

Astrophysical constraints on the high-density equation of state

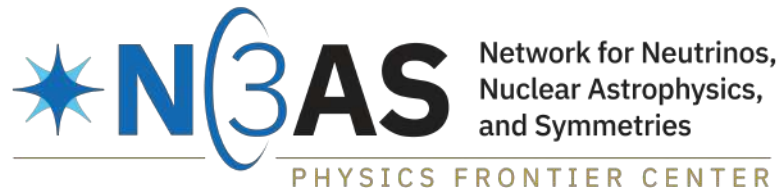
Sophia Han, TDLI Fellow

T.D. Lee Institute, Shanghai Jiao Tong Univ.



李政道研究所
TSUNG-DAO LEE INSTITUTE

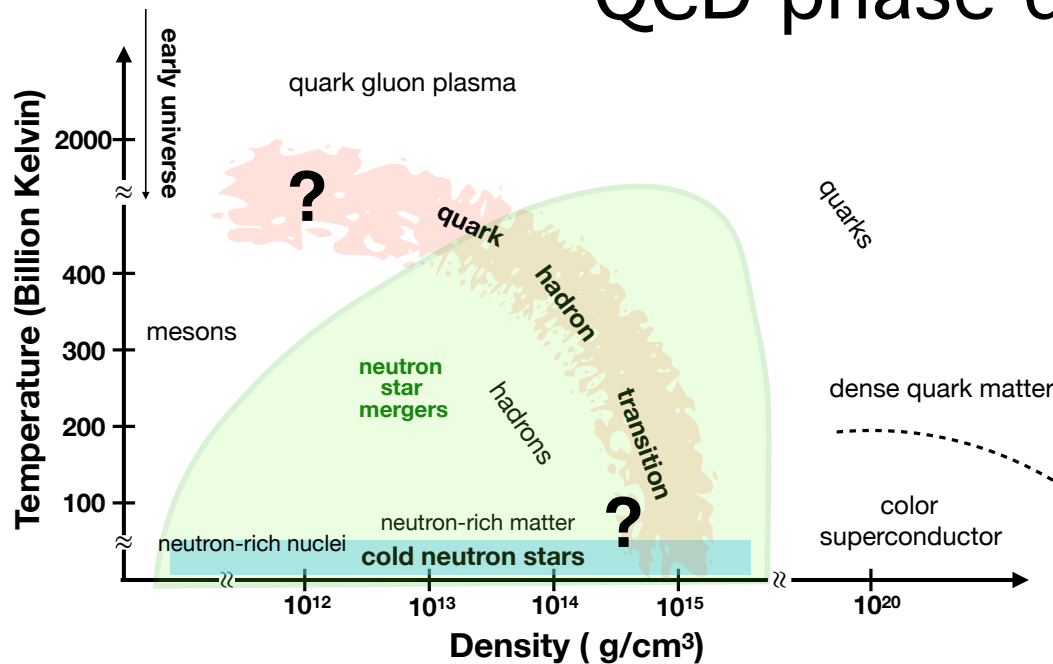
Theoretical Physics Colloquium
May 18, 2022 @ ASU



Outline

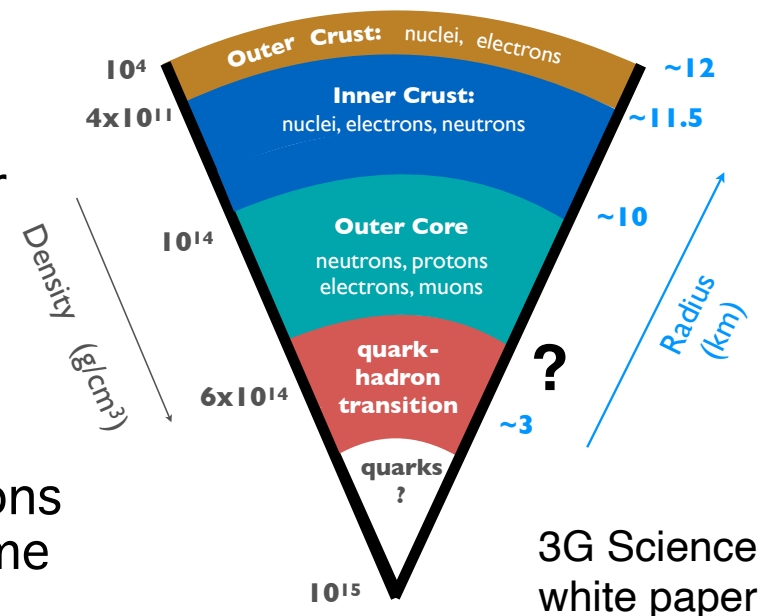
- **Introduction - dense matter and neutron stars**
- **Neutron star structure and the equation of state (EoS)**
- **Multi-messenger constraints on the EoS: what have we learned so far?**
- **Future directions**

QCD phase diagram



- lattice QCD gives good result at finite temperature, but is stymied currently at finite density
- perturbative QCD: only valid at asymptotically high densities
- can't calculate properties of cold dense matter, must observe!

- properties of ultra-dense matter in the inner cores of neutron stars (NSs)
- a challenging problem as no terrestrial experiments can probe such high densities
- also because reliable first-principle calculations break down at the strongly-interacting regime

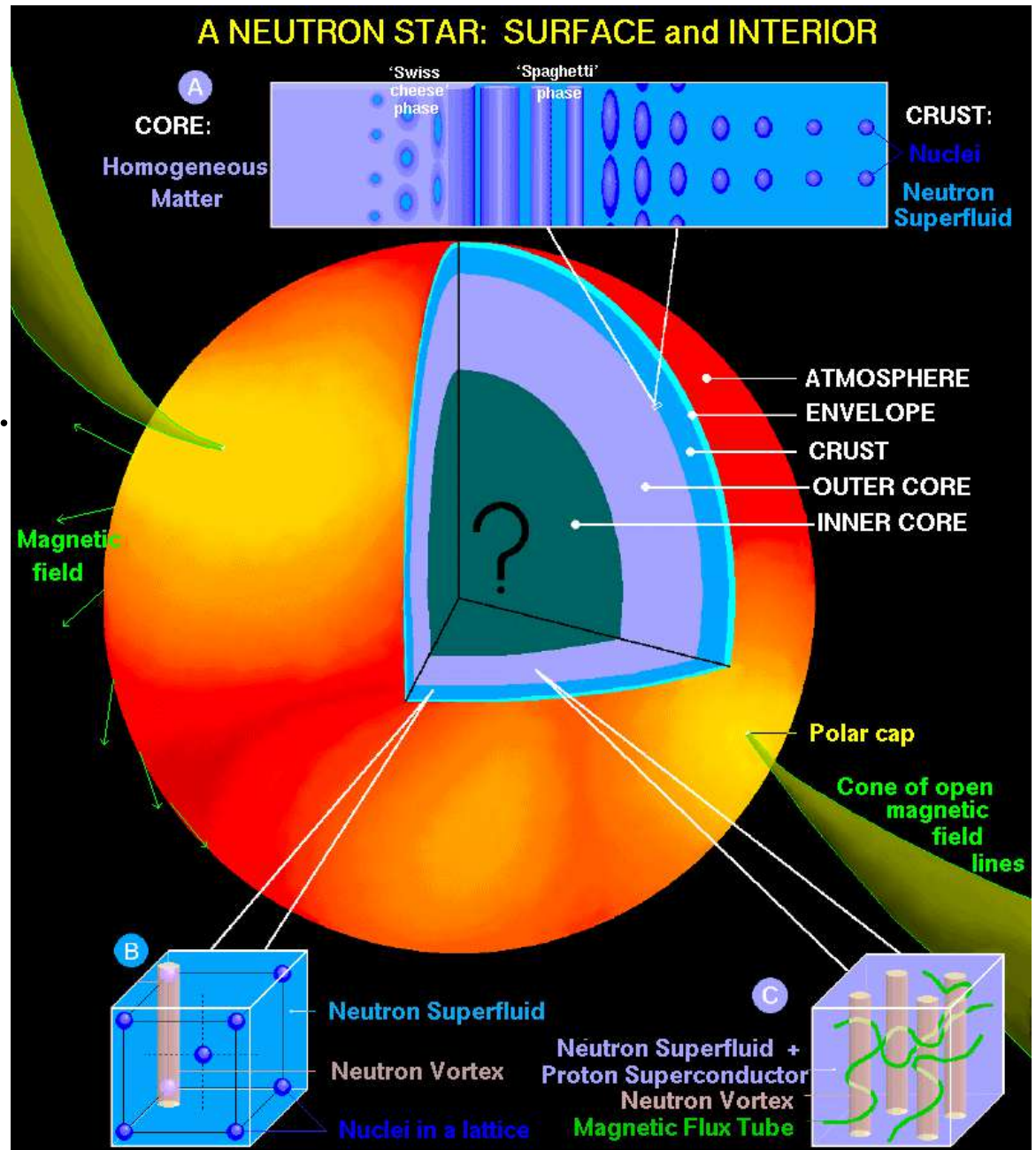


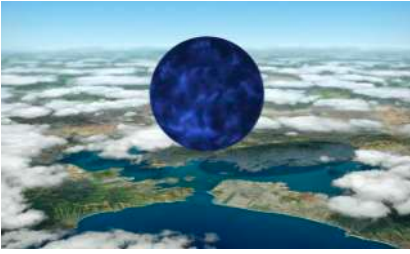
Dense matter in NSs

- stable nuclei
- neutron-rich nuclei
- neutron-rich nuclei with quasi-free neutrons
-
- homogeneous nucleonic matter (liquid)
- exotica

Fundamental questions

- what are the most relevant lower-energy degrees of freedom?
- how does deconfinement evolve as $T \rightarrow 0$ on the QCD phase diagram?



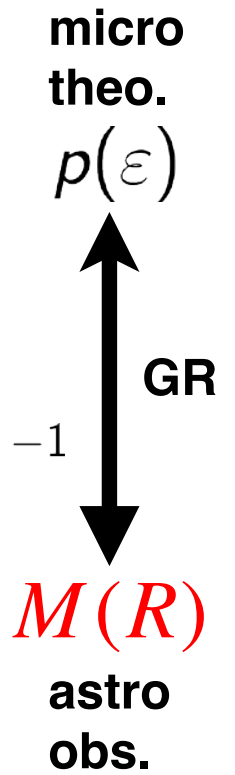


Nature's extreme labs

- for the interior of a spherical, static, relativistic star

$$\frac{dp}{dr} = -\varepsilon(r) \frac{Gm(r)}{r^2} \left[1 + \frac{p(r)}{\varepsilon(r)} \right] \left[1 + \frac{4\pi r^3 p(r)}{m(r)} \right] \left[1 - \frac{2Gm(r)}{r} \right]^{-1}$$

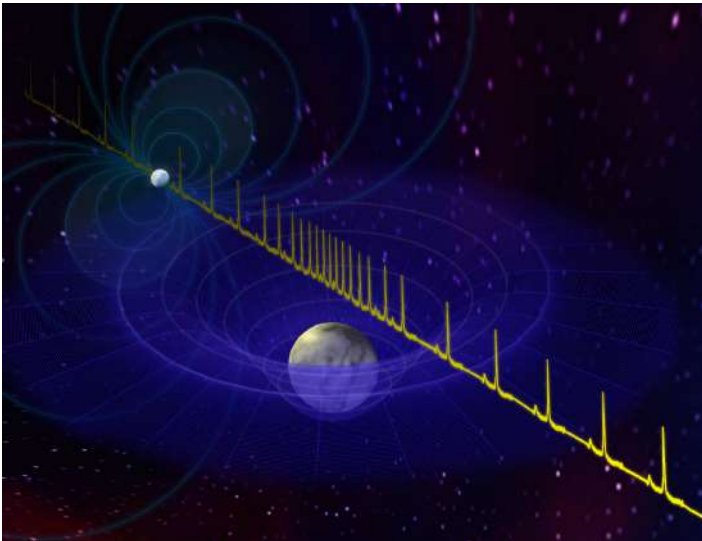
$$m(r) \equiv 4\pi \int_0^r \varepsilon(r) r^2 dr$$



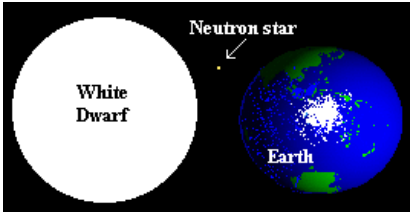
massive neutron stars $\sim 2 M_{\odot}$ do exist!

radio pulsar timing

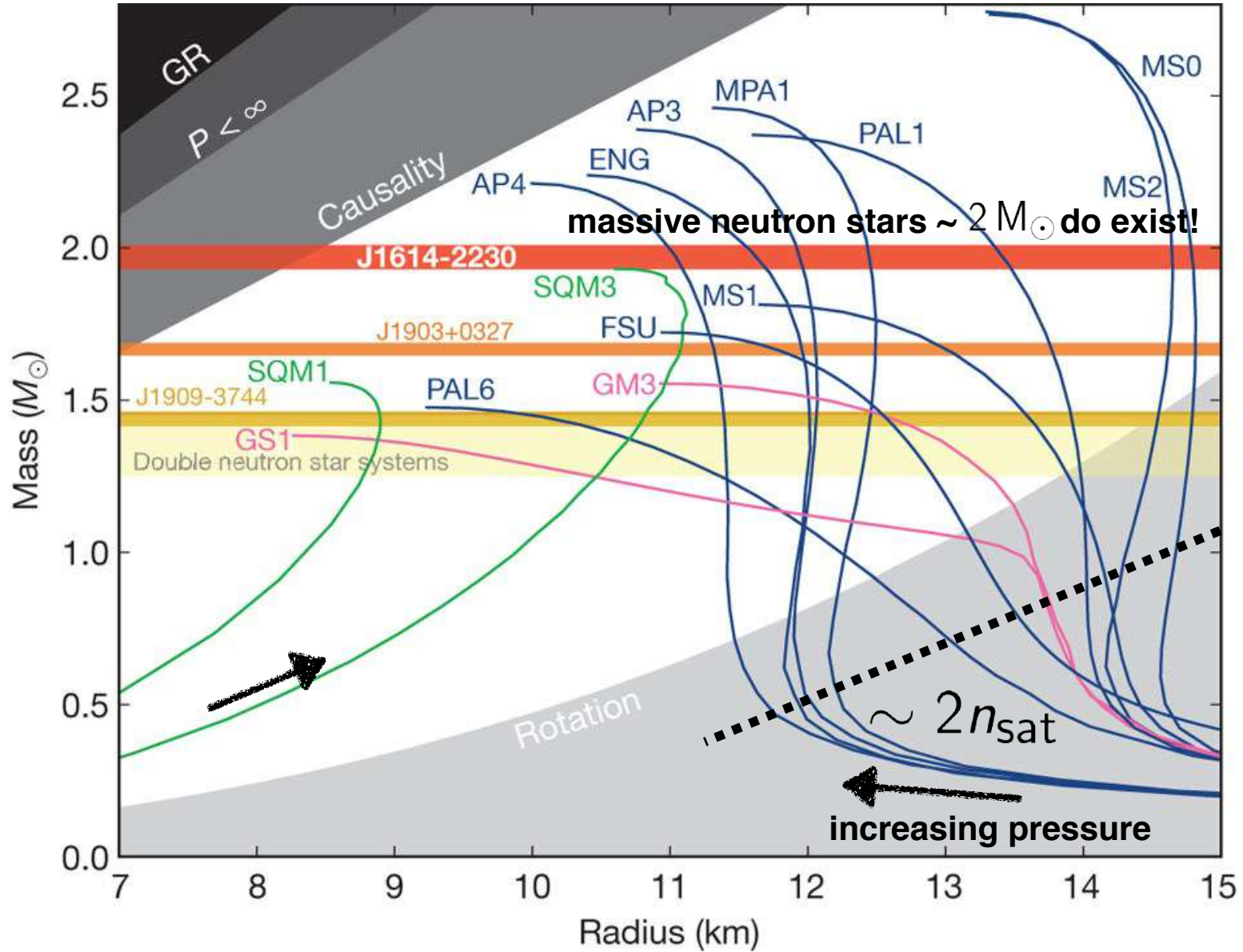
©NRAO



Source	Mass (M_{\odot})	References
PSR J1614-2230	1.97 ± 0.04	Demorest et al. (2010)
	1.928 ± 0.017	Fonseca et al. (2016)
	1.908 ± 0.016	Arzoumanian et al. (2018)
PSR J0438+0432	2.01 ± 0.04	Antoniadis et al. (2013)
PSR J0740+6620	$2.14^{+0.10}_{-0.09}$	Cromartie et al. (2019)
	2.08 ± 0.07	Fonseca et al. (2021)



NS mass-radius diagram



massive neutron stars $\sim 2 M_{\odot}$ do exist!

micro
theo.
 $p(\epsilon)$

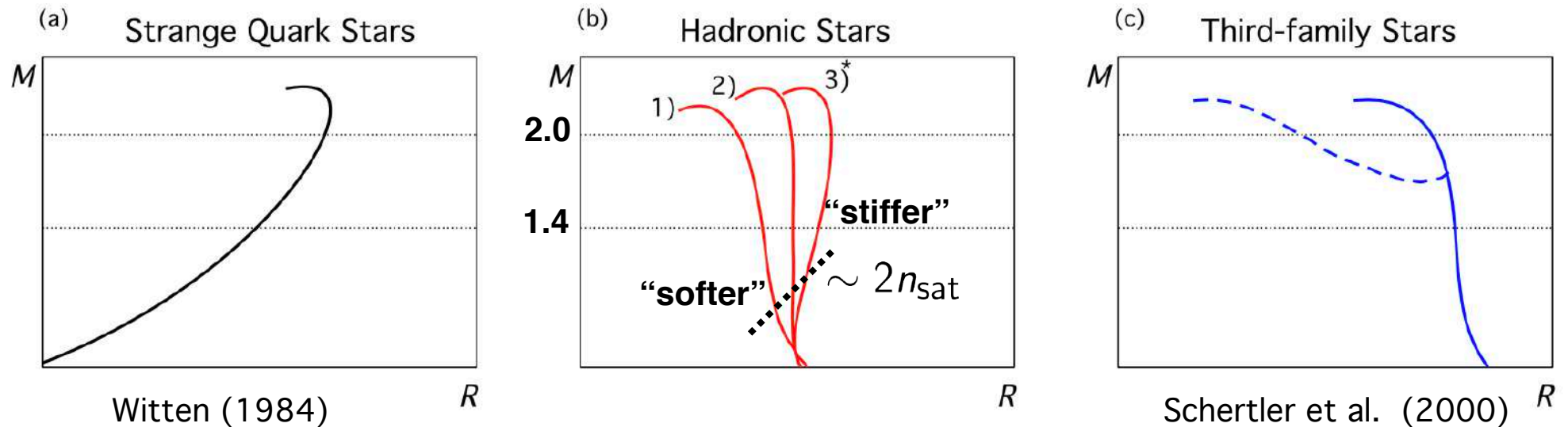
GR

$M(R)$

astro
obs.

