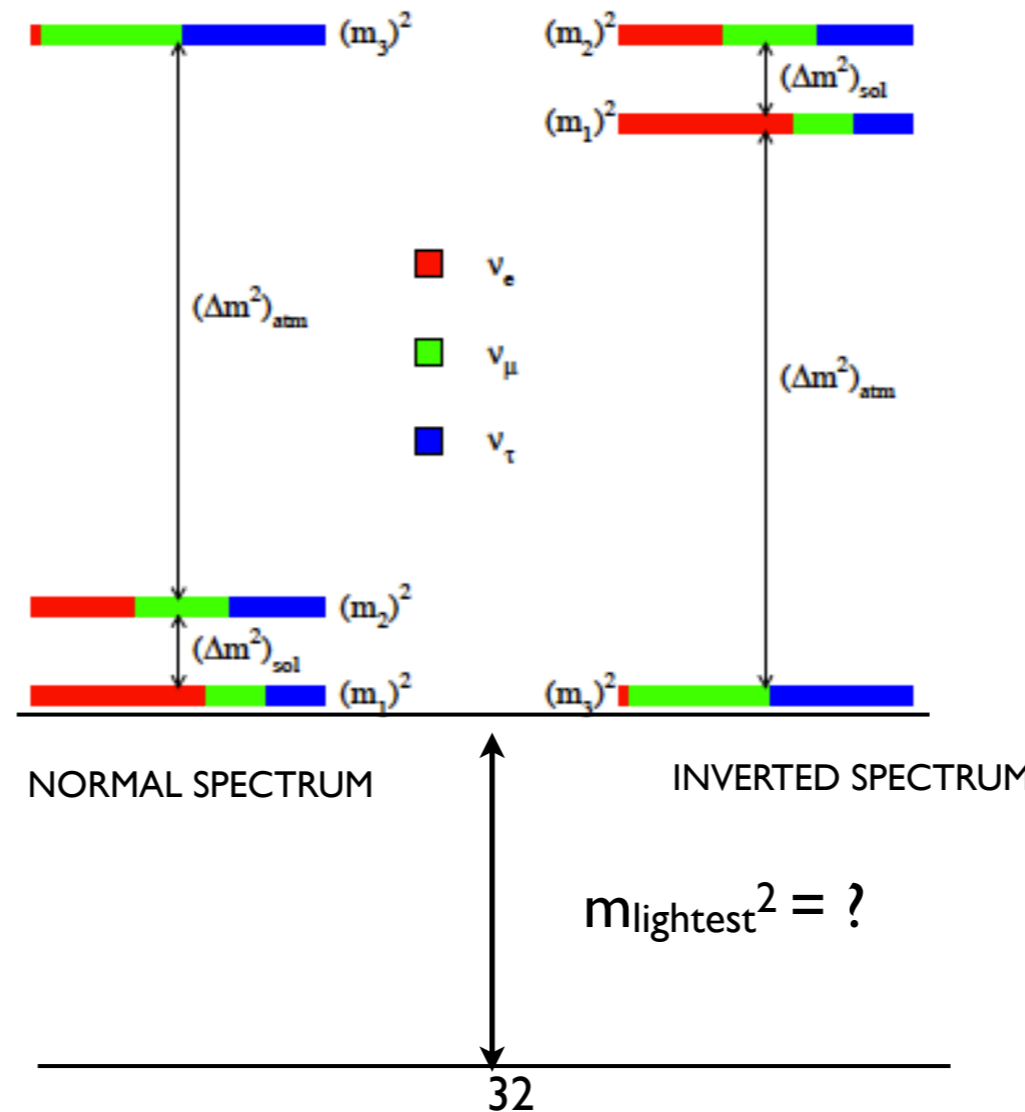
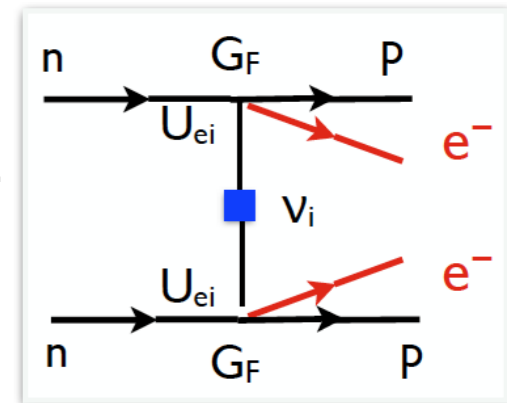
$$\frac{g}{\sqrt{2}} W_{\mu}^{-} \bar{e}_{L}^{\alpha} \gamma^{\mu} U^{\alpha i} \nu_{L}^{i}$$

Unitary mixing in CC vertex
(Pontecorvo-Maki-Nakagawa-Sakata):
3 angles (known), **1+2** phases (unknown)

Discovery potential / target

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$$\langle m_{\beta\beta} \rangle^2 = \left| \sum_i U_{ei}^2 m_{\nu i} \right|^2$$

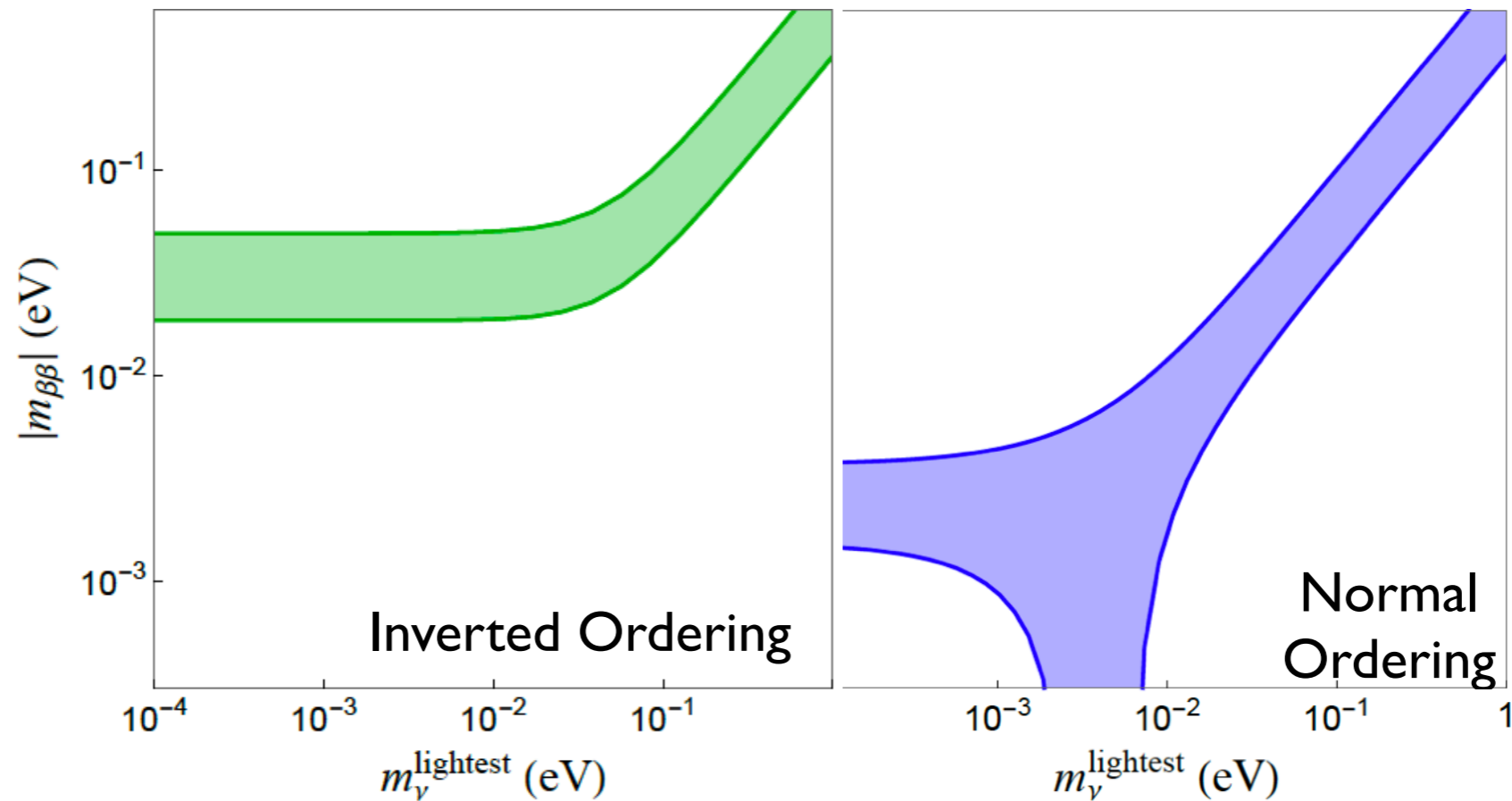
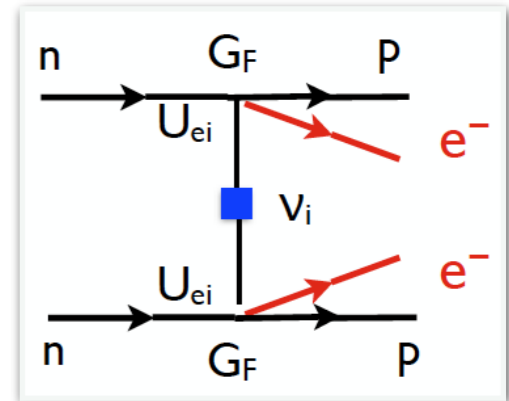


Mass ordering still
not fixed by
oscillation data

Discovery potential / target

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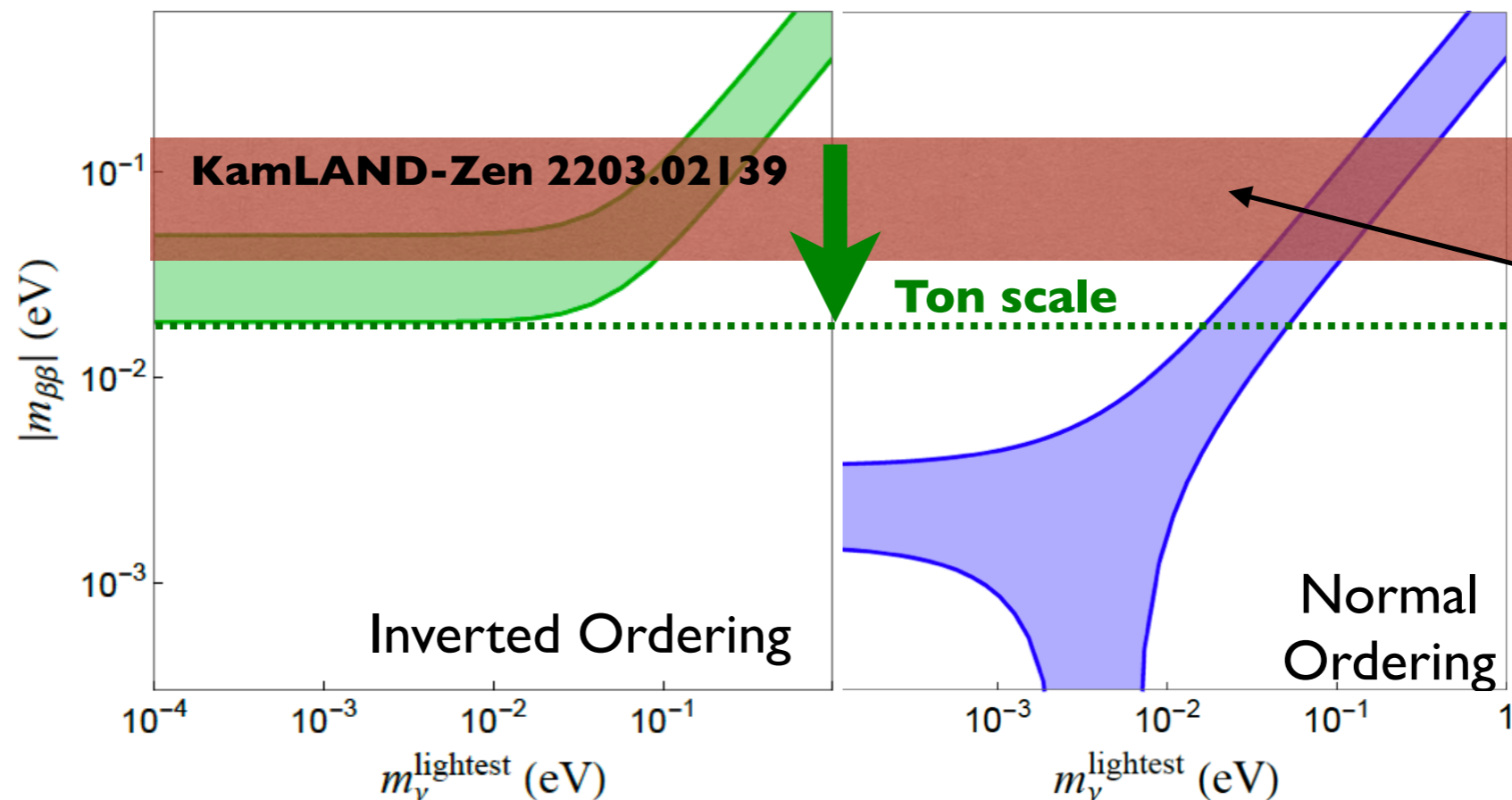
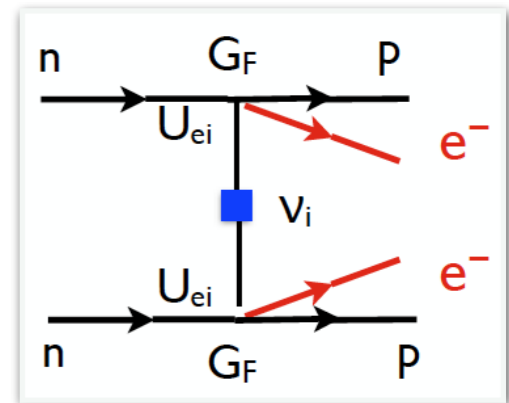


Bands: unknown
Majorana phases

Discovery potential / target

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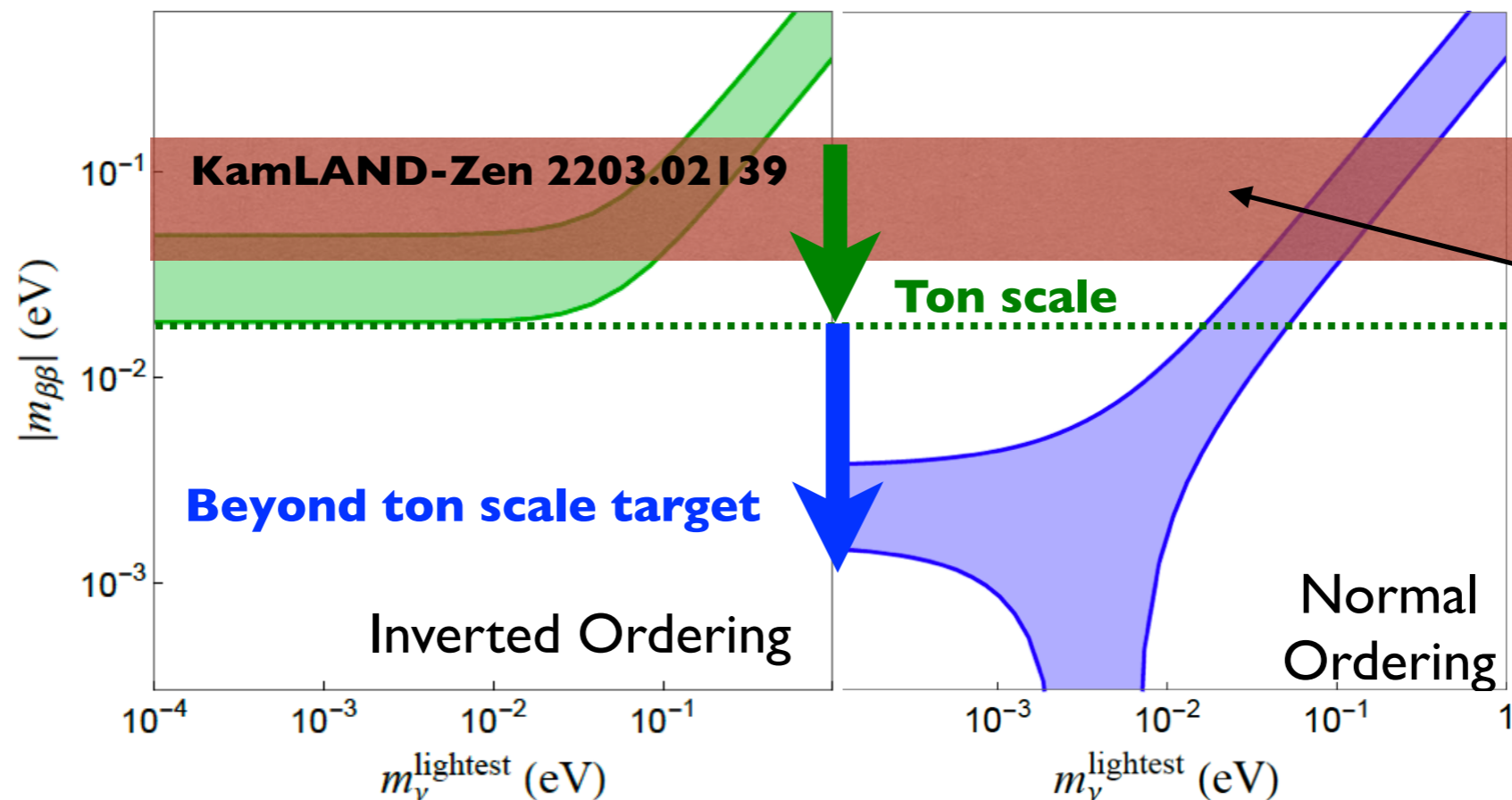
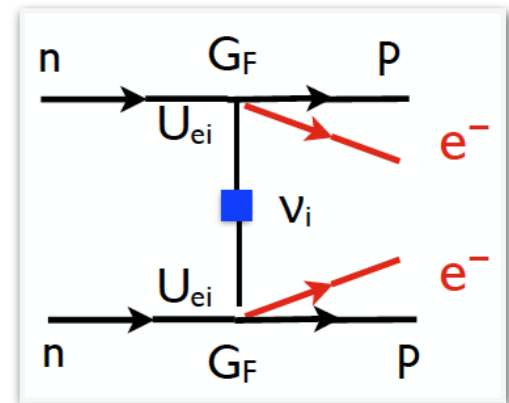
Bands: unknown Majorana phases

Assuming current range for matrix elements, discovery @ ton-scale possible for **inverted spectrum** or **$m_{\text{lightest}} > 50 \text{ meV}$**

Discovery potential / target

- $0\nu\beta\beta$ can be predicted in terms of ν mass parameters: $\Gamma \propto |M_{0\nu}|^2 (m_{\beta\beta})^2$

$$\langle m_{\beta\beta} \rangle^2 = \left| \sum_i U_{ei}^2 m_{\nu i} \right|^2$$



Natural (but challenging!) beyond ton-scale target is $m_{\beta\beta} \sim \text{meV}$

Diagnosing power

- High scale seesaw implies falsifiable correlations with other ν mass probes. Future data can unravel new LNV sources or physics beyond “ Λ CDM + m_ν ”