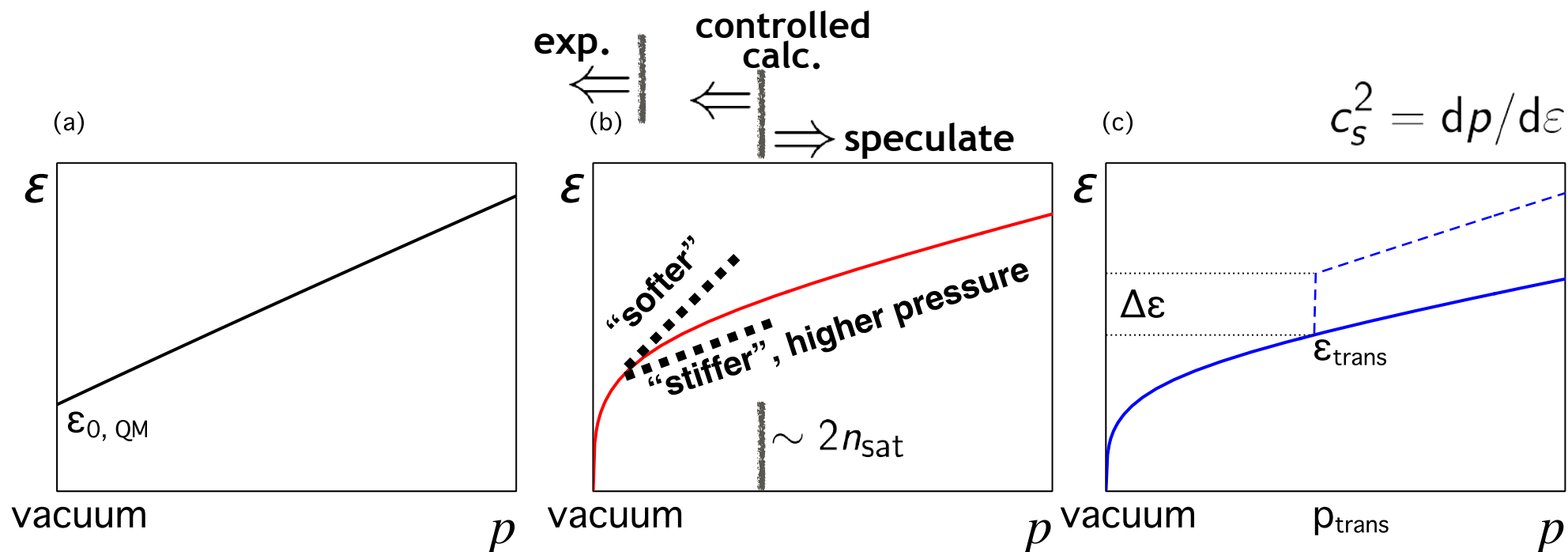
Categories of the M-R relation

SH & Prakash,
arXiv:2006.02207



- self-bound stars with a bare surface e.g. strange matter hypothesis
- continuous (and mostly smooth) profile for normal hadronic EoSs; *also possible with weak/mild phase transition or crossover
- substantial softening e.g. discontinuity in the energy density induced by a strong sharp phase transition

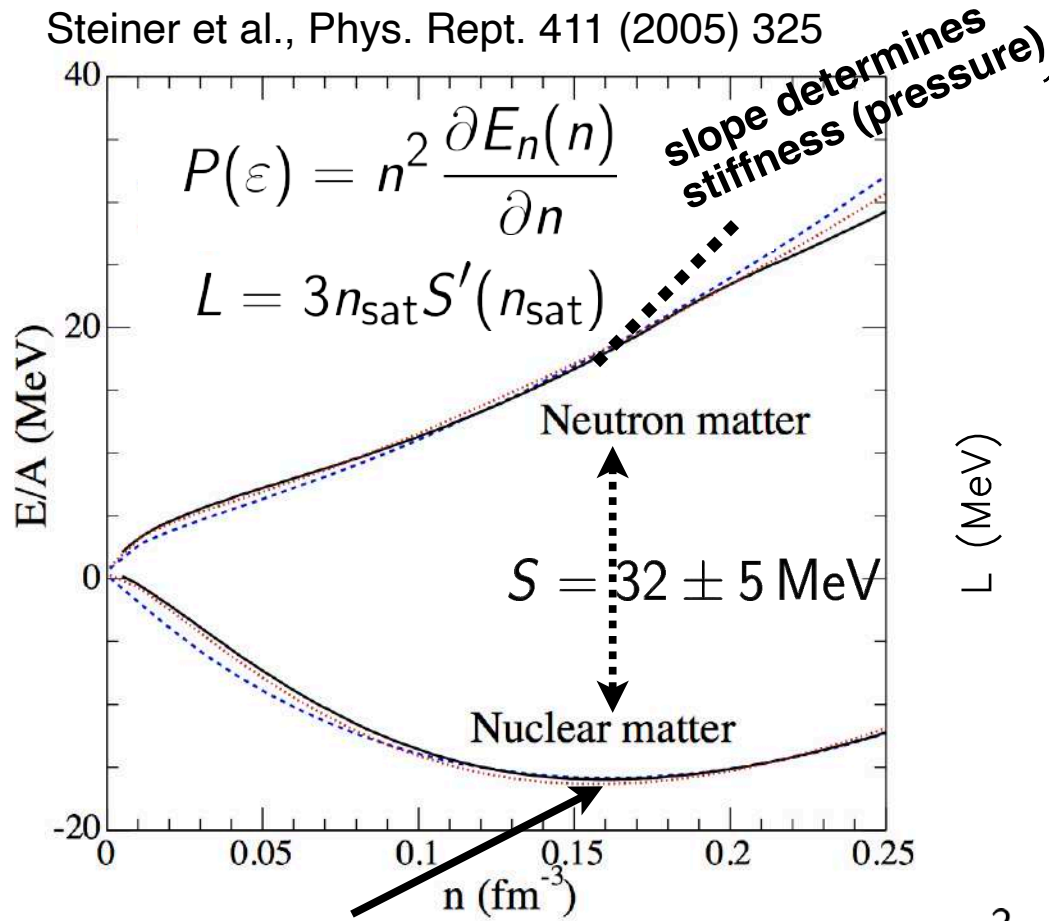
Schematic EoSs from theory



- self-bound stars with a bare surface e.g. strange matter hypothesis
- continuous (and mostly smooth) profile for normal hadronic EoSs; *also possible with weak/mild phase transition or crossover
- substantial softening e.g. discontinuity in the energy density induced by a strong sharp phase transition

From nuclei to neutron stars

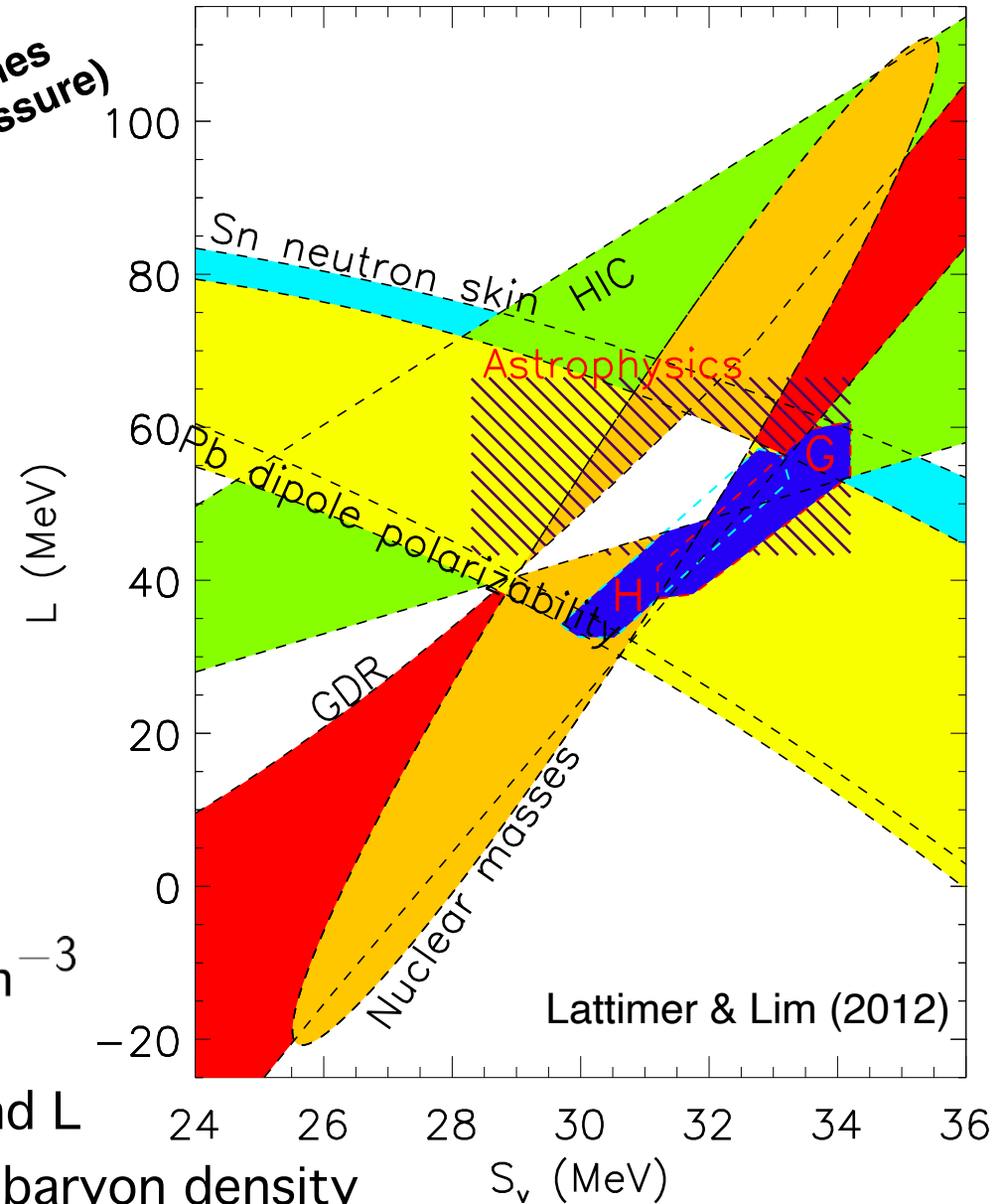
Steiner et al., Phys. Rept. 411 (2005) 325



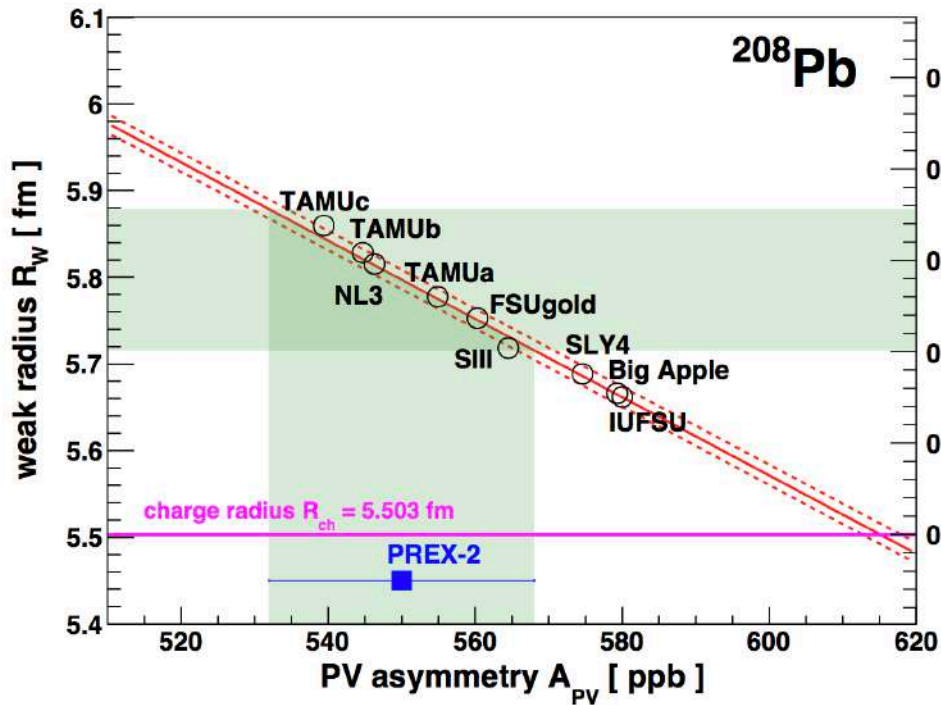
stable nuclei is here
[self-bound system]

$$n_{\text{sat}} \simeq 0.16 \text{ fm}^{-3}$$

- nuclear experiments correlate S and L
- theory extrapolates in isospin and baryon density