$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_i \right|$$

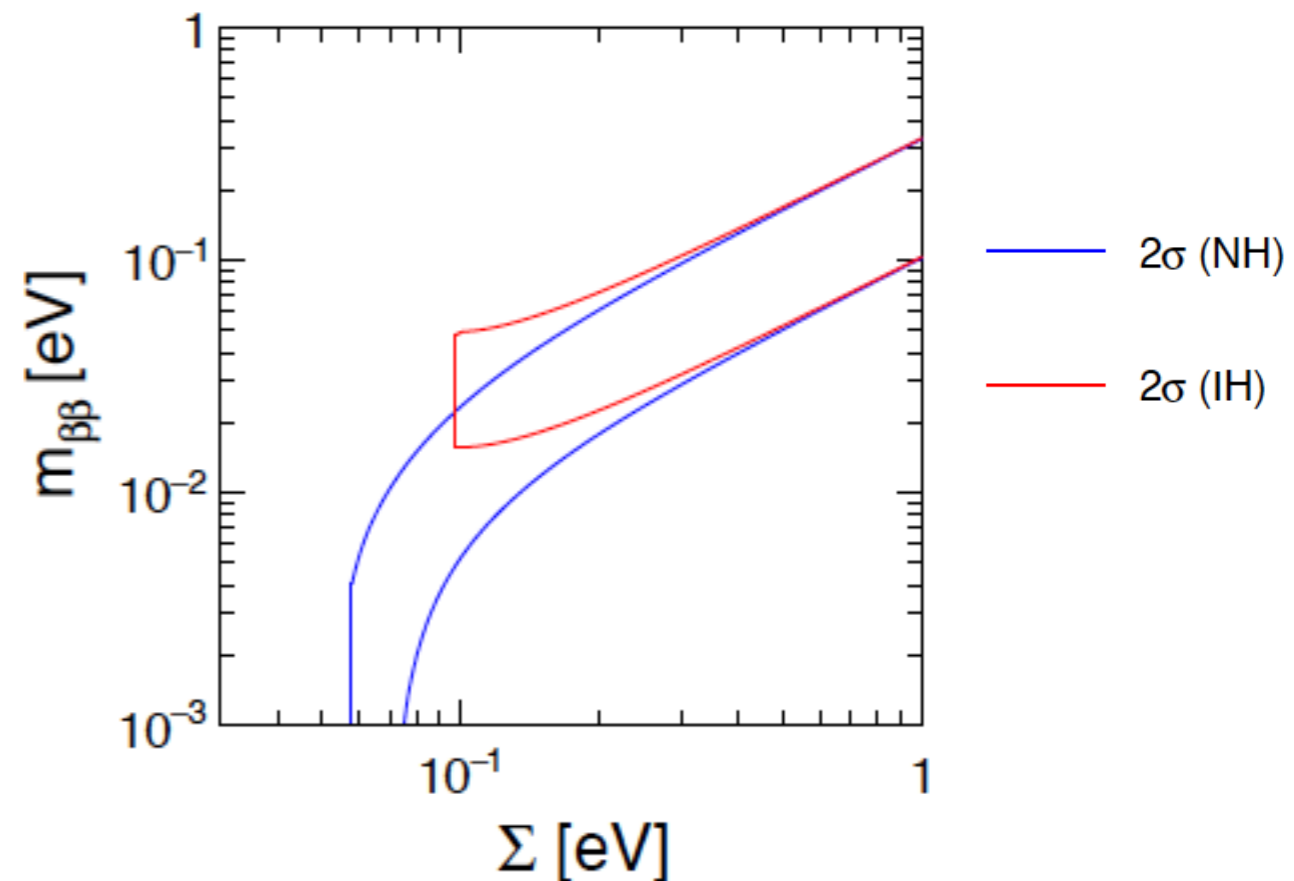
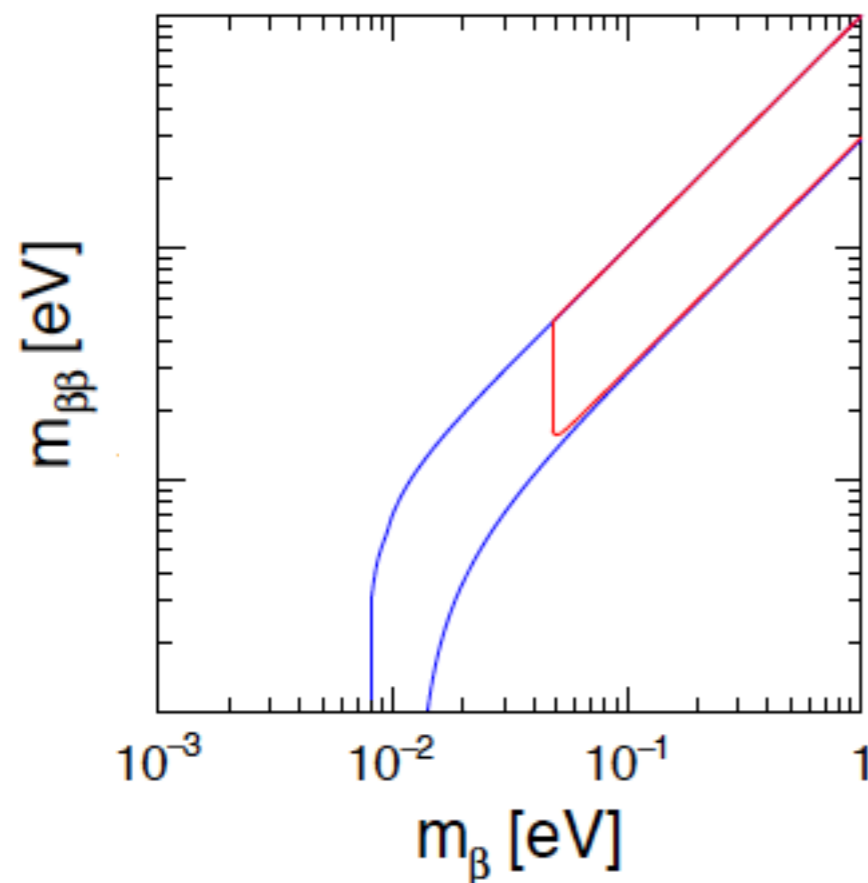
$0\nu\beta\beta$ decay

$$m_\beta = \sqrt{\sum_i |U_{ei}|^2 m_i^2}$$

Tritium β decay

$$\Sigma = \sum_i m_i$$

Cosmology



Diagnosing power

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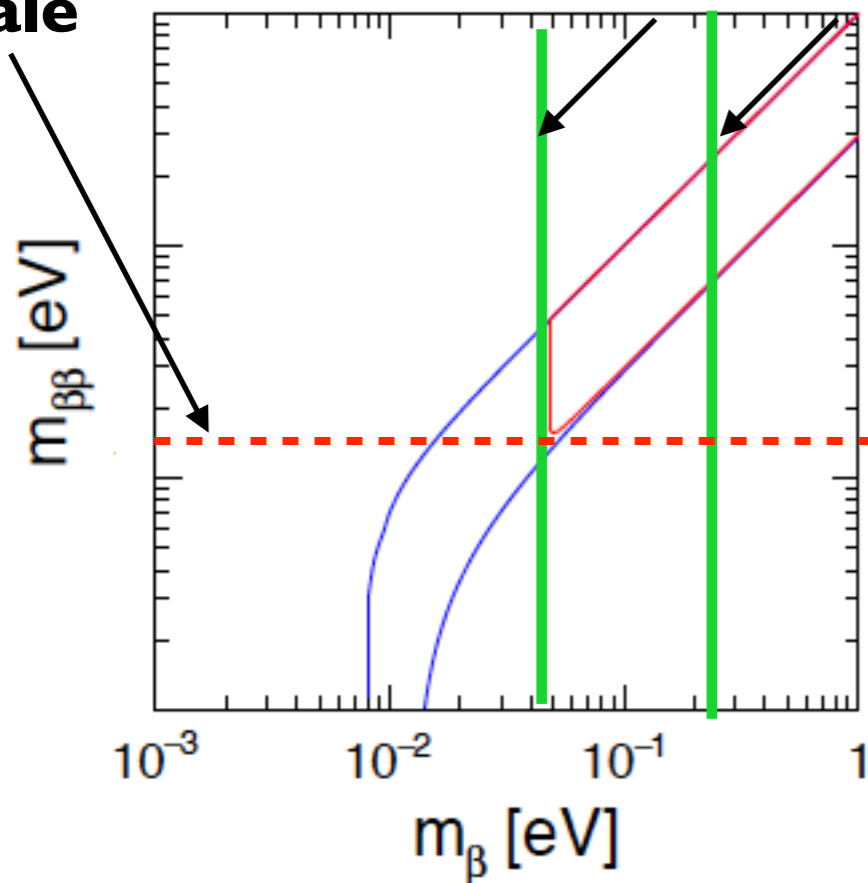
Tritium β decay