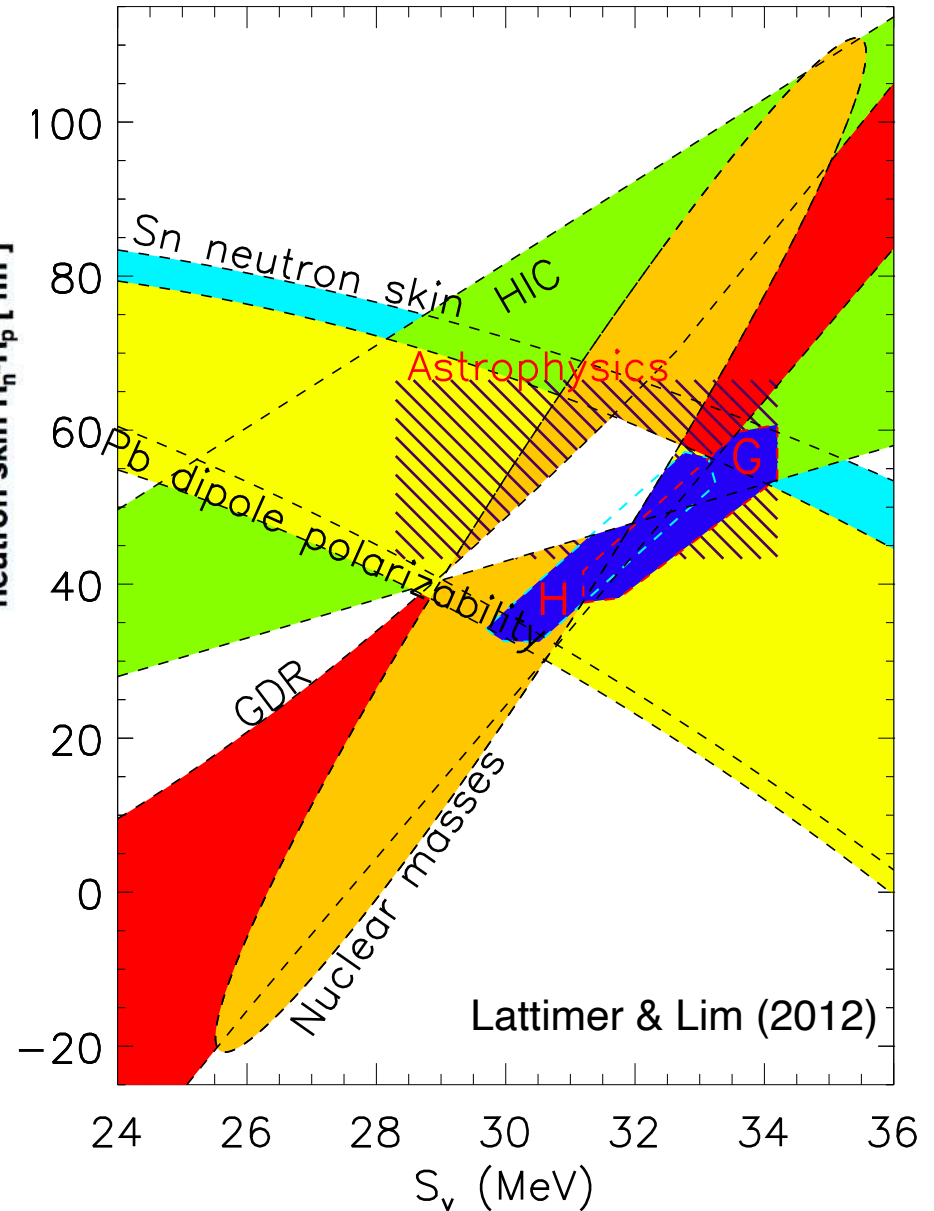
From nuclei to neutron stars



PREX Collaboration, arXiv: 2102.10767

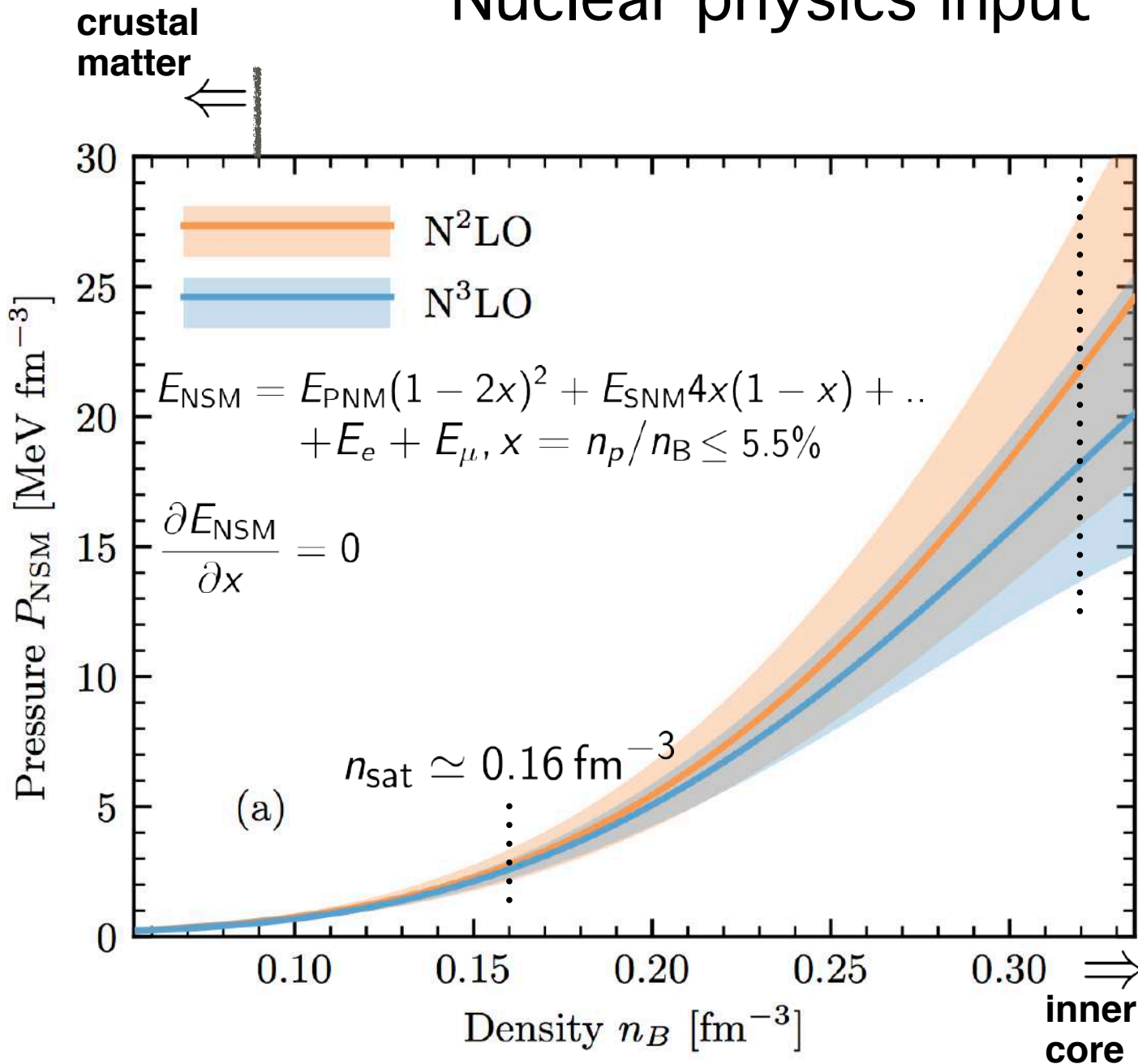
new! neutron skin measurement @JLab

- under hot debate: pressures at sub-nuclear densities and the inferred value of L



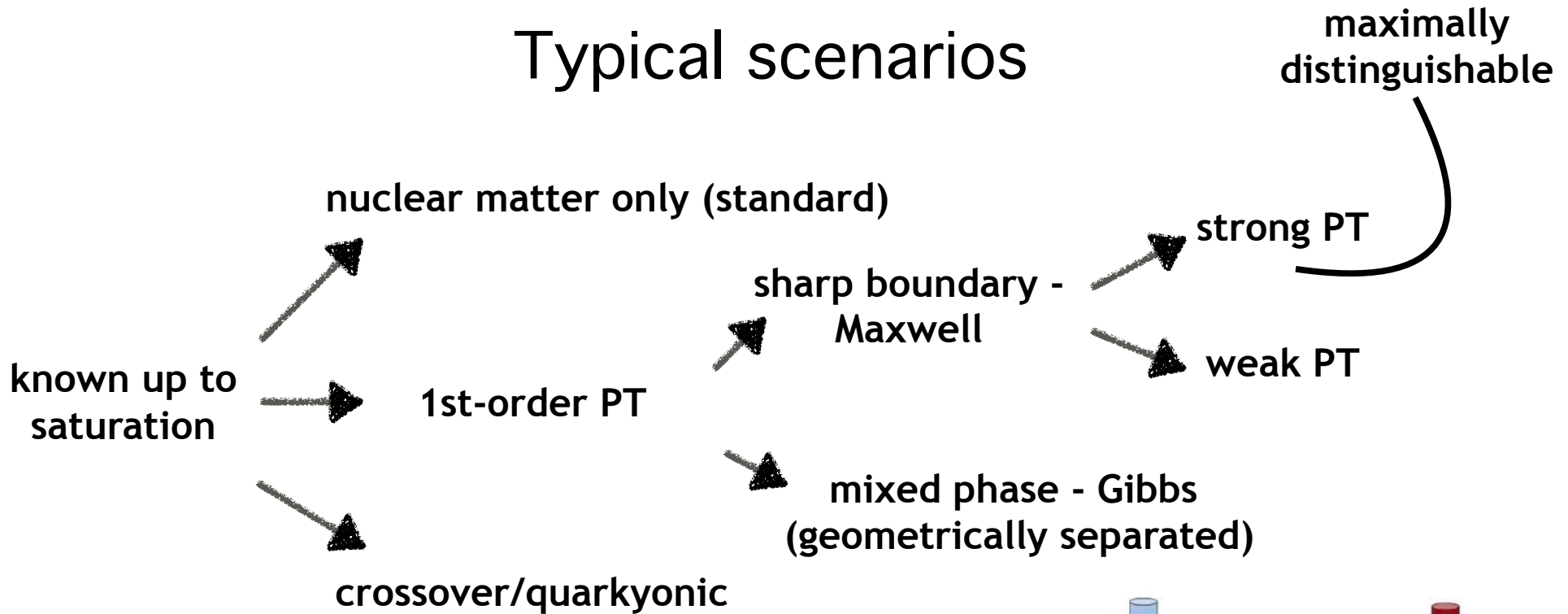
Nuclear physics input

Drischler, **SH**,
Lattimer, Prakash,
Reddy and Zhao,
arXiv:2009.06441

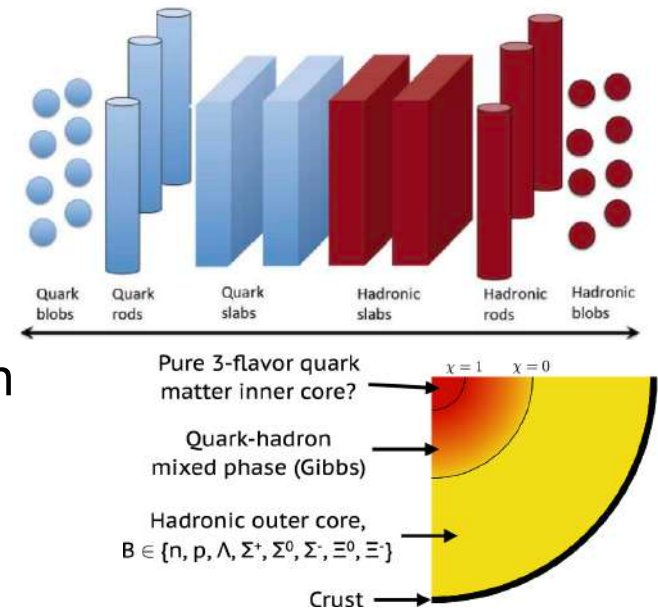


- pressure at low densities (**outer core**) controls typical NS radii: stiff or soft?
- reliably **quantified** uncertainties from chiEFT for beta-equilibrated NSM
- less than $\sim 5\%$ deviation from PNM pressures
- to **extrapolate** or match at higher densities in the inner core

Typical scenarios



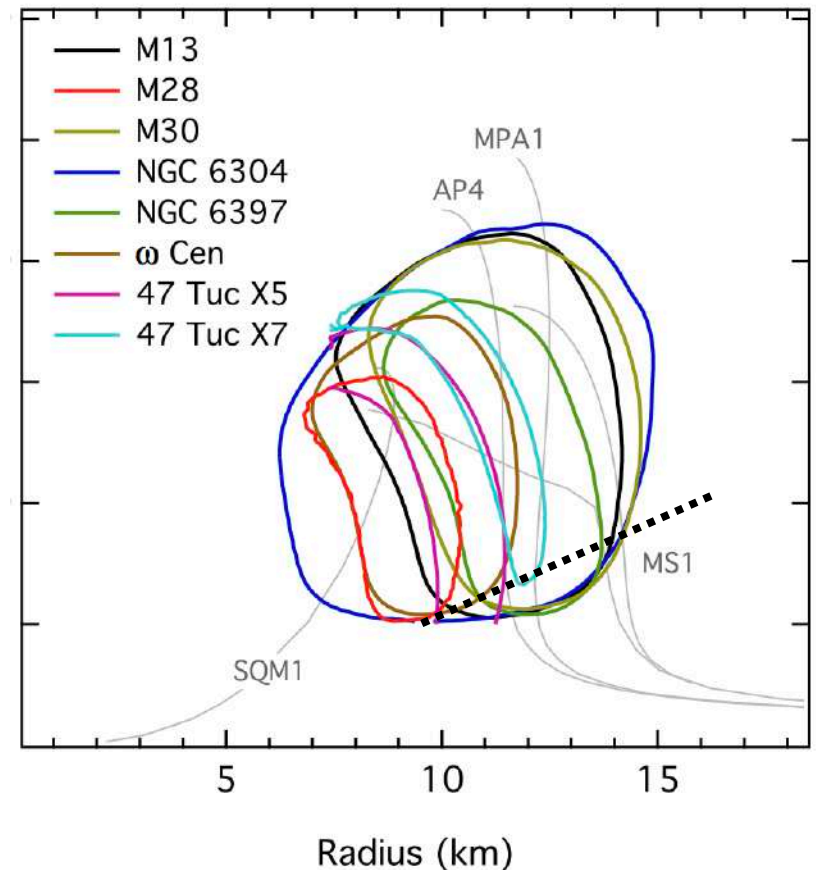
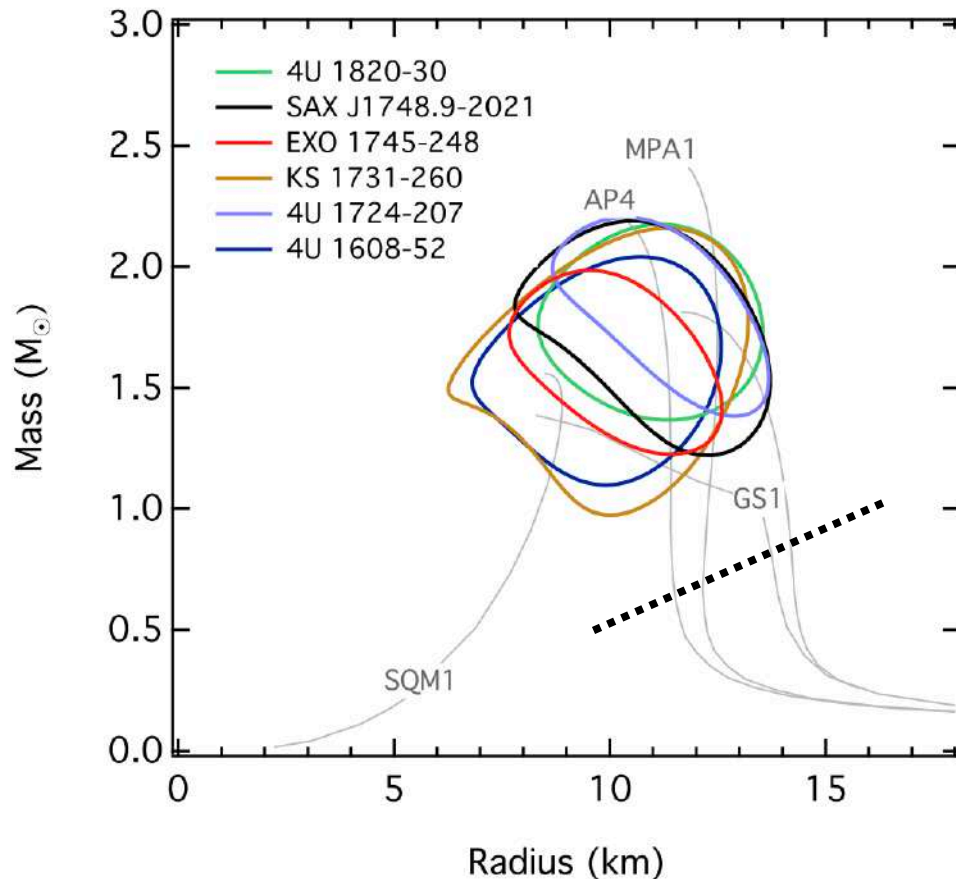
- masquerade problem: likely indistinguishable through observations that constrain M-R only
- smooth crossover: no more easily understood in terms of hadrons than in terms of quarks
- 1st-PT: mixed phase (Gibbs) is favored if the hadron/quark surface tension is small



X-ray probes of NS radii

- conventional methods of radius estimates through surface photon emission detection suffer from large uncertainties

Ozel & Freire (2016);
Steiner et. al (2016)



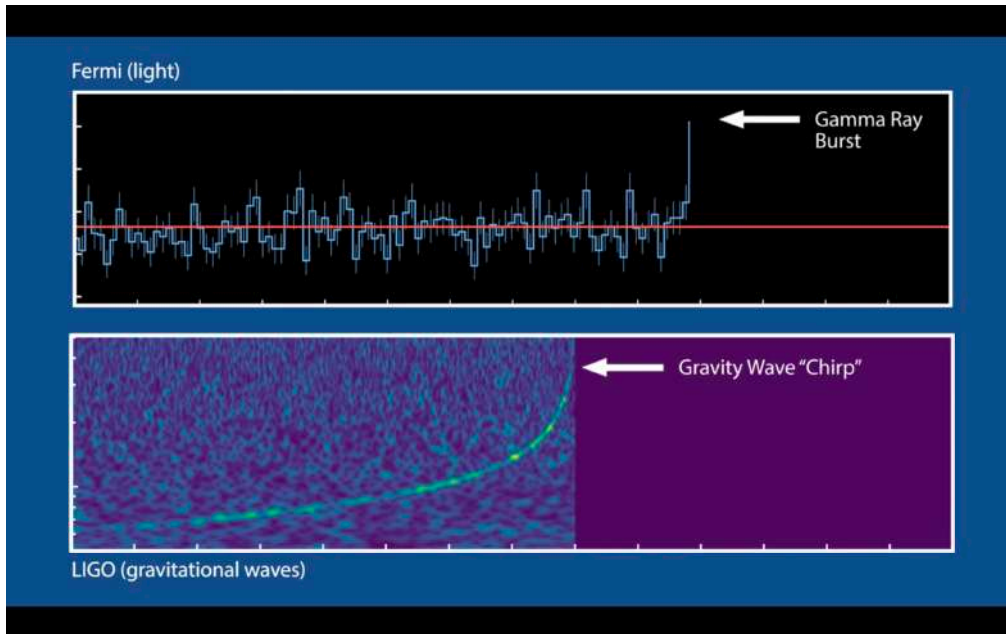
- photospheric radius expansions (PREs)

- quiescent low-mass x-ray binaries (QLMXBs)

First BNS merger detection

GW170817 that unveiled the multi-messenger era

- “hear” cosmic collisions between densest astronomical objects
- follow-up E&M signals; “see” e.g. evidence for nucleosynthesis



credit: Karan Jani/Georgia Tech

EoS affects GW emission during inspiral

- tidal deformability

$$\Lambda \equiv \frac{\lambda}{M^5} \equiv \frac{2}{3} k_2 \left(\frac{Rc^2}{GM} \right)^5$$