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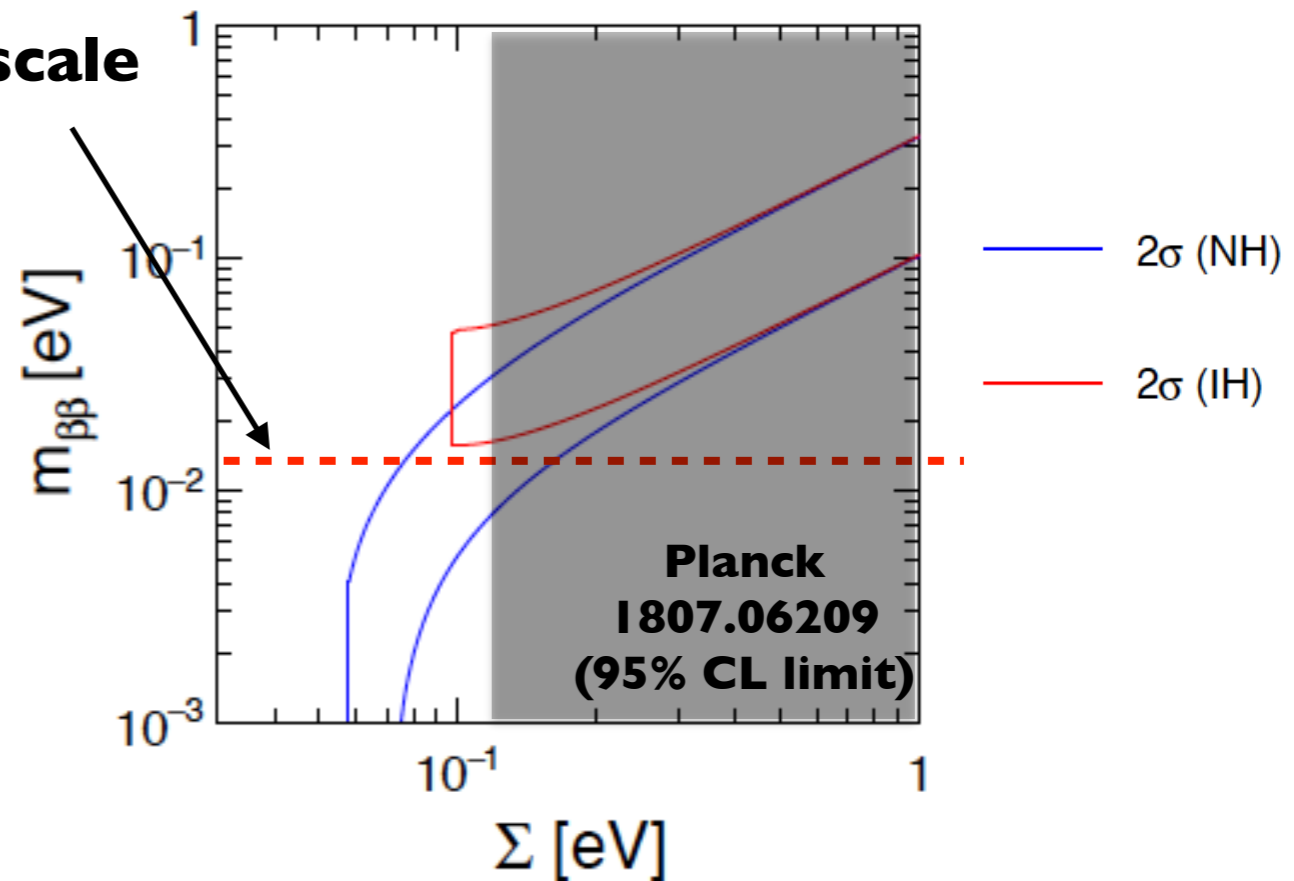
Cosmology

Project8 **KATRIN**

Ton scale



Ton scale



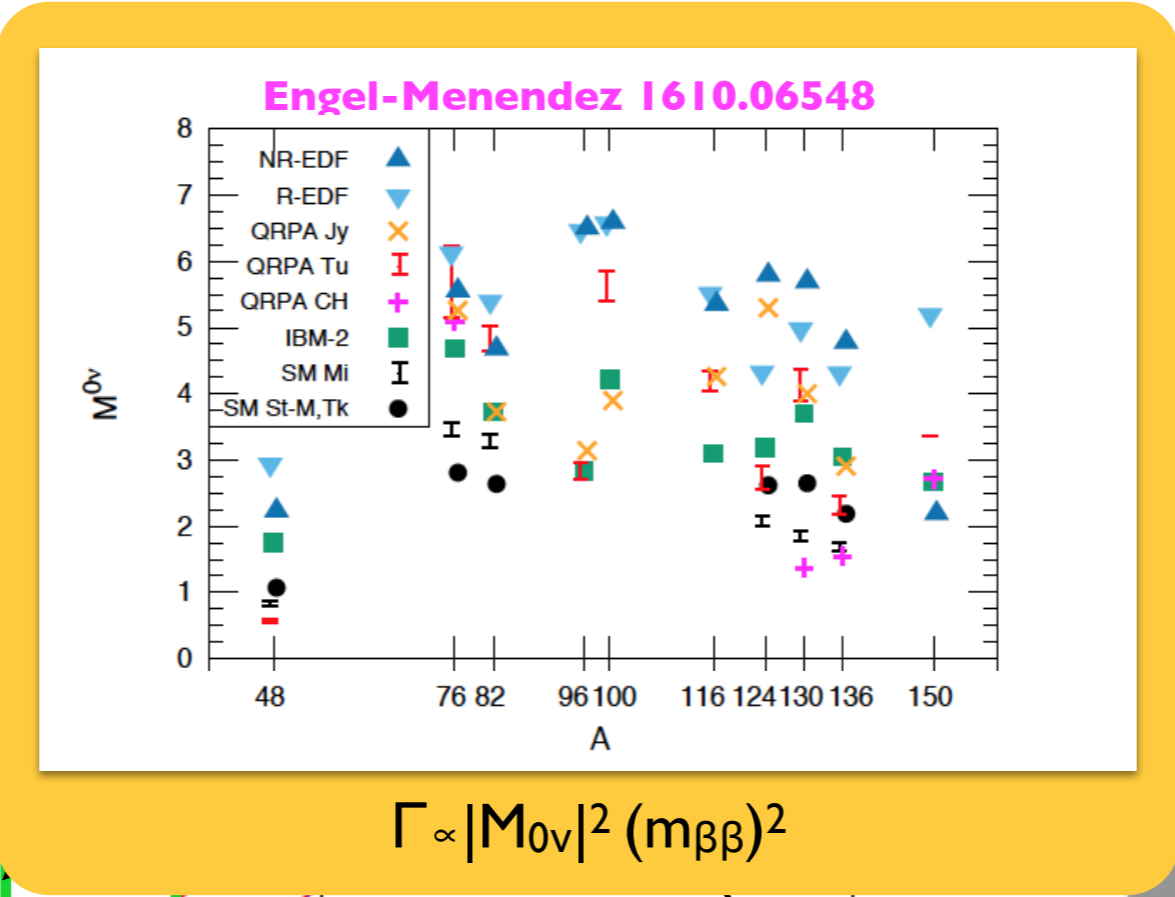
Diagnosing power

- High scale seesaw
- Future data can use

er ν mass probes.
 ond “ Λ CDM + m_ν ”

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$0\nu\beta\beta$ decay

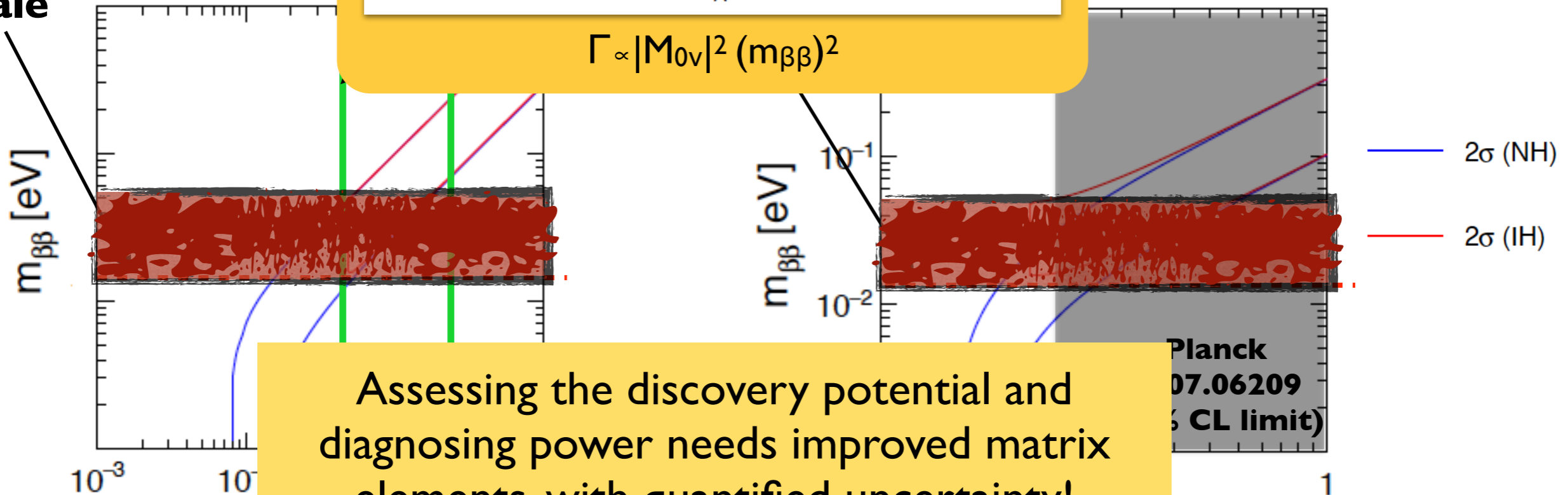


$$\Sigma = \sum_i m_i$$

Cosmology

Pr

Ton scale



Assessing the discovery potential and diagnosing power needs improved matrix elements, with quantified uncertainty!

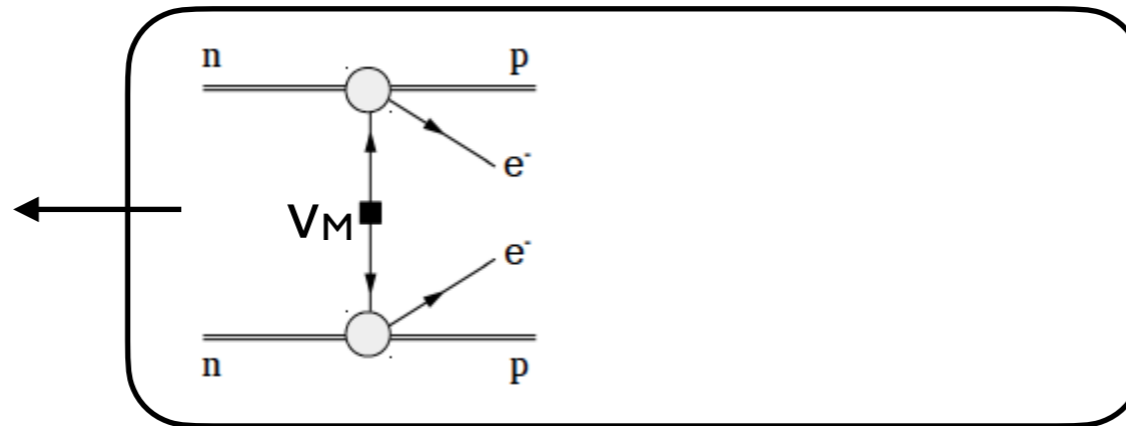
New insights from EFT

VC, W. Dekens, E. Mereghetti, A. Walker-Loud, 1710.01729

VC, W. Dekens, J. de Vries, M. Graesser, E. Mereghetti, S. Pastore, U. van Kolck 1802.10097

- Transition operator to leading order in Q/Λ_χ ($Q \sim k_F \sim m_\pi$, $\Lambda_\chi \sim \text{GeV}$)

'Usual' V_M exchange
 $\sim 1/k_F^2 \sim 1/Q^2$
Coulomb-like potential



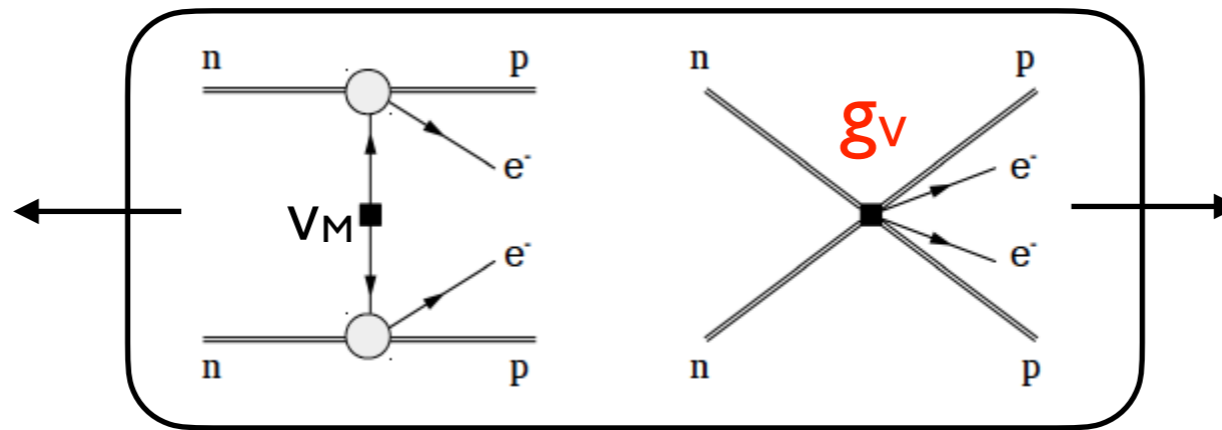
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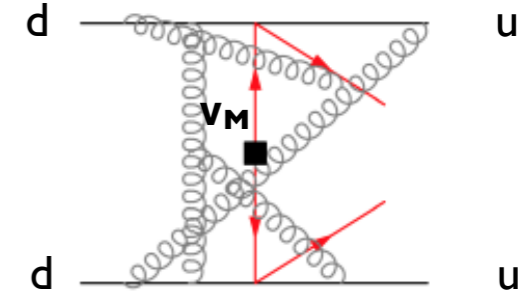
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'New': short-range
 coupling $g_V \sim 1/Q^2$



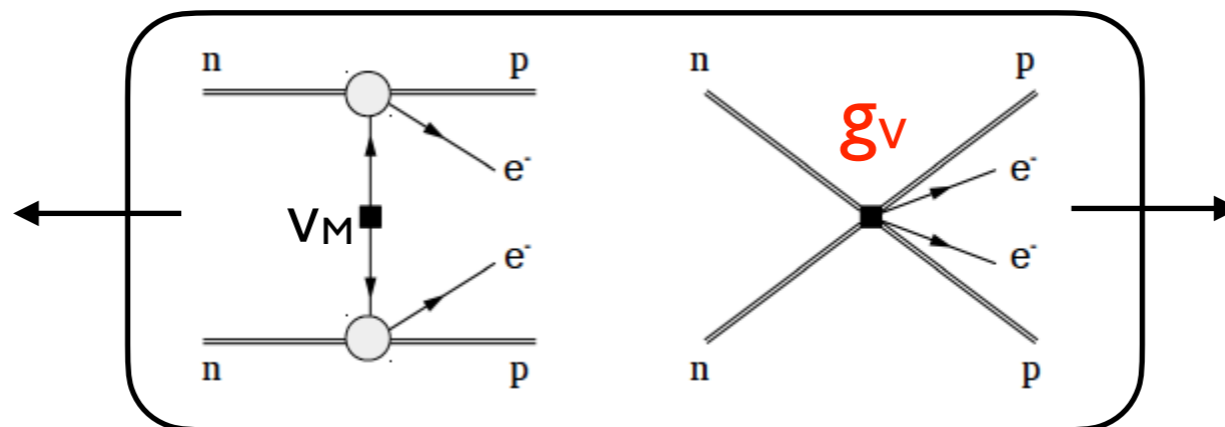
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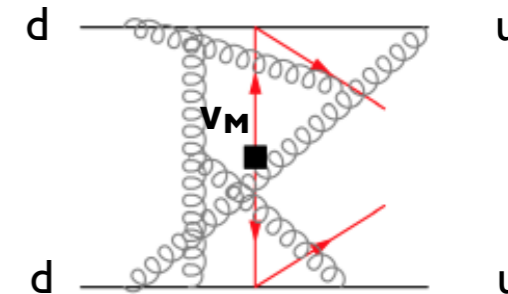
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'New': short-range coupling $g_V \sim 1/Q^2$



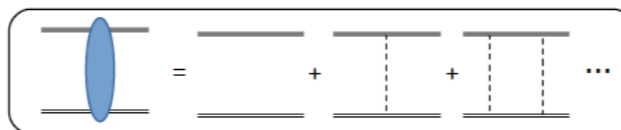
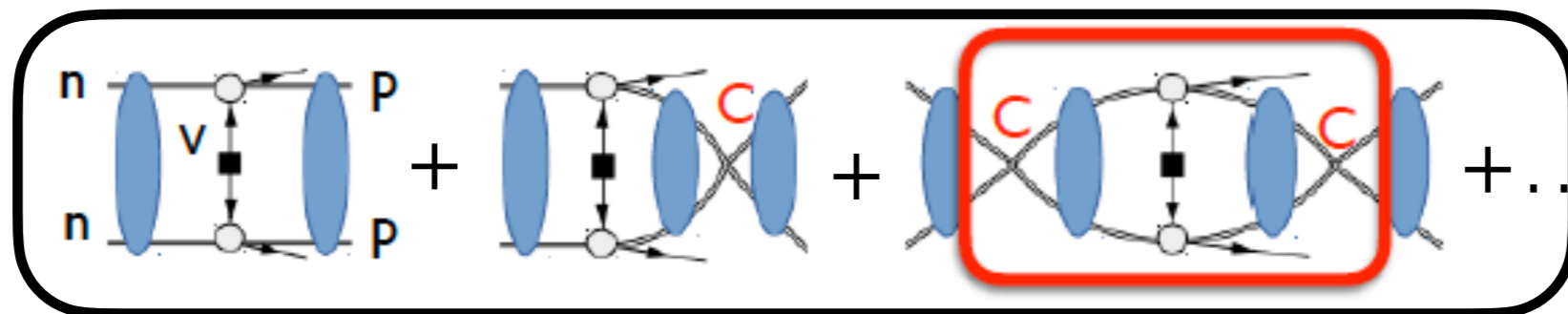
- Required by renormalization of $nn \rightarrow pp$ amplitude in presence of strong interactions

UV divergence $\propto (m_N C / 4\pi)^2 \sim 1/Q^2$

LO strong potential



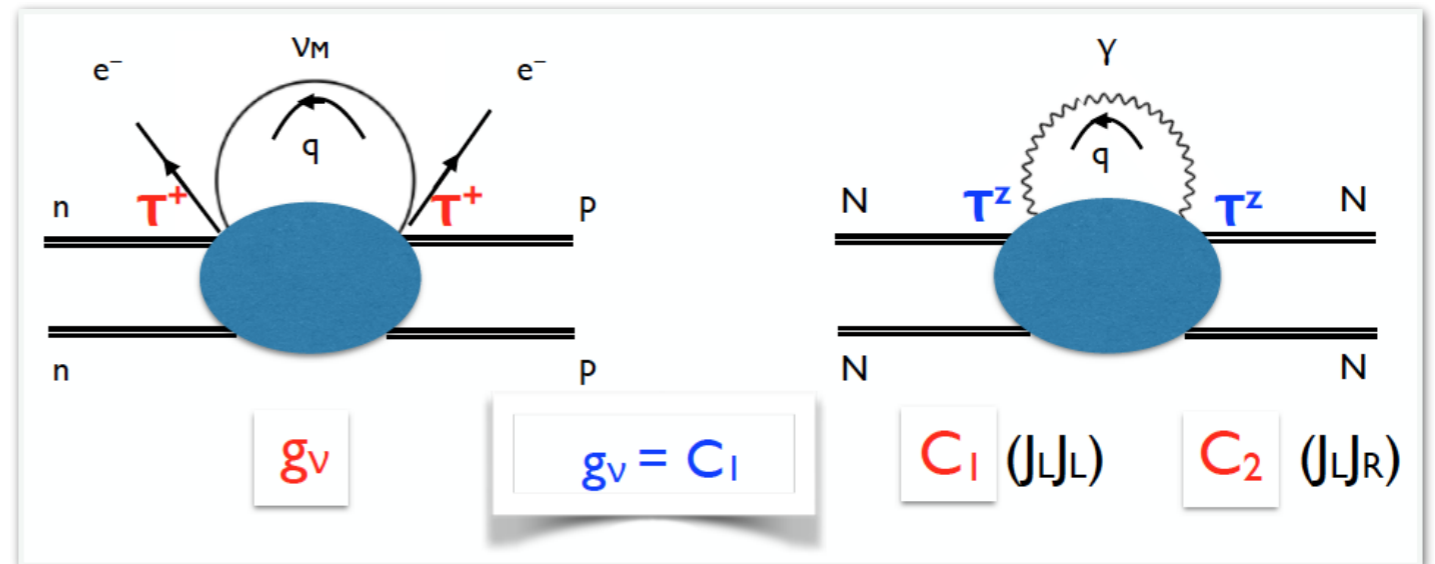
$C \sim 4\pi / (m_N Q)$



Impact on nuclear matrix elements

VC, W. Dekens, J. de Vries, M. Graesser, E. Mereghetti, S. Pastore, U. van Kolck 1802.10097