- compactness $k_2 \propto (M/R)^{-1}$

$$\beta \equiv M/R$$

$$Q_{ij} = -\lambda \varepsilon_{ij}$$

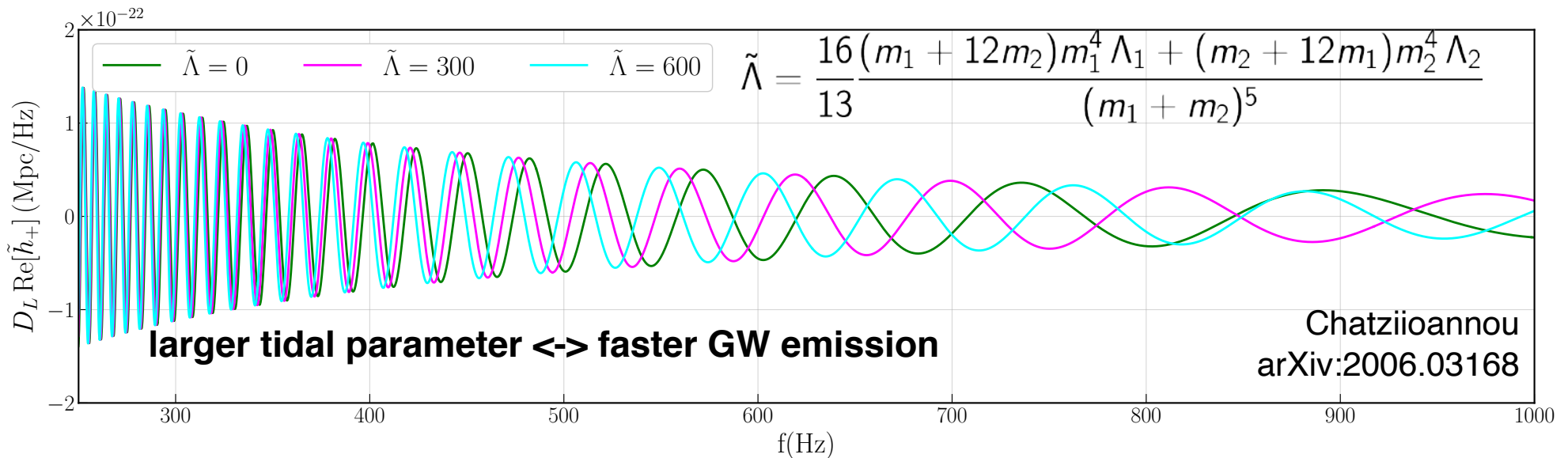
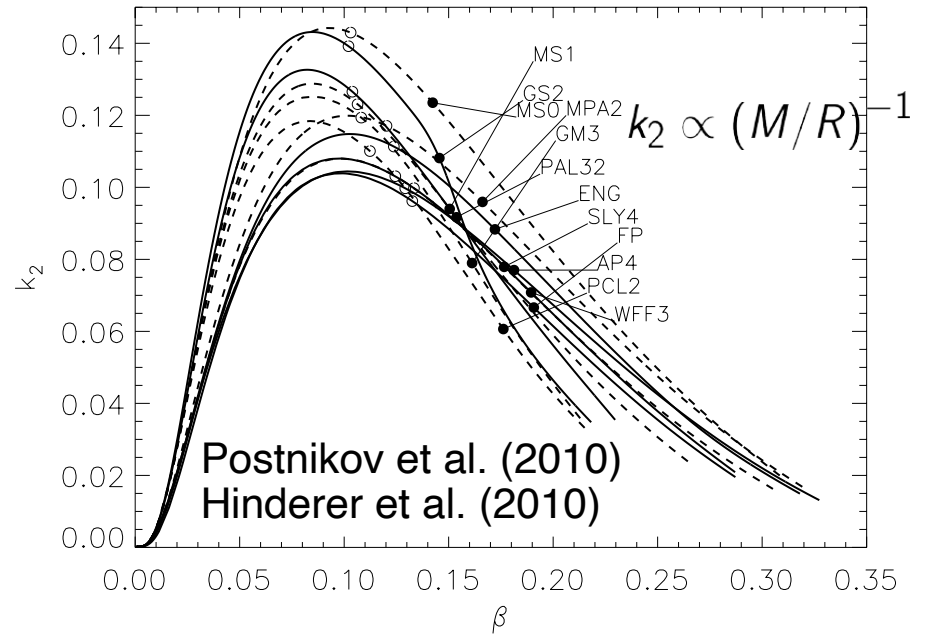


tidal effects



Impact on pre-merger GW signal

- tidal Love number depends on the EoS and the compactness M/R
- matter effects (NSs) leave imprints in the waveform - distinguish from point-particles (BHs)
- much cleaner systematics

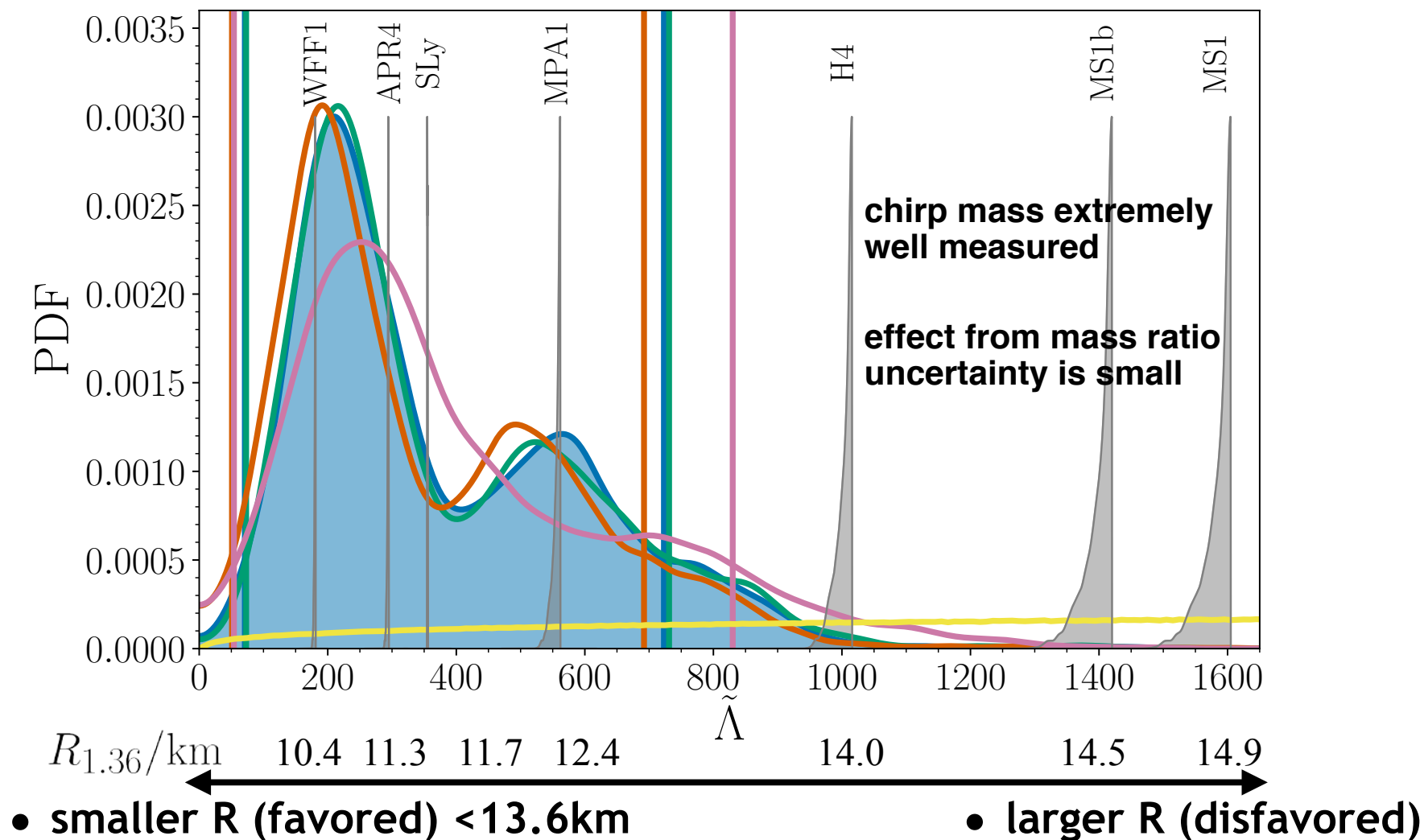


Measuring NS radius with GWs

$$m_{\text{tot}} = 2.73_{-0.01}^{+0.04} M_{\odot}$$

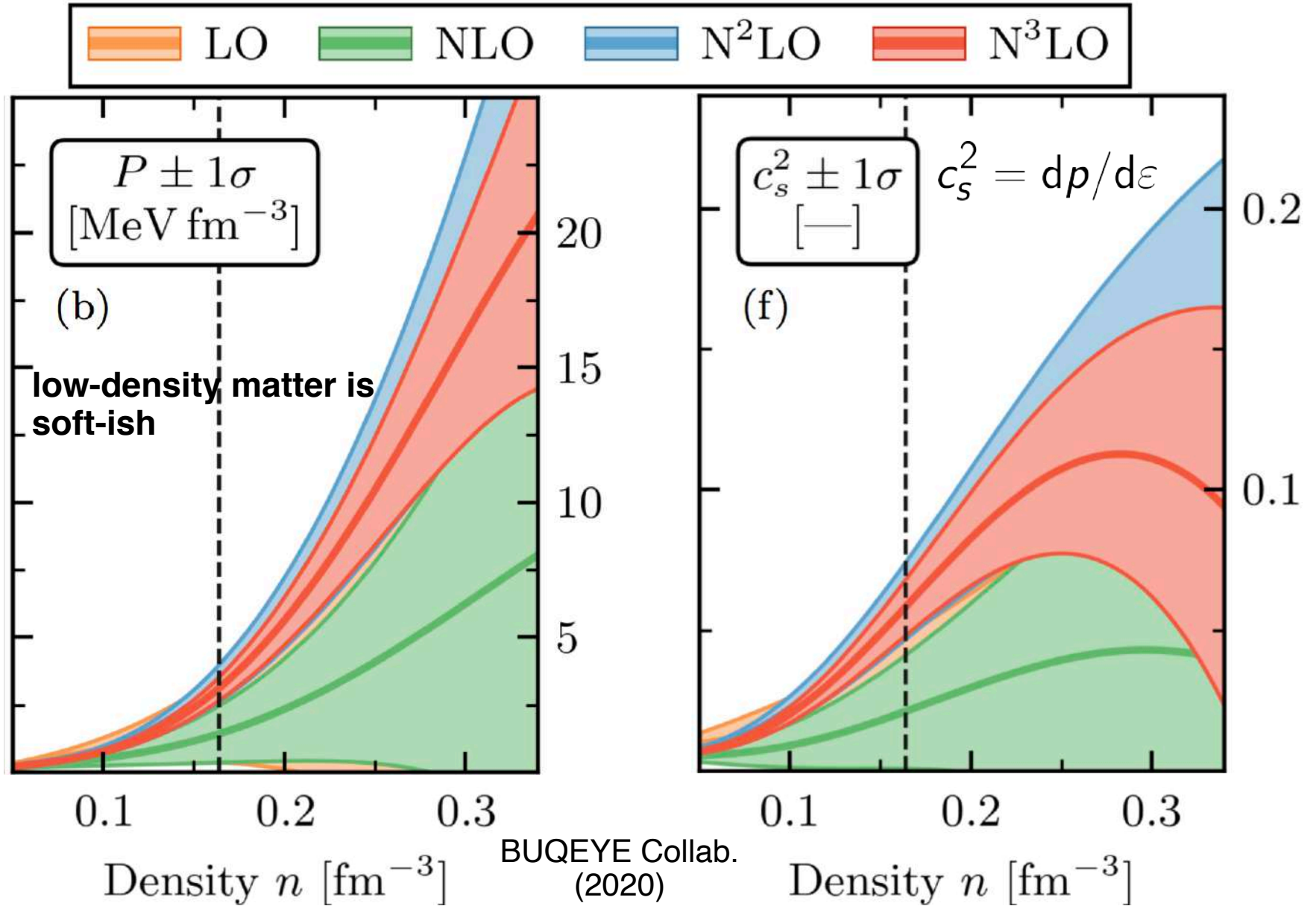
$$\mathcal{M} = (m_1 m_2)^{3/5} / (m_1 + m_2)^{1/5} = 1.186_{-0.001}^{+0.001} M_{\odot}$$

LVC collaboration, arXiv:1805.11579



Pure neutron matter (PNM)

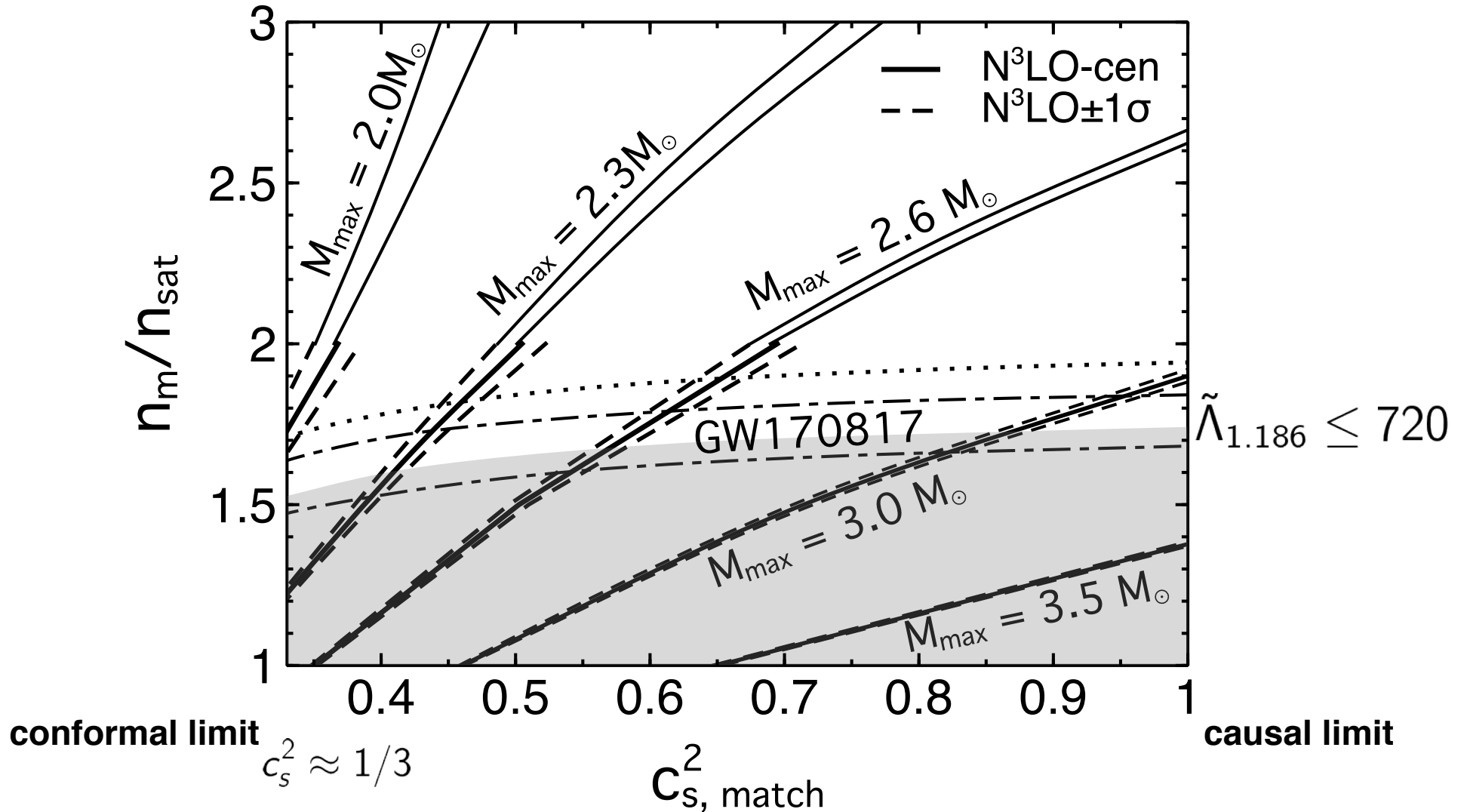
Drischler et al.
PRL 125, 202702
arXiv:2004.07232



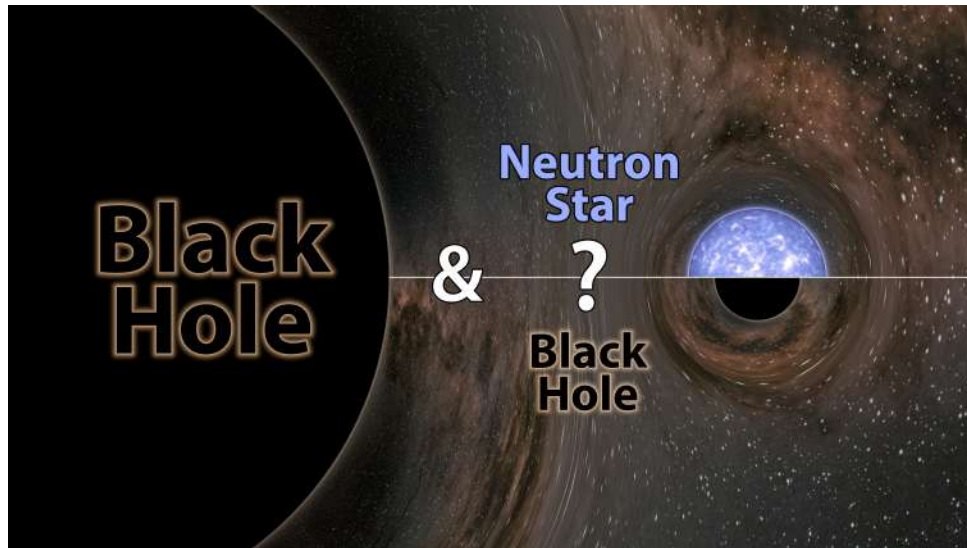
GW + heavy pulsars

Drischler, **SH**,
Lattimer, Prakash,
Reddy and Zhao,
arXiv:2009.06441

- sound speed in the **core** and **when** rapid stiffening in the EoS begins



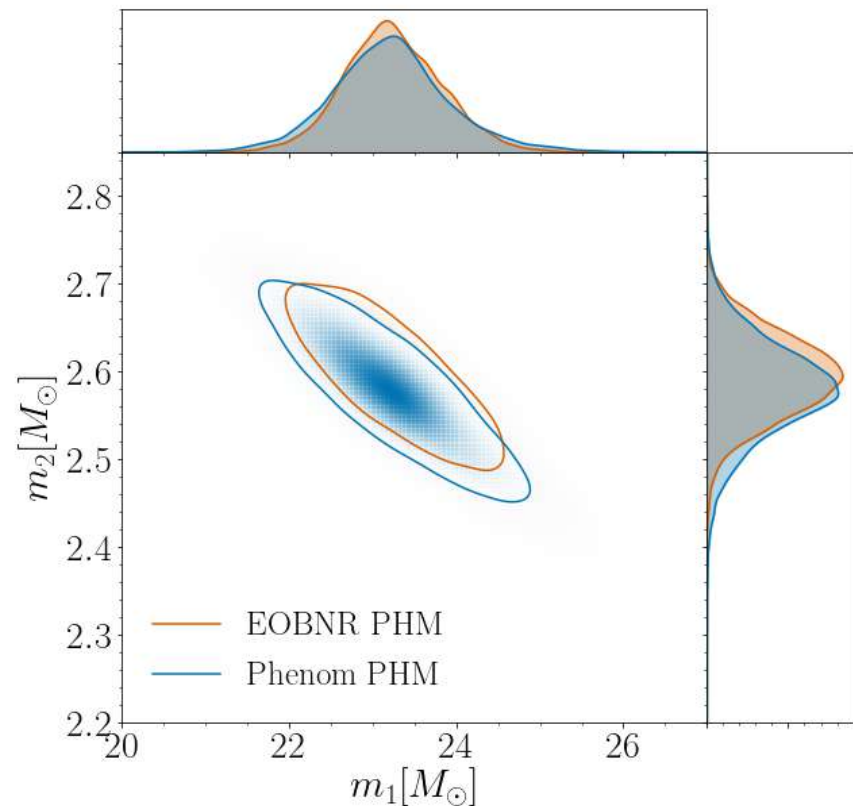
GW190814



- extremely loud event produced by the inspiral and merger of two compact objects -- one, a black hole, and the other of **undetermined** nature
- the mass measured for the lighter compact object makes it either the lightest black hole or the heaviest neutron star ever discovered

the most asymmetric system observed

$$m_1 = 23.2^{+1.1}_{-1.0} M_{\odot} \quad m_2 = 2.59^{+0.08}_{-0.09} M_{\odot}$$