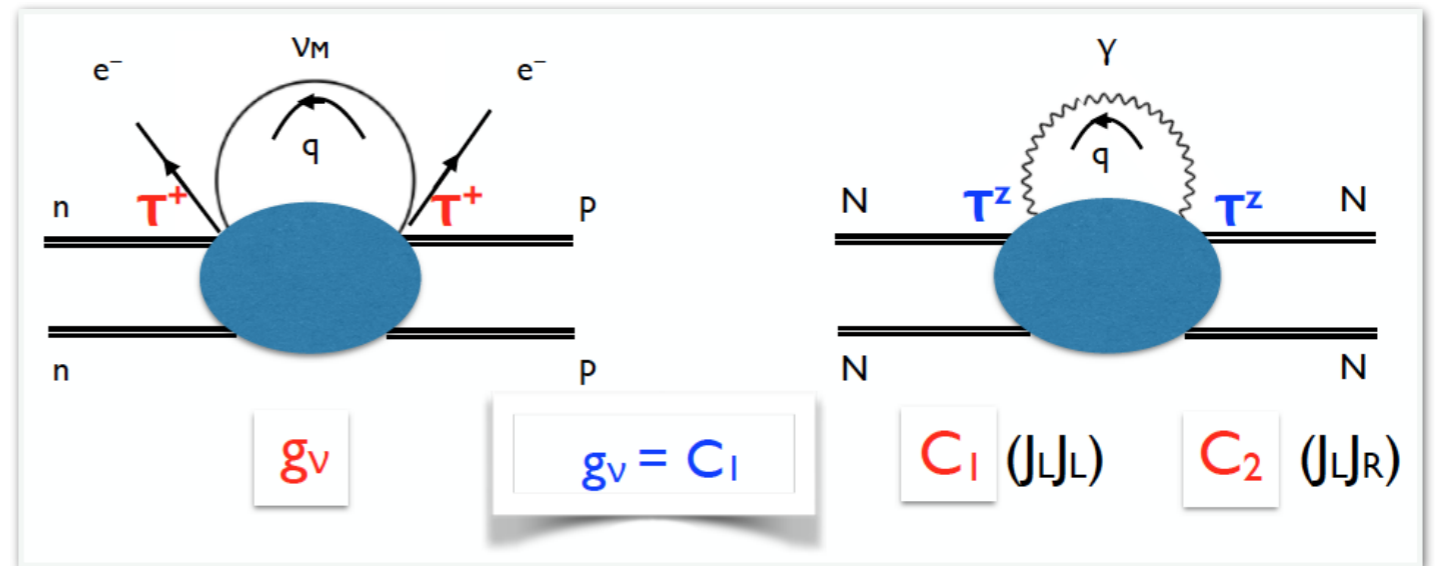
- Chiral+isospin symmetry relate g_V to one of two $I=2$ e.m. couplings (hard γ 's & v 's)
- NN data ($a_{nn}+a_{pp}-2a_{np}$) determine C_1+C_2 , confirming LO scaling



Impact on nuclear matrix elements

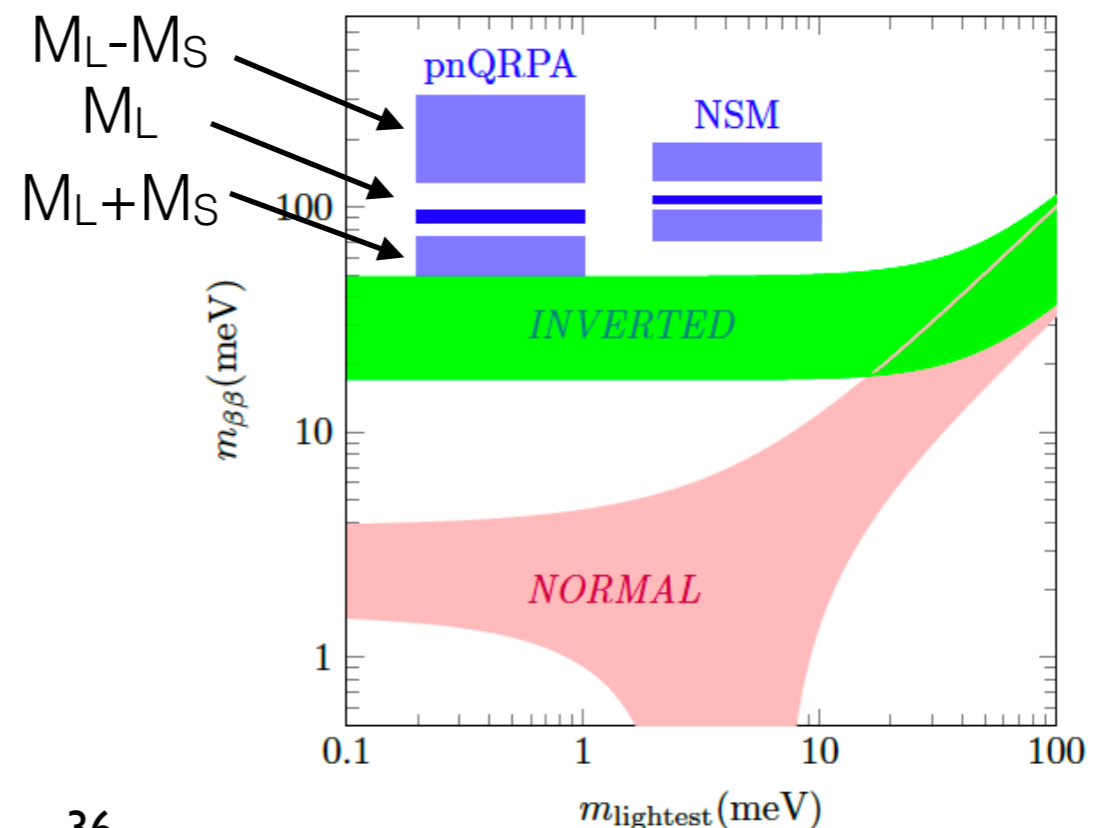
VC, W. Dekens, J. de Vries, M. Graesser, E. Mereghetti, S. Pastore, U. van Kolck 1802.10097

- Chiral+isospin symmetry relate g_v to one of two $I=2$ e.m. couplings (hard γ 's & v 's)
- NN data ($a_{nn}+a_{pp}-2a_{np}$) determine C_1+C_2 , confirming LO scaling



- Assuming $g_v \sim (C_1 + C_2)/2 \rightarrow$ $O(1)$ impact on m.e. and $m_{\beta\beta}$ extraction

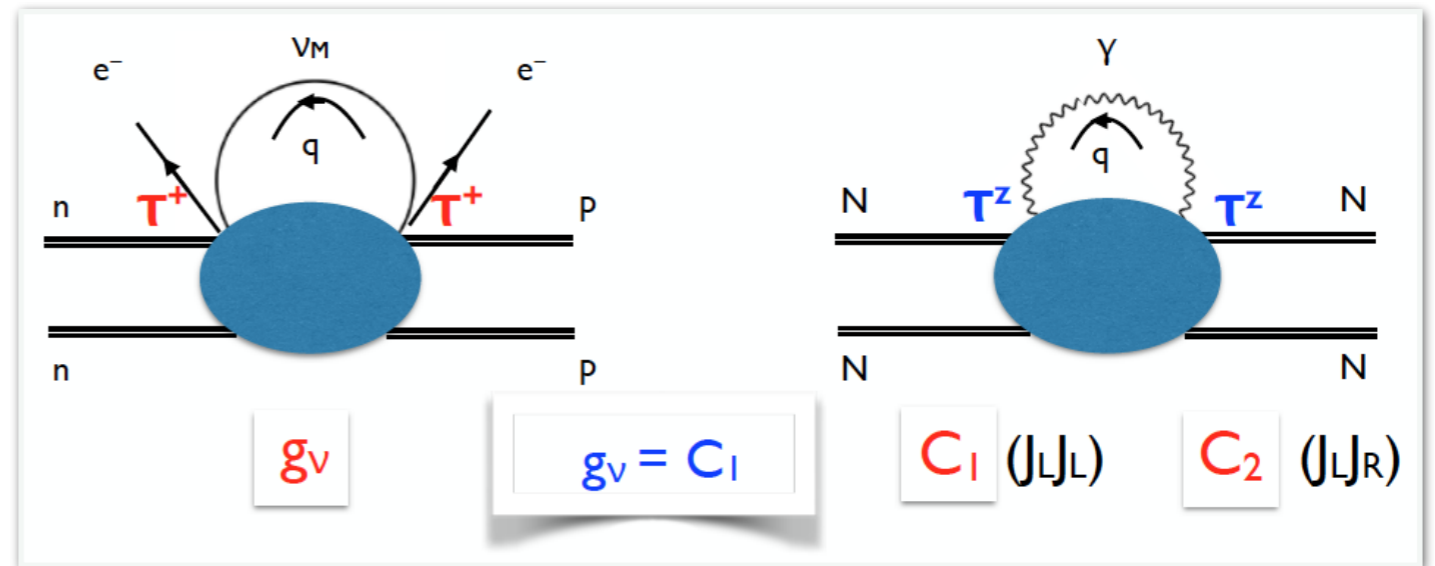
Iokiniemi-Soriano-Menendez, 2107.13354



Impact on nuclear matrix elements

VC, W. Dekens, J. de Vries, M. Graesser, E. Mereghetti, S. Pastore, U. van Kolck 1802.10097

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Iokiniemi-Soriano-Menendez, 2107.13354