LVC collaboration, arXiv:2006.12611

Sound speed in the core

$$c_s^2(r) \equiv dp(r)/d\varepsilon(r)$$

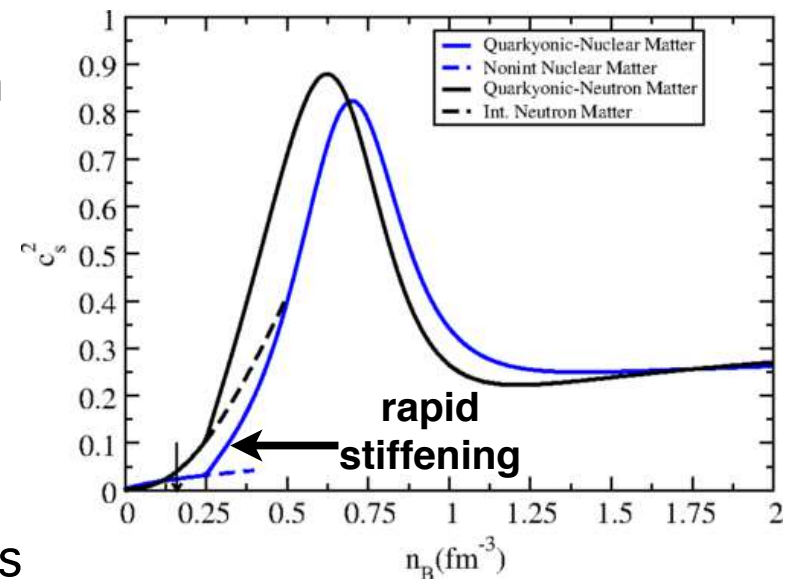
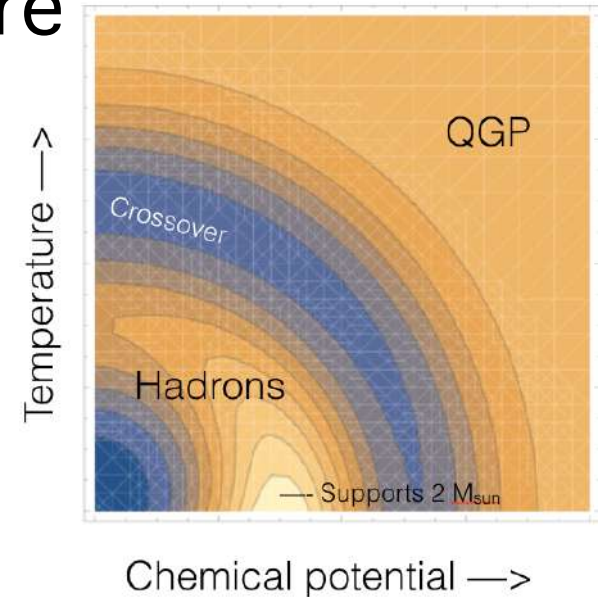
how fast pressure rises with energy density

Possible behavior in neutron star interiors

- minimal scenario of normal nuclear matter: (smoothly) continuous function of pressure
- first-order phase transition scenario: finite energy density discontinuity induces sudden softening near the phase boundary
- crossover scenario/quarkyonic matter

Limits

- asymptotically high density: $\sim 1/3$
- $\sim 4-8$ times saturation: supports massive NSs
- high-T: matches lattice calc./heavy-ion data



McLerran & Reddy,
PRL 122, 122701 (2019)

Sound speed in the core

$$c_s^2(r) \equiv dp(r)/d\varepsilon(r)$$

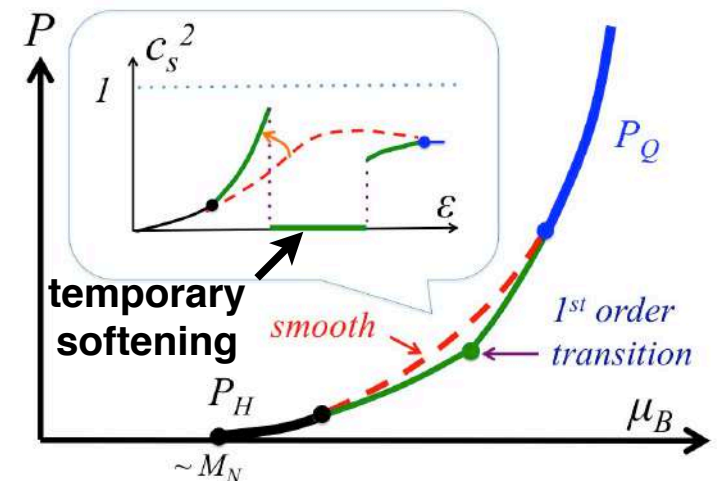
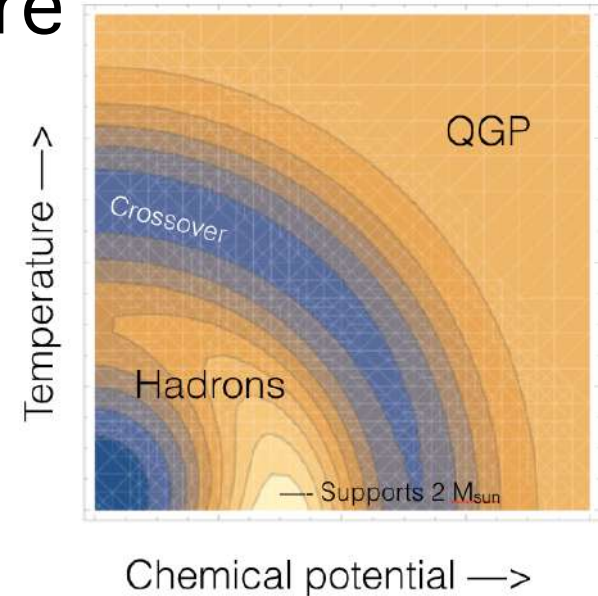
how fast pressure rises with energy density

Possible behavior in neutron star interiors

- minimal scenario of normal nuclear matter: (smoothly) continuous function of pressure
- first-order phase transition scenario: finite energy density **discontinuity** induces sudden **softening** near the phase boundary
- crossover scenario/quarkyonic matter

Limits

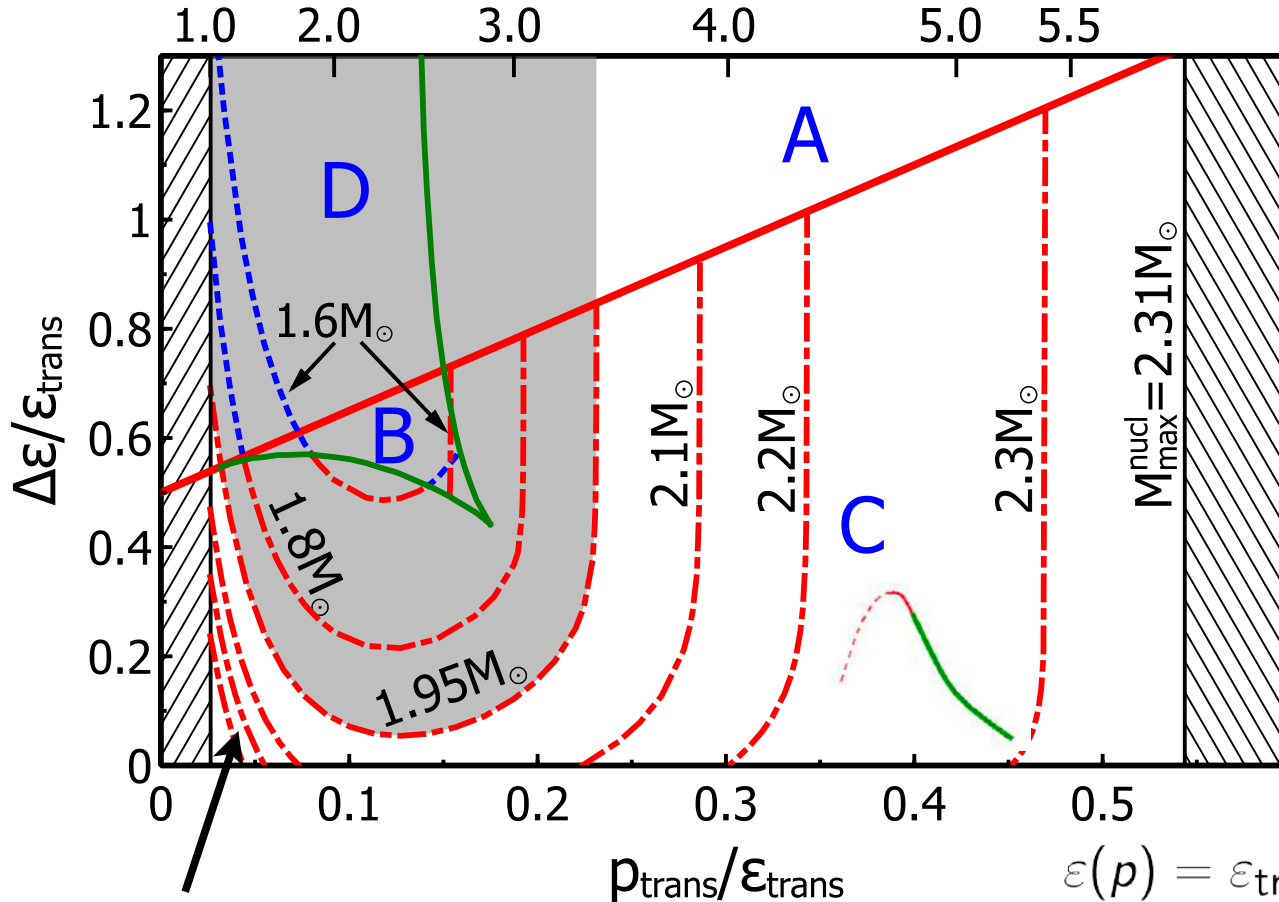
- asymptotically high density: $\sim 1/3$
- $\sim 4-8$ times saturation: supports massive NSs
- high-T: matches lattice calc./heavy-ion data



Baym et al. arXiv:1707.04966
Rept. Prog. Phys. 81, 056902

Constraints from max. mass

DBHF (stiff) NM, $c_{\text{QM}}^2 = 1/3$
 n_{trans}/n_0



- with weakly interacting quarks, very limited to reach two solar masses
- high transition density scenario: resembles no PT; short extension
- low transition density scenario: no twin stars

Generic ansatz

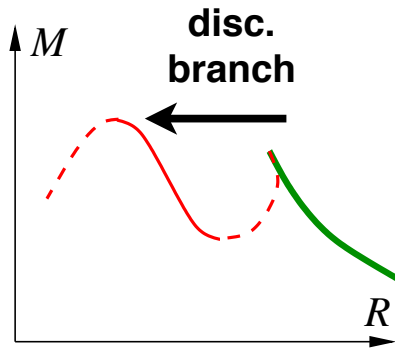
$$\varepsilon(p) = \varepsilon_{\text{trans}} + \Delta\varepsilon + c_{\text{QM}}^{-2}(p - p_{\text{trans}})$$

still survives the conformal limit

e.g. a population of BNS events

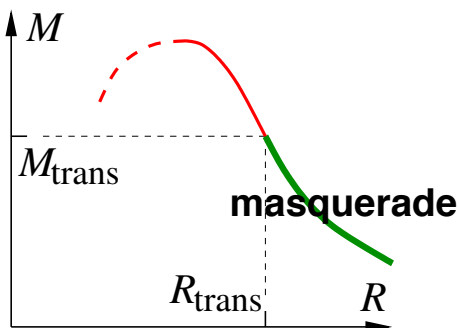
Chatziioannou & **SH**, arXiv:1911.07091

SH & Steiner, arXiv:1810.10967



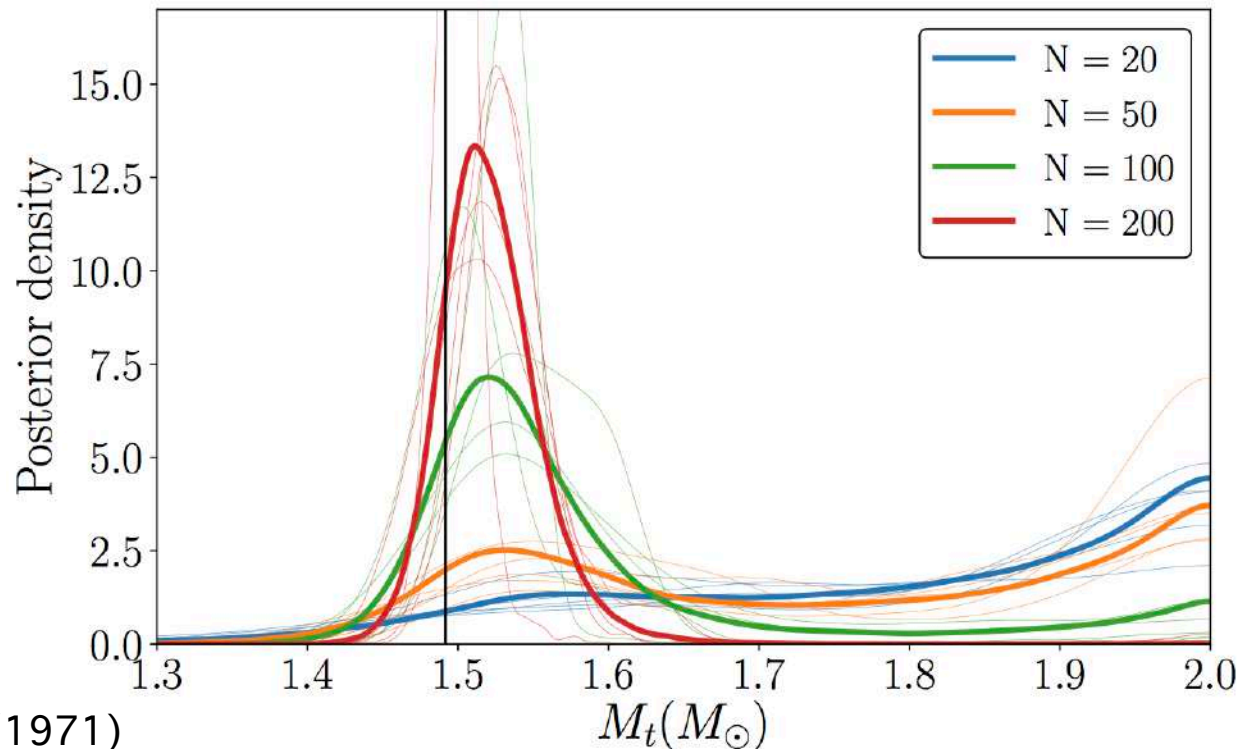
Alford, **SH** & Prakash

PRD 88, 083013 (2013)



critical strength to trigger an instability

$$\frac{\Delta \epsilon_{\text{crit}}}{\epsilon_{\text{trans}}} = \frac{1}{2} + \frac{3}{2} \frac{p_{\text{trans}}}{\epsilon_{\text{trans}}} \quad \text{Seidov (1971)}$$



- might identify third-family stars [strong 1st-OPT] with **pre-merger** GWs
- requires multiple (**N~50-100**) future detections to separate different families: NS-NS, NS-HS, HS-HS mergers

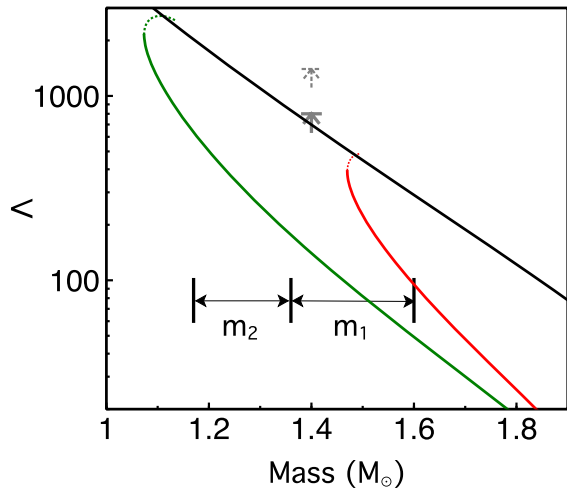
e.g. a population of BNS events

$$\Lambda \equiv \frac{\lambda}{M^5} \equiv \frac{2}{3} k_2 \left(\frac{Rc^2}{GM} \right)^5$$

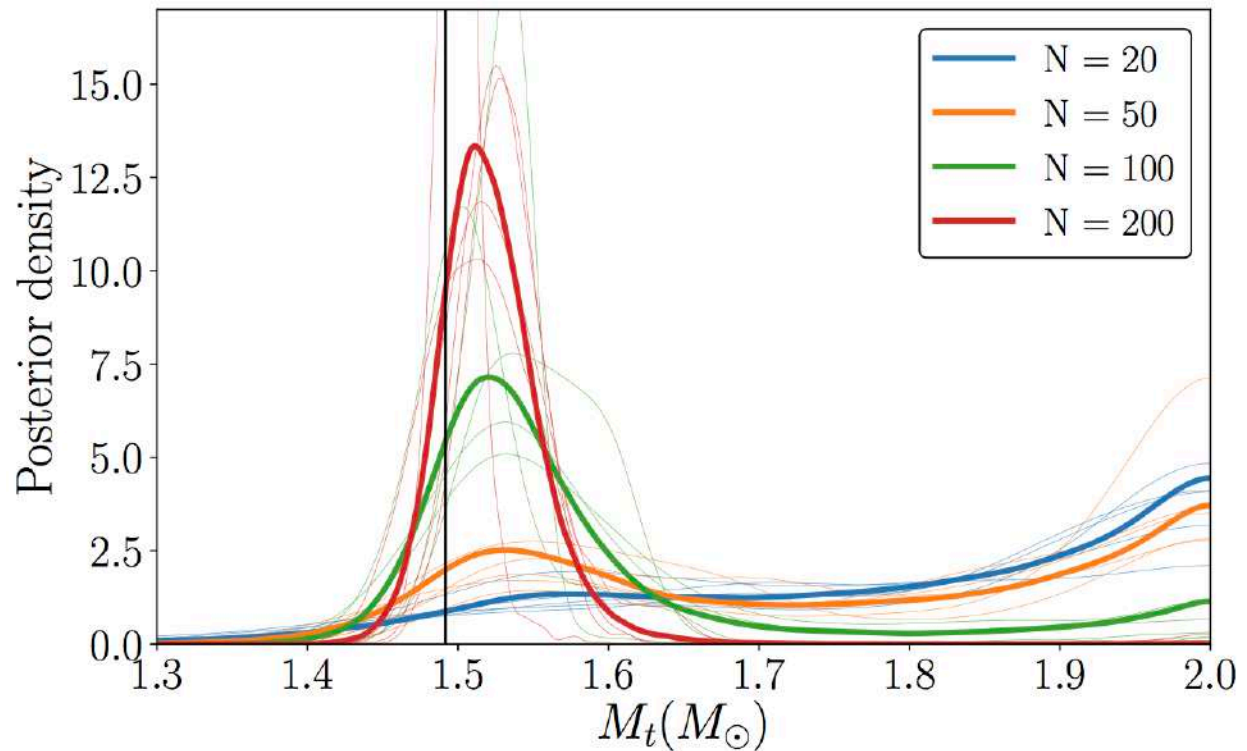
DBHF + CSS ($c_{QM}^2=1$)

Chatziioannou & SH, arXiv:1911.07091

SH & Steiner, arXiv:1810.10967



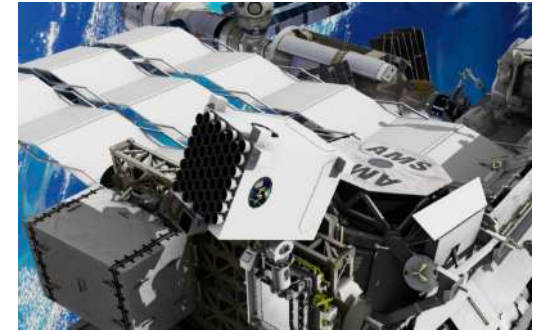
most populated if the normal branch > 13 km and the high density matter is still **strongly interacting** $c_s^2 \gtrsim 0.4$



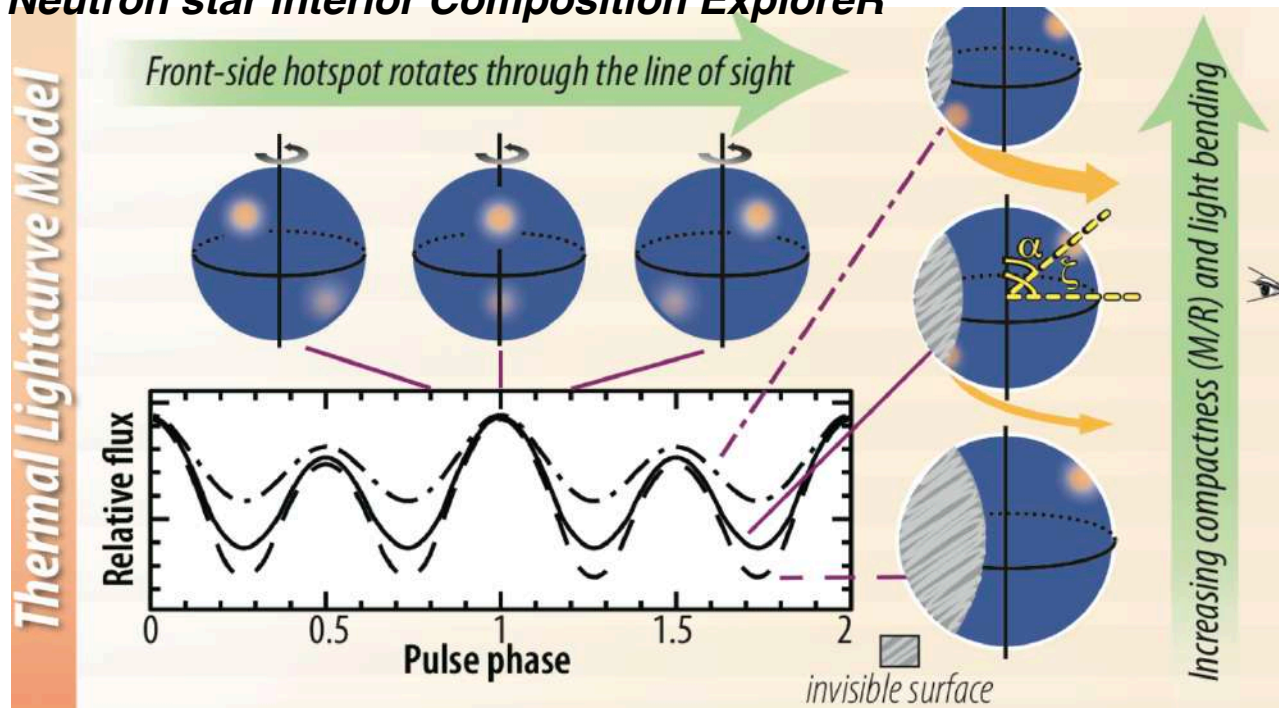
- might identify third-family stars [strong 1st-OPT] with **pre-merger** GWs
- requires multiple (**N~50-100**) future detections to separate different families: NS-NS, NS-HS, HS-HS mergers

NS radii from hotspots

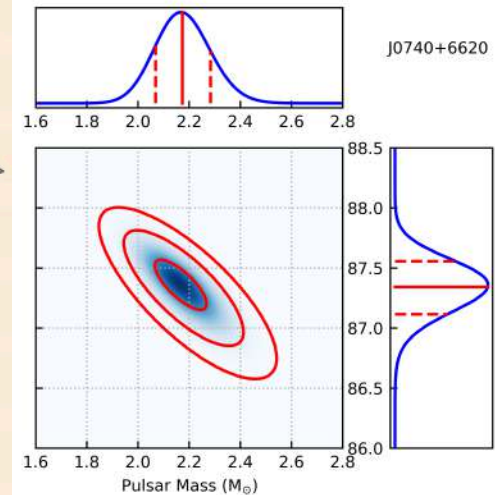
- light-curve modeling of x-ray pulse profiles that are sensitive to the stellar compactness M/R



Neutron star Interior Composition ExploreR



PSR J0740+6620

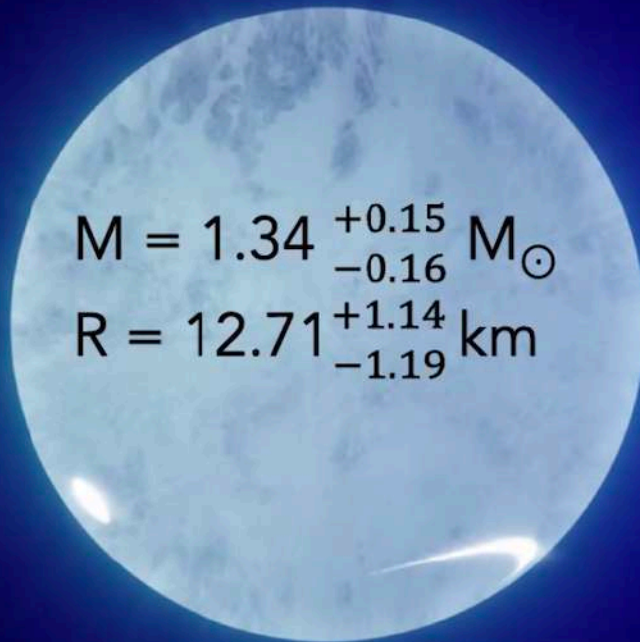


Cromartie et al.
Nature Astronomy (2019)

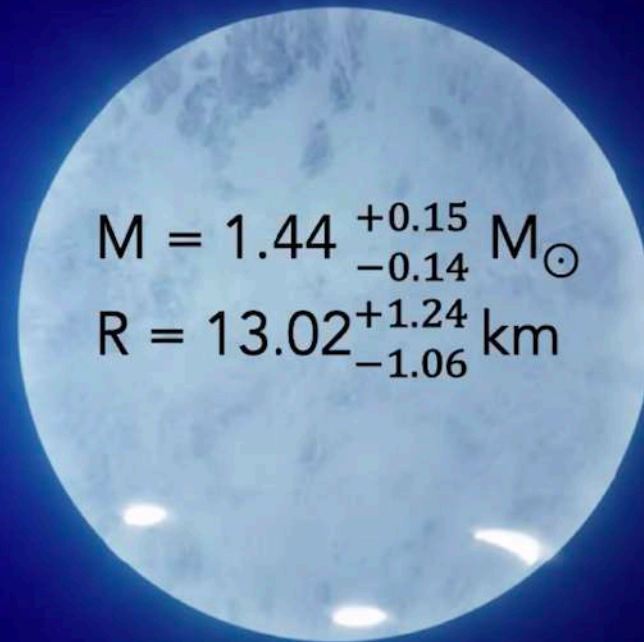
- most recent data on the heaviest NS known so far: combined information with precise mass measurements through Shapiro delay (radio)

Independent NICER team analyses

Riley et al. 2019



Miller et al. 2019

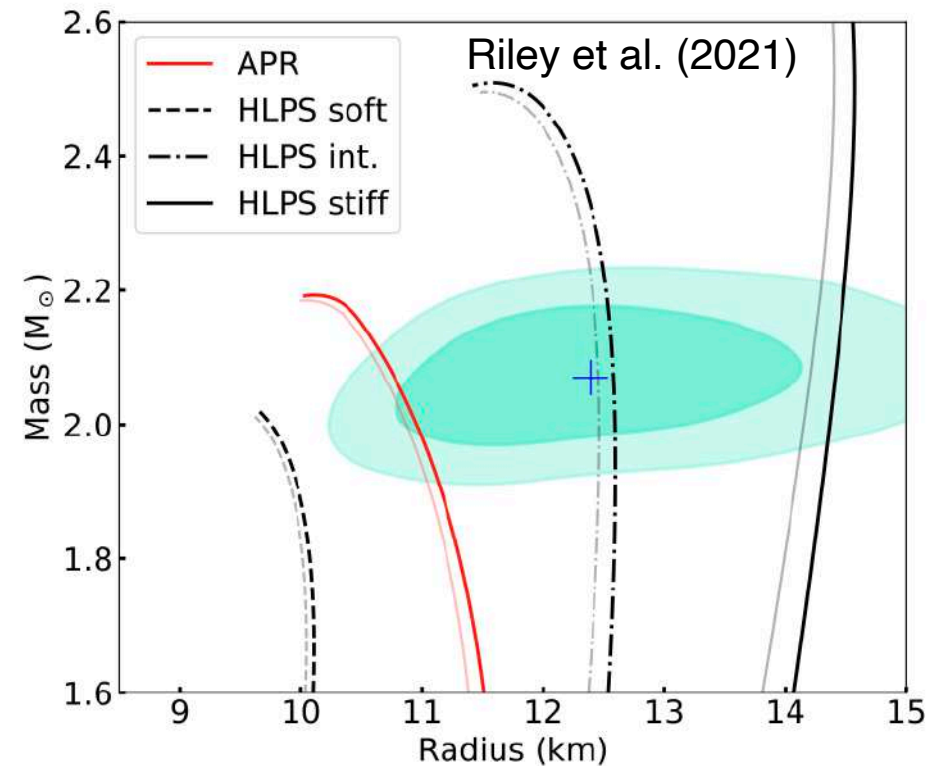
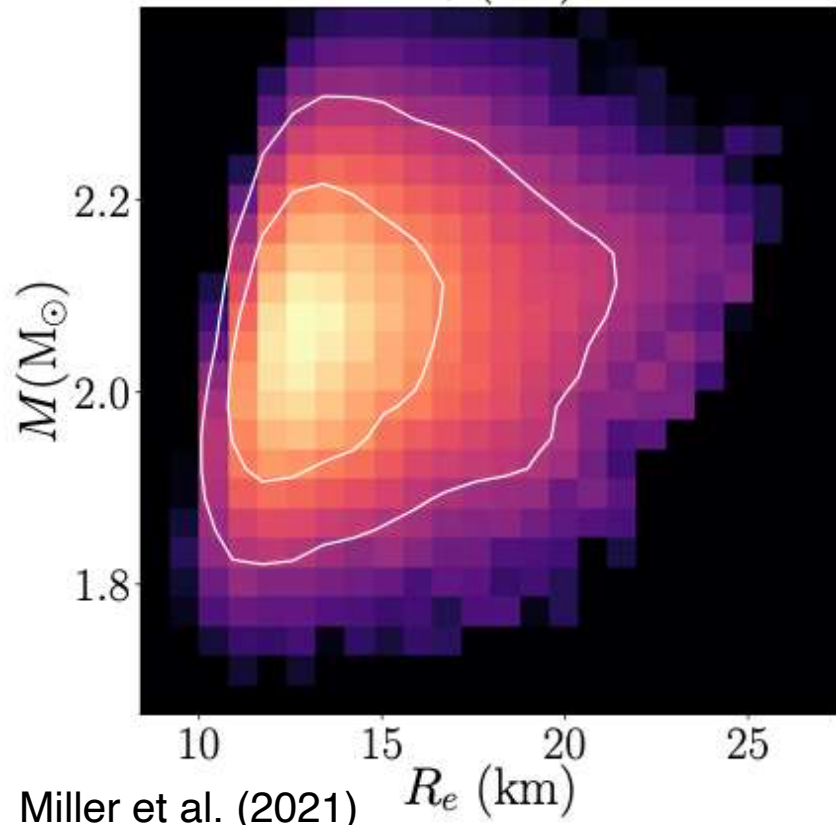