$M_l - M_s$ vs m_{lightest}

Several approaches to determine g_v

- Large- N_c arguments point to $g_v \sim (C_1 + C_2)/2$

Richardson, Shindler, Pastore, Springer, 2102.02814

- Lattice QCD — gearing up

Tuo et al. 1909.13525;
Detmold, Murphy 2004.07404

Davoudi, Kadam, 2012.02083

- Dispersive approach

VC, Dekens, deVries, Hoferichter, Mereghetti, 2012.11602, 2102.03371

0.1 1 10 100

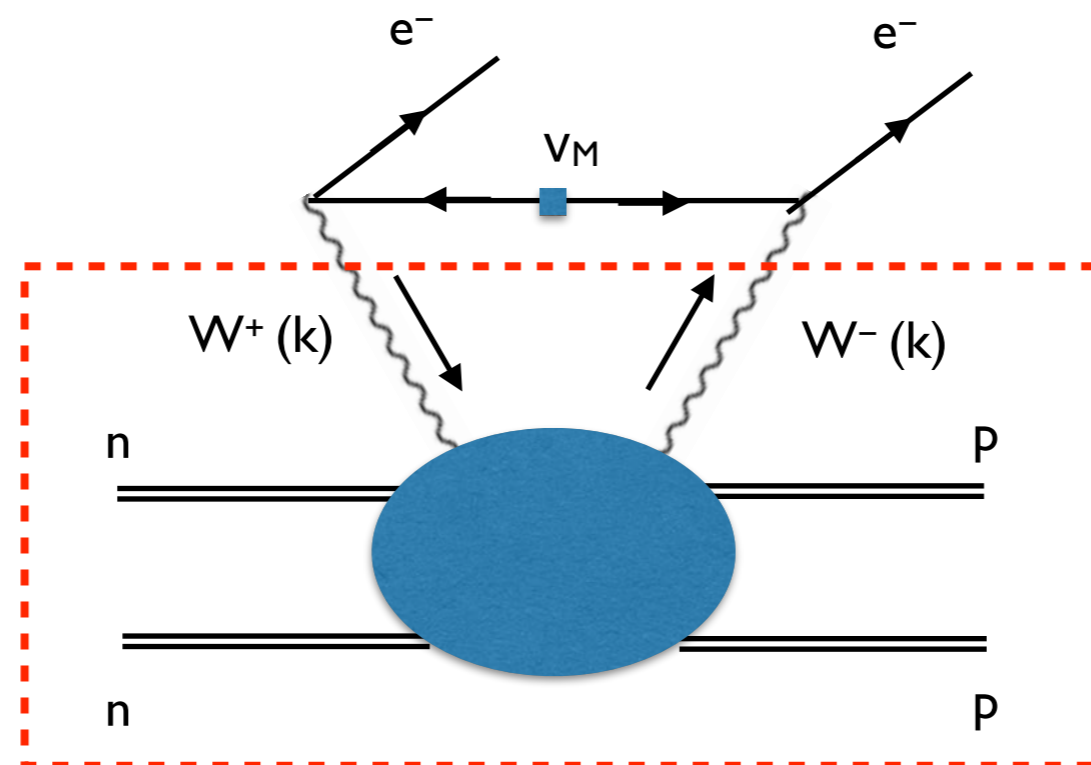
$m_{\text{lightest}}(\text{meV})$

Estimating the contact term (I)

VC, Dekens, deVries, Hoferichter, Mereghetti, 2012.11602, 2102.03371

- Useful representation of the amplitude

$$A_\nu \propto \int \frac{d^4 k}{(2\pi)^4} \frac{g_{\alpha\beta}}{k^2 + i\epsilon} \int d^4 x e^{ik \cdot x} \langle pp | T \{ j_W^\alpha(x) j_W^\beta(0) \} | nn \rangle$$



Forward “Compton” amplitude

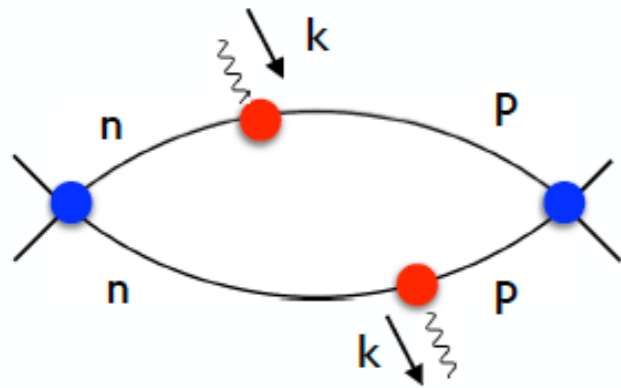
Estimating the contact term (I)

VC, Dekens, deVries, Hoferichter, Mereghetti, 2012.11602, 2102.03371

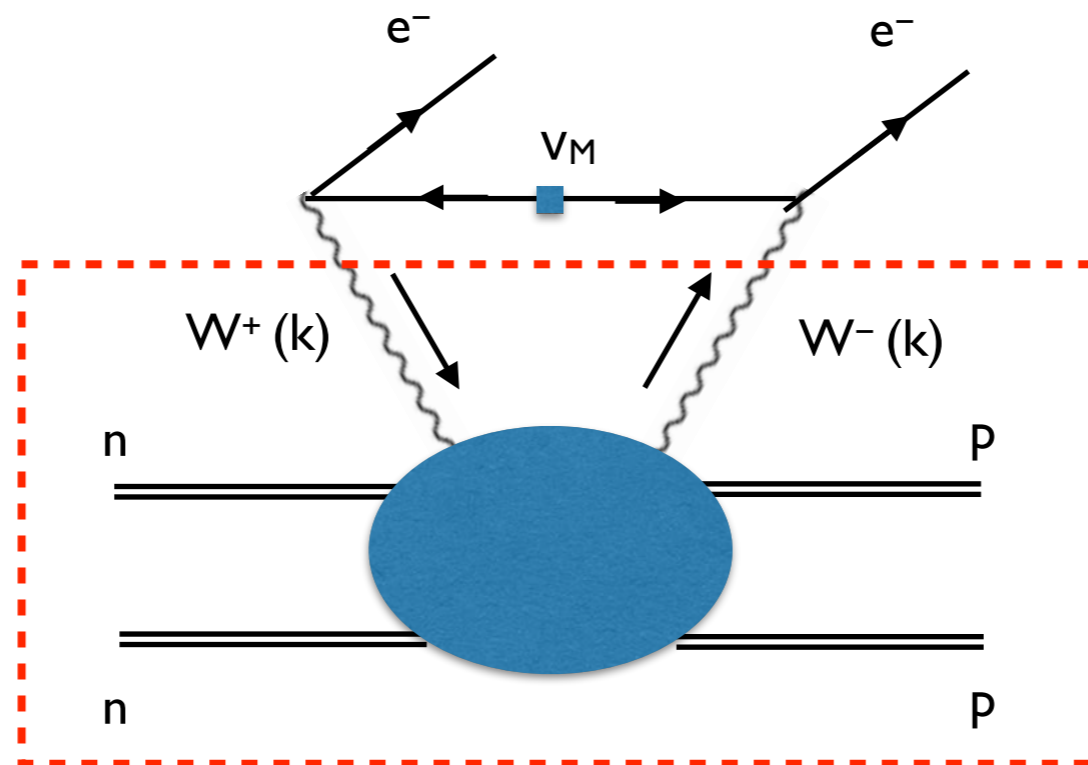
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Low k: chiral EFT to NLO

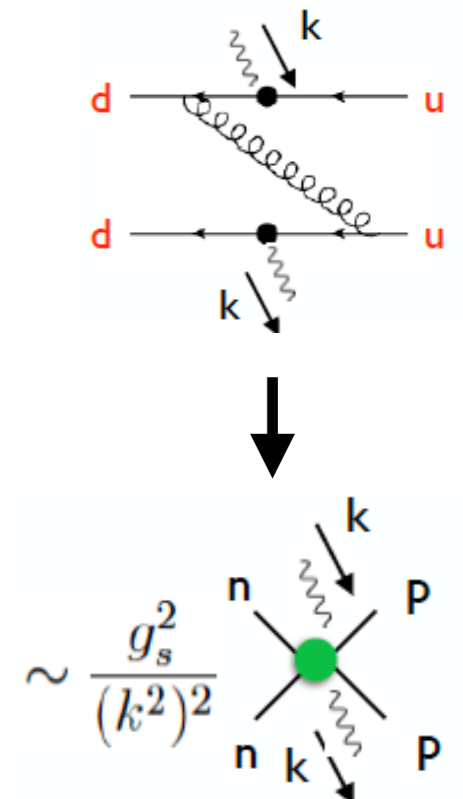


Intermediate k: resonance contributions in πNN intermediate state, ...



Forward "Compton" amplitude

High k: QCD OPE



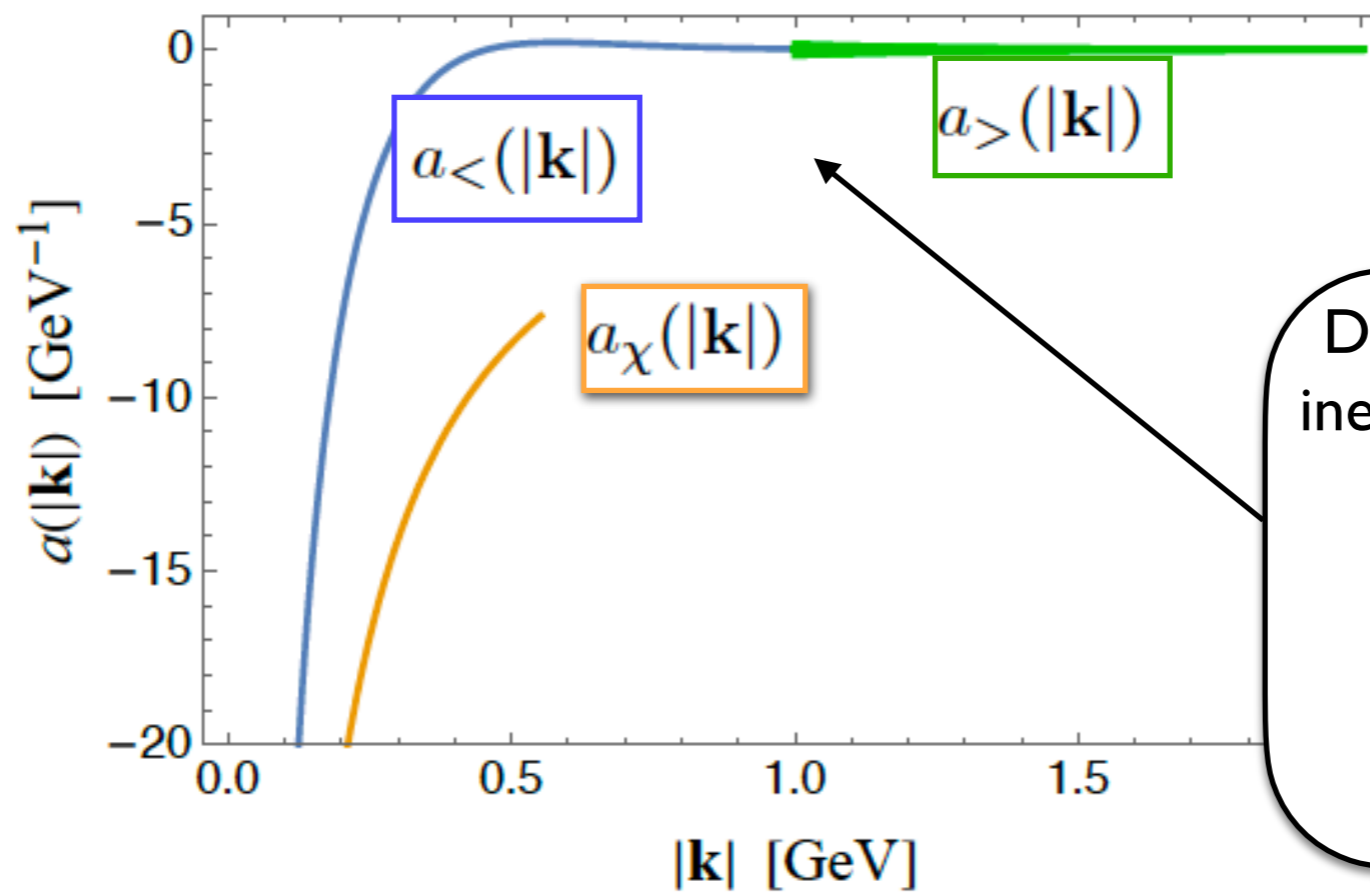
$$\sim \frac{g_s^2}{(k^2)^2}$$

Estimating the contact term (2)

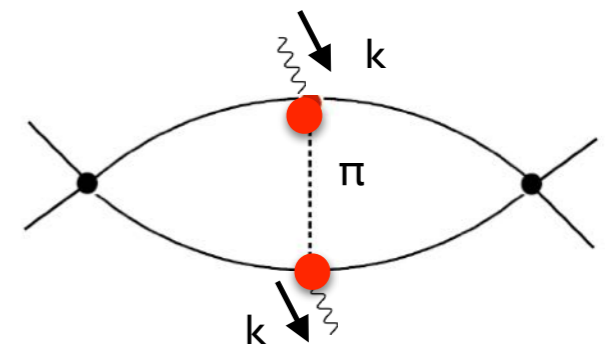
VC, Dekens, deVries, Hoferichter, Mereghetti, 2012.11602, 2102.03371

- Determine $C_{1,2}$ with $\sim 30\%$ uncertainty (dominated by intermediate k)

$$\mathcal{A}_\nu \propto \int_0^\Lambda d|\mathbf{k}| a_{<}(|\mathbf{k}|) + \int_\Lambda^\infty d|\mathbf{k}| a_{>}(|\mathbf{k}|)$$



Dominant uncertainty from inelastic channels ($NN\pi, \dots$):



Estimating the contact term (2)

VC, Dekens, deVries, Hoferichter, Mereghetti, 2012.11602, 2102.03371

- Determine $C_{1,2}$ with $\sim 30\%$ uncertainty (dominated by intermediate k)
- Validation: $C_1 + C_2 \Rightarrow (a_{nn} + a_{pp})/2 - a_{np} = 15.5(4.5)$ fm versus $10.4(2)$ fm (exp)
- Provided 'synthetic data' for the $nn \rightarrow pp$ amplitude at threshold
- First calculation of $^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$ with contact fitted to synthetic data \Rightarrow **contact term enhances nuclear matrix element by $(43 \pm 7)\%$**

Wirth, Yao, Hergert, 2105.05415

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Good news, while we wait for lattice results and first-principles calculations in heavier nuclei

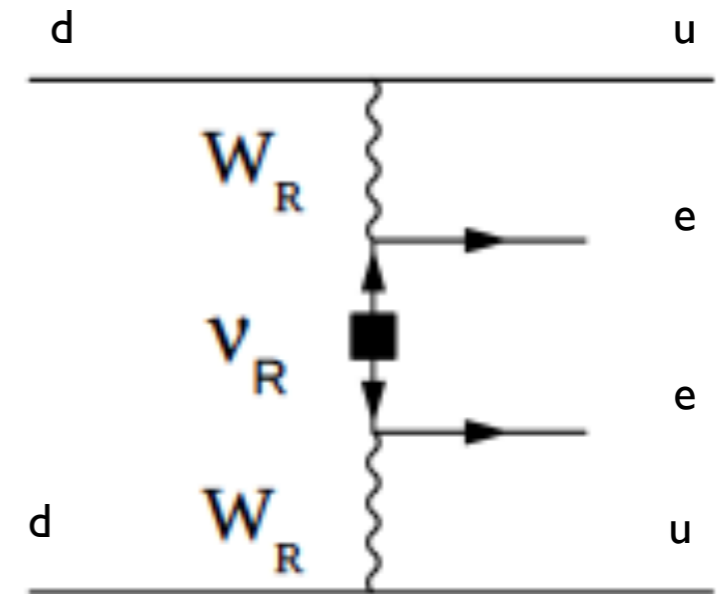
LNV from multi-TeV scale physics

- Observable contributions to $0\nu\beta\beta$ *not directly related to the exchange of light neutrinos:*

$$m_{\beta\beta} G_F^2 / Q^2 \sim 1 / \Lambda^5$$

(if $m_{\beta\beta} \sim 0.1$ eV and $\Lambda \sim \text{TeV}$)

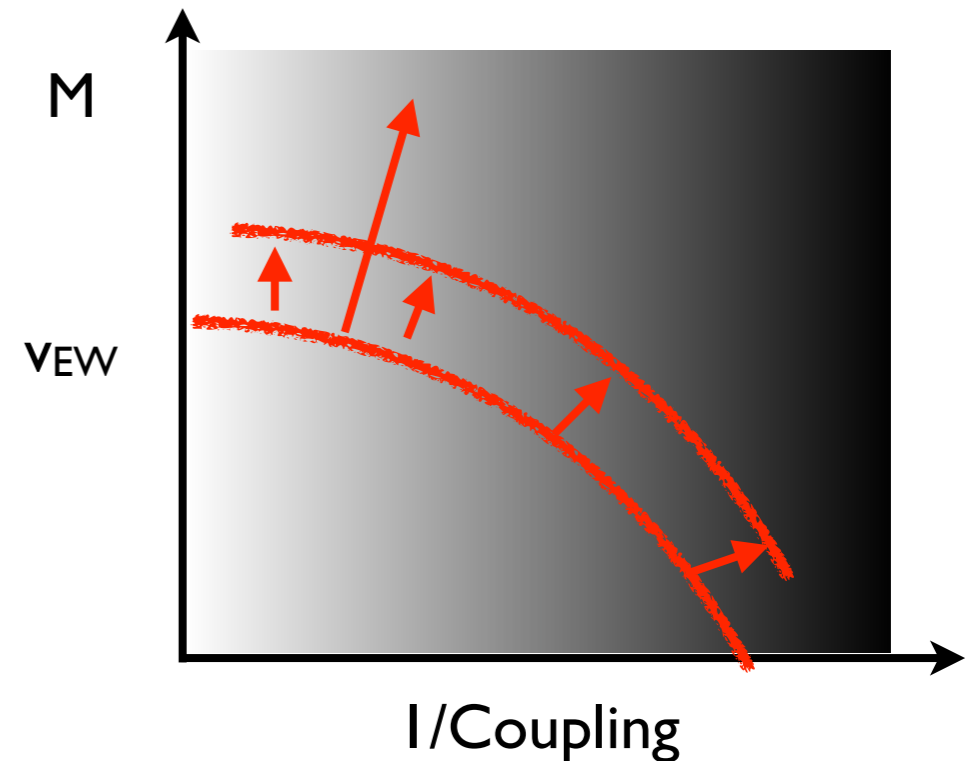
$$(\Lambda \sim M_{\nu R} \sim M_{W R})$$



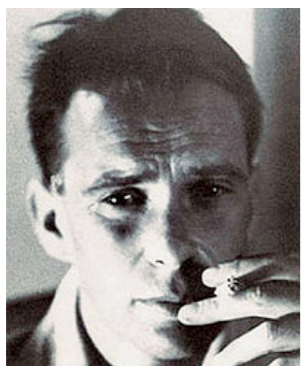
- New contributions *can interfere with* $m_{\beta\beta}$ or add incoherently, thus changing the interpretation of experimental results in terms of m_ν
- Correlated (or precursor!) signal at LHC: $pp \rightarrow ee jj$

Concluding comments

- Precision frontier experiments offer powerful ways to search for new physics
- Illustrated impact through two examples:
- β decays as a precision electroweak test
 - CKM unitarity test sensitive to new physics in multi-TeV range. Discovery window exists well into the LHC era
- $0\nu\beta\beta$ decay and lepton number violation
 - Ton-scale searches have great discovery potential — we simply don't know origin of neutrino mass and scale Λ associated with LNV
 - Theory progress through synergy of EFT, lattice QCD, and nuclear structure



Thank you!



A drawing by
Bruno Touschek