Credit: NASA's Goddard Space Flight Center/CI Lab

PSR J0030+0451

Results published together in an ApJ Letters Focus Issue in December 2019

NS radii from hotspots



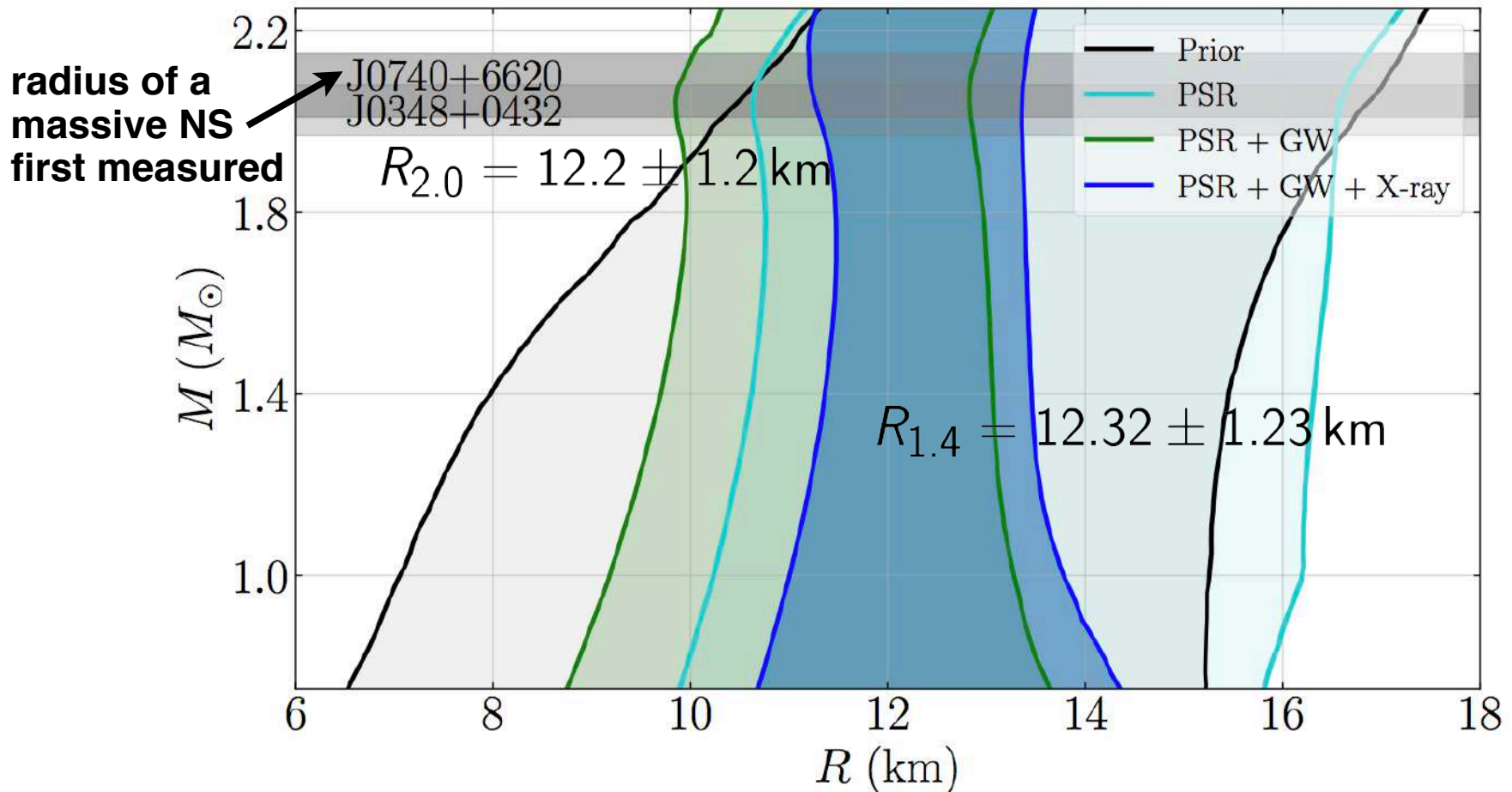
new! radius of PSR J0740+6620 $\sim 2 M_\odot$ $13.7^{+2.6}_{-1.5}$ km vs. $12.4^{+1.3}_{-1.0}$ km

previously: PSR J0030+0451 $\sim 1.4 M_\odot$

- analyses of waveforms produced by hotspots of rotation-powered pulsars
- tend to favor relatively stiffer EoS at intermediate ($2 \sim 3 n_{\text{sat}}$) densities

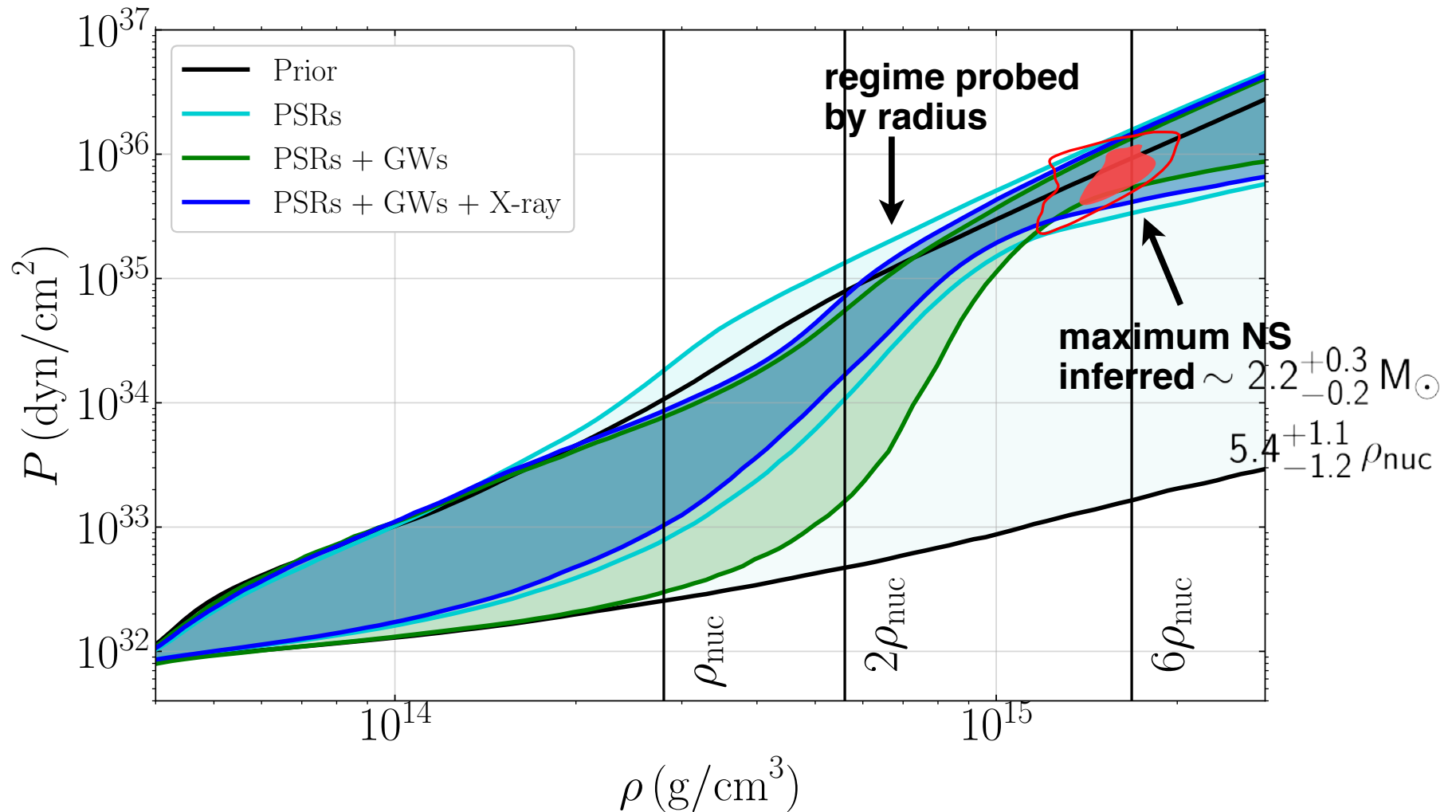
Multimessenger constraints

Legred et al.
(including **SH**),
arXiv:2106.05313



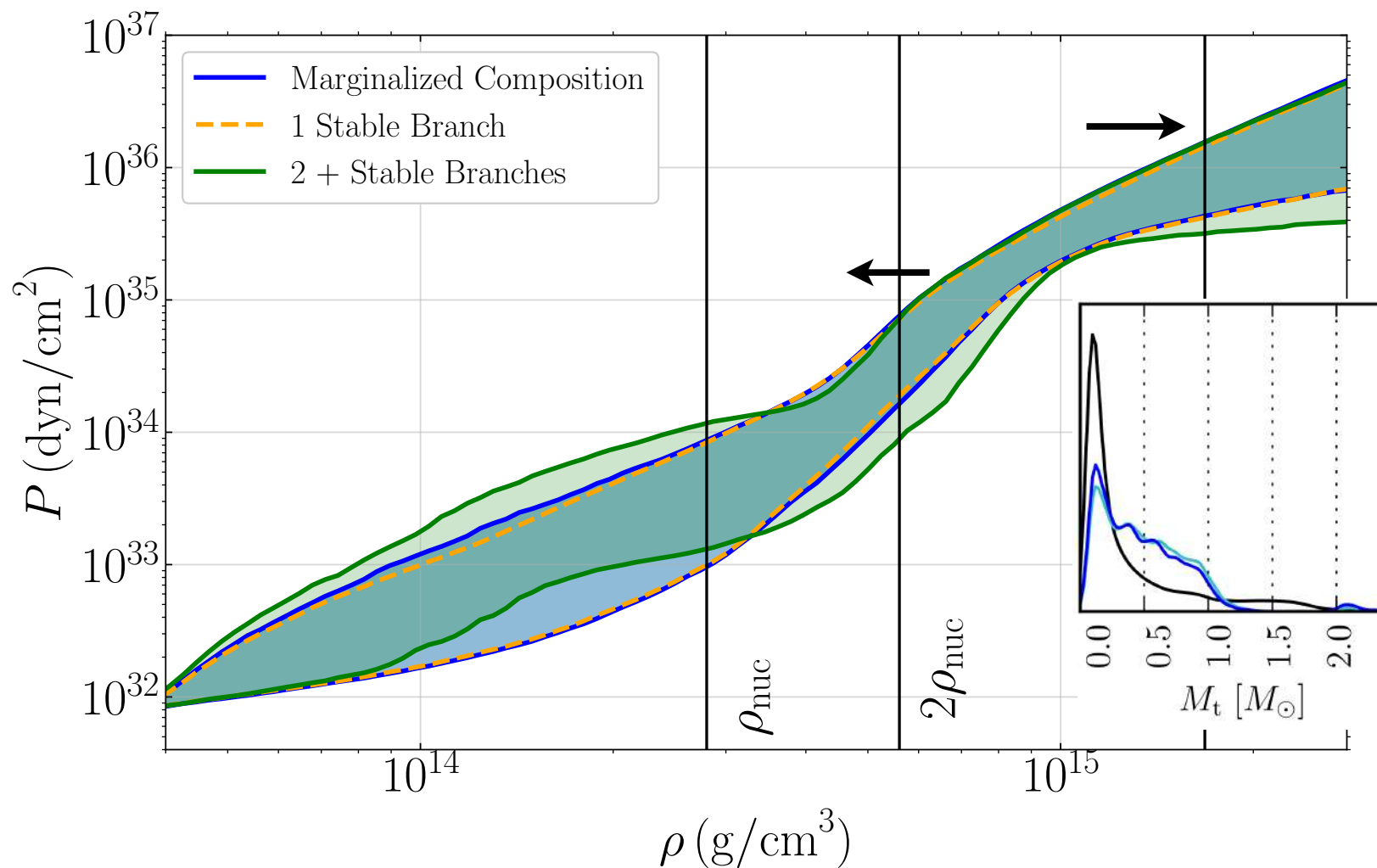
- nonparametric survey conditioned on ensembles of existing model EoSs
- GW170817+190425, NICER J0030 & J0740, and massive pulsars

Pressure vs. density



- tightening the pressure constraint at intermediate densities
- (90% symmetric credible intervals) best compatibility with data

Single branch vs. multiple branches



driven by radius

expected from
max. mass

- full posterior is dominated by EoSs with a single stable branch
- onset for the unstable branch (extra softening) pushed to two ends

Summary

- multimessenger constraints point to NS radii around 12.5 km \pm 1.5 km
- most extreme phase transitions that lead to drastic >2 -3km reduction seem disfavored; onset restricted to either low or high densities
- milder PTs or smooth crossovers are fairly consistent with data; requires high sound speed in the inner core
- pressure or stiffness in nuclear EoS up to twice saturation density is crucial for interpretations of high-density behavior: the golden window

Looking forward

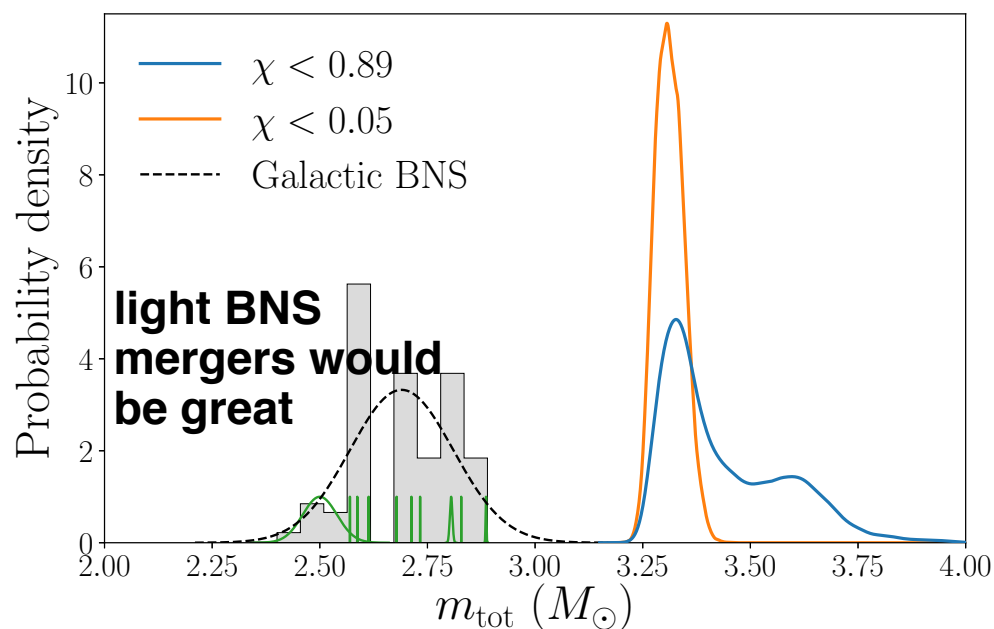
LVC collaboration
arXiv:2006.12611

GW190425

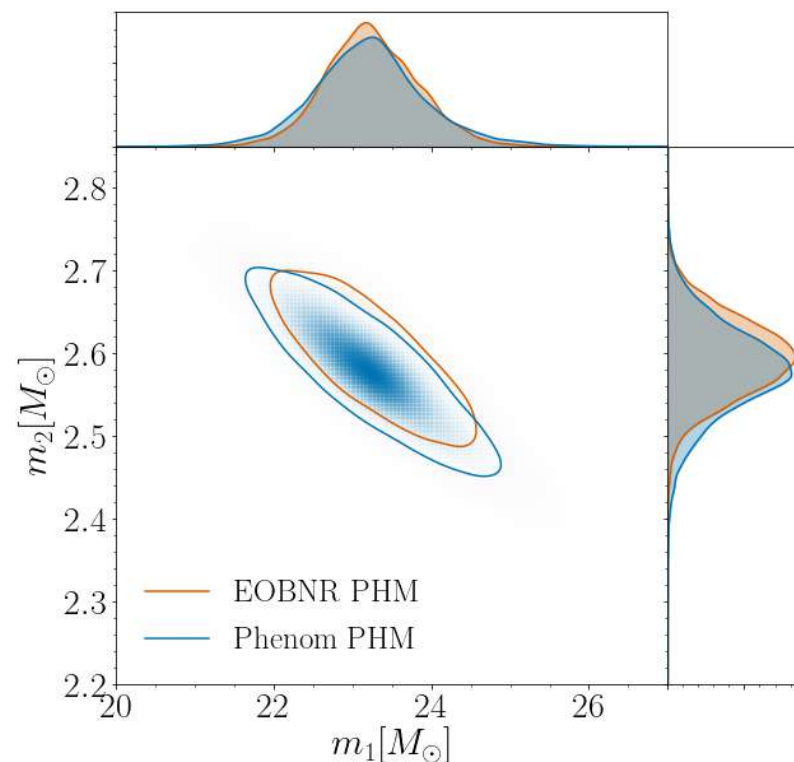
hunting for surprises..

- total mass ~ 3.4 solar masses
- signal too weak to provide further EoS constraints $R < 16$ km

LVC collaboration
arXiv:2001.01761



see events of GWTC-2: arXiv:2010.14527



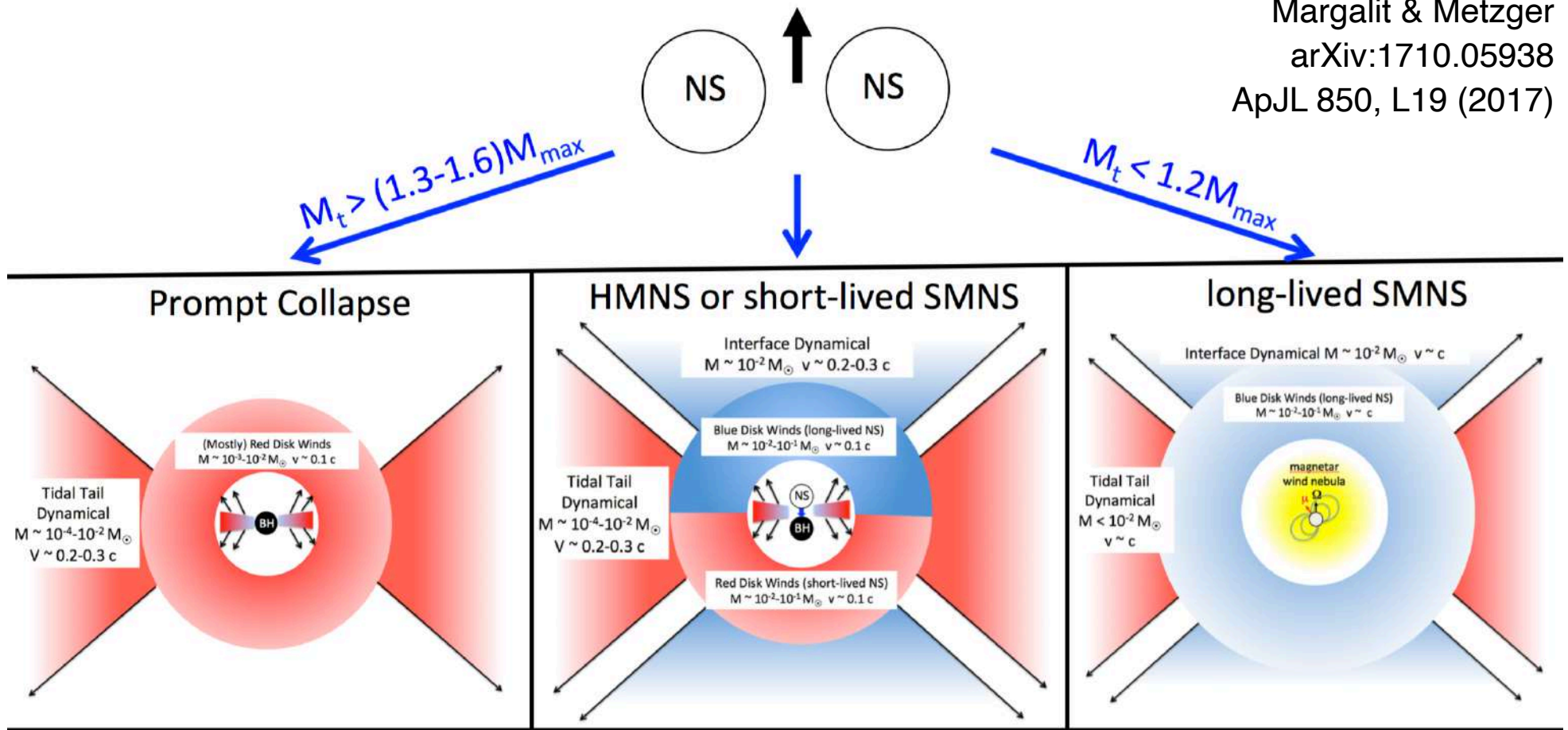
GW190814

more mass-gap objects?

- component of ambiguous nature
- most asymmetric system observed

Fate of merger remnant

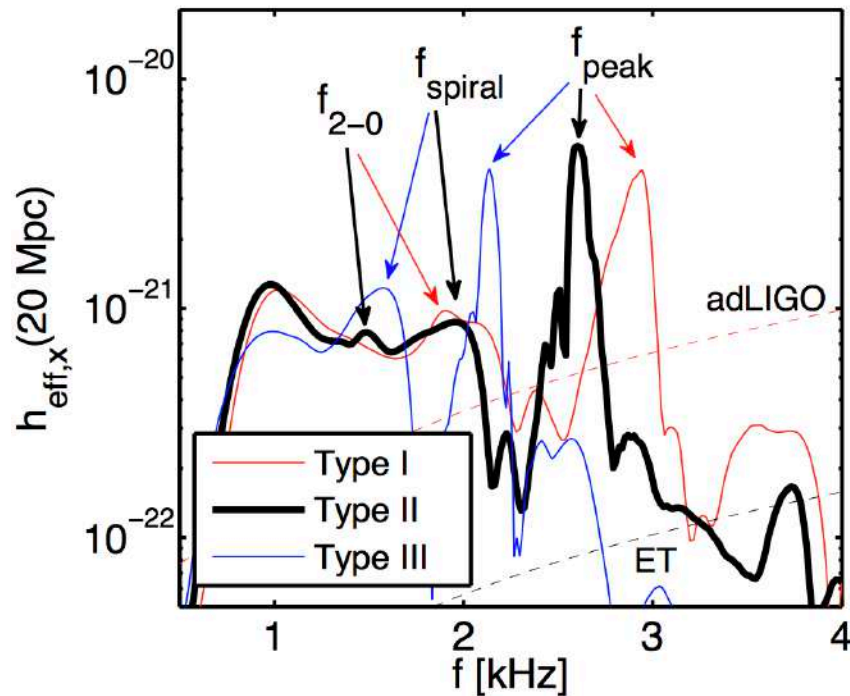
Margalit & Metzger
arXiv:1710.05938
ApJL 850, L19 (2017)



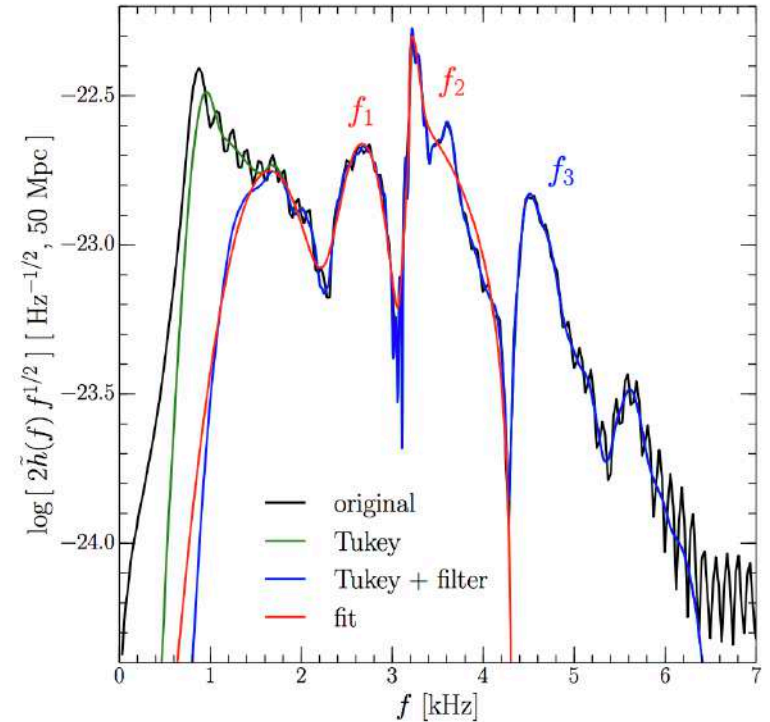
- GW + EM constraints from 170817 seem to favor $M_{\text{max}} < 2.16 \sim 2.3$ solar masses Ruiz et al. (2018), Rezzolla et al. (2018), Shibata et al. (2019)
- radius > 10.68 km to prevent prompt collapse Bauswein et al. (2017)

Post-merger dynamics

Bauswein & Stergioulas, arXiv:1502.03176



Takami et al., arXiv:1412.3240

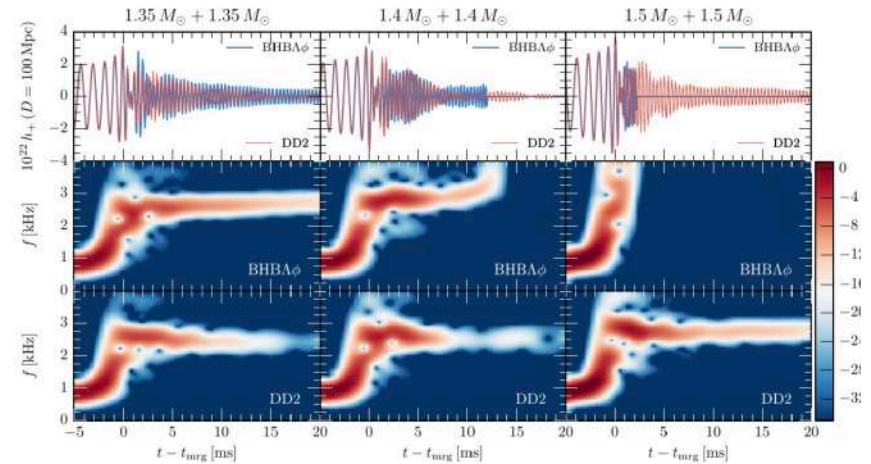
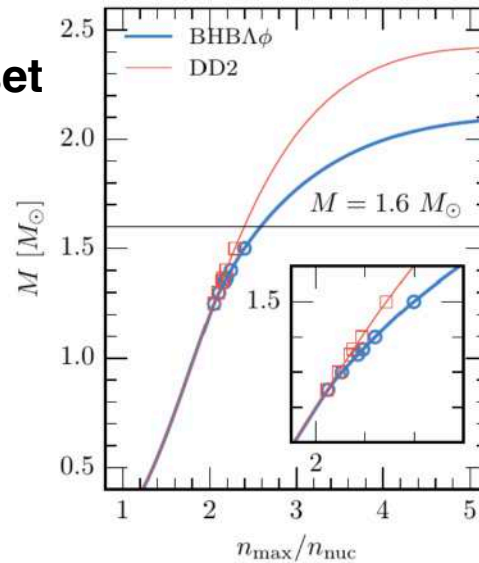


- complicated spectra of excited modes depend on the EoS
- location of the dominant peak strongly correlated with NS radii
- within reach of next generation GW detectors (~ 10 times more sensitive)

e.g. softening effects on post-merger GW

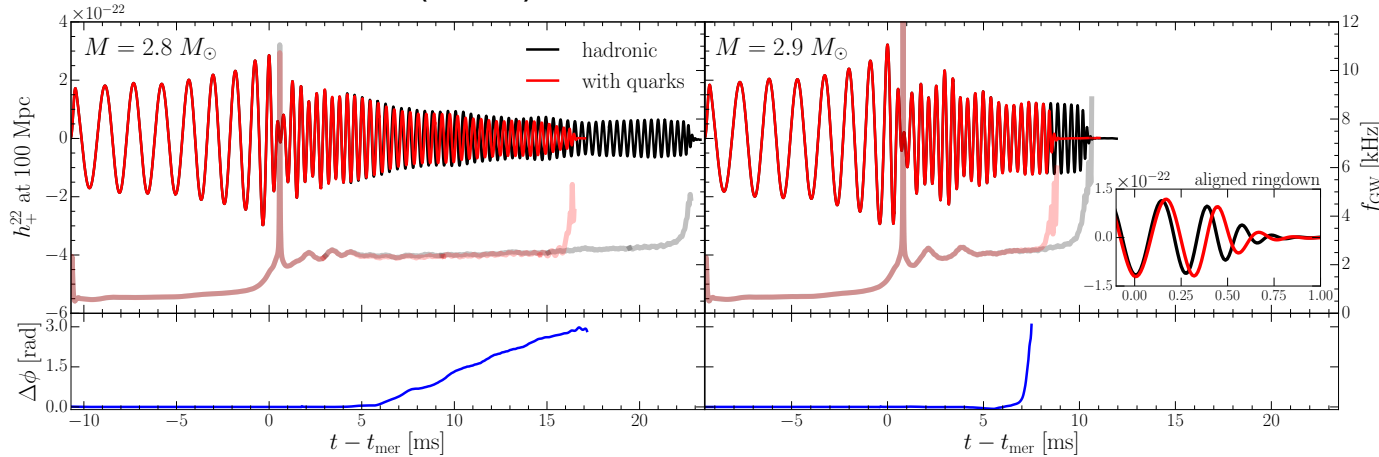
hyperon onset

- more compact remnant (higher central density)
- earlier collapse; higher frequency



Most et al. arXiv:1807.03684
PRL 122, 061101 (2019)

Radice et al. arXiv:1612.06429
ApJL 842, L10 (2017)

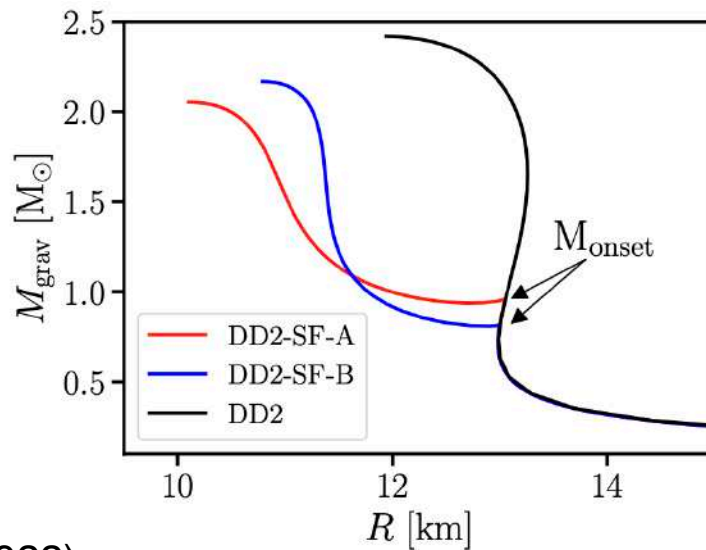


**1st-OPT to soft
quark matter
after merger**

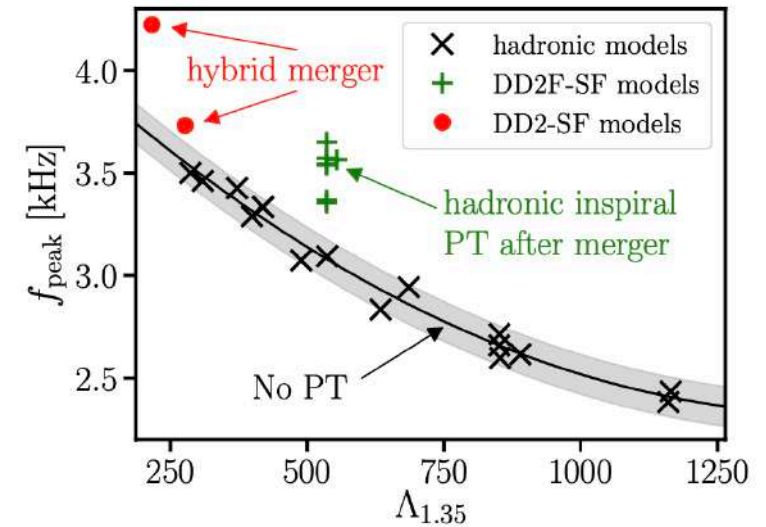
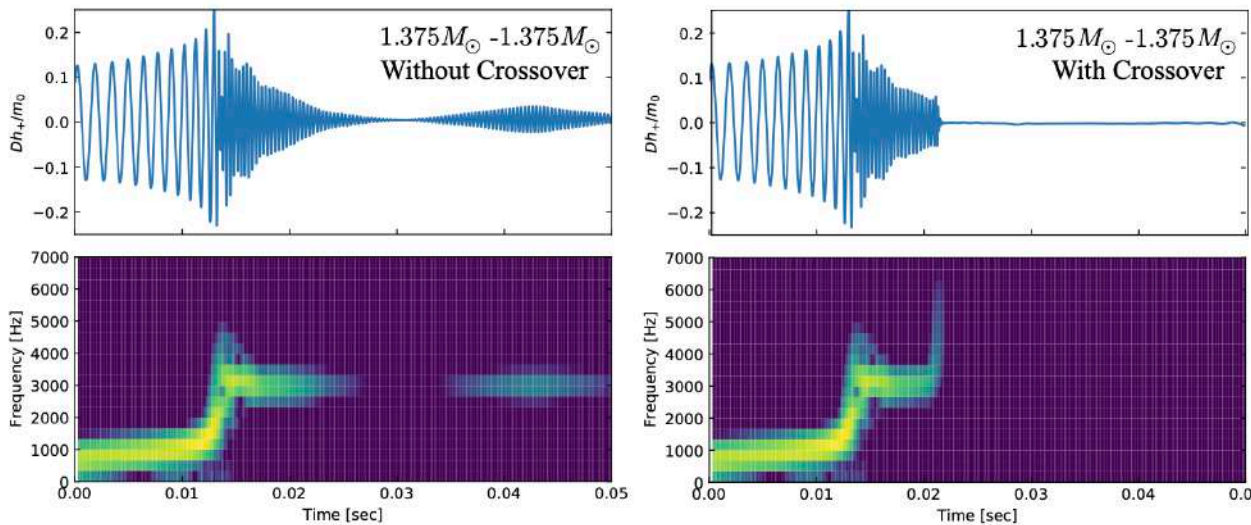
e.g. softening effects on post-merger GW

third-family stars

- stiff EoS at low density -DD2
- strong 1st-OPT to stiff quark matter



Fujimoto et al. (2022)



Bauswein & Blacker (2020)

- soft EoS at low density \sim N3LO chiEFT

- rapid stiffening in the crossover regime

crossover to soft quark matter after merger

NSBH mergers

LVK collaboration
arXiv:2106.15163

THE ASTROPHYSICAL JOURNAL LETTERS, 915:L5 (24pp), 2021 July 1

© 2021. The Author(s). Published by the American Astronomical Society.

OPEN ACCESS

<https://doi.org/10.3847/2041-8213/ac082e>



CrossMark

Observation of Gravitational Waves from Two Neutron Star–Black Hole Coalescences

Table 2
Source Properties of GW200105 and GW200115

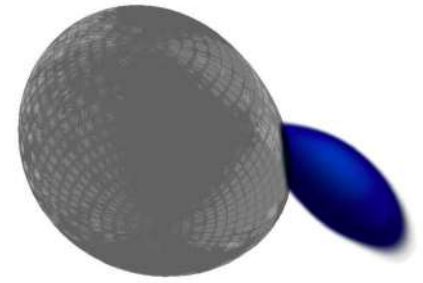
	GW200105		GW200115	
	Low Spin ($\chi_2 < 0.05$)	High Spin ($\chi_2 < 0.99$)	Low Spin ($\chi_2 < 0.05$)	High Spin ($\chi_2 < 0.99$)
Primary mass m_1/M_\odot	$8.9^{+1.1}_{-1.3}$	$8.9^{+1.2}_{-1.5}$	$5.9^{+1.4}_{-2.1}$	$5.7^{+1.8}_{-2.1}$
Secondary mass m_2/M_\odot	$1.9^{+0.2}_{-0.2}$	$1.9^{+0.3}_{-0.2}$	$1.4^{+0.6}_{-0.2}$	$1.5^{+0.7}_{-0.3}$
Mass ratio q	$0.21^{+0.06}_{-0.04}$	$0.22^{+0.08}_{-0.04}$	$0.24^{+0.31}_{-0.08}$	$0.26^{+0.35}_{-0.10}$
Total mass M/M_\odot	$10.8^{+0.9}_{-1.0}$	$10.9^{+1.1}_{-1.2}$	$7.3^{+1.2}_{-1.5}$	$7.1^{+1.5}_{-1.4}$
Chirp mass \mathcal{M}/M_\odot	$3.41^{+0.08}_{-0.07}$	$3.41^{+0.08}_{-0.07}$	$2.42^{+0.05}_{-0.07}$	$2.42^{+0.05}_{-0.07}$
Detector-frame chirp mass $(1+z)\mathcal{M}/M_\odot$	$3.619^{+0.006}_{-0.006}$	$3.619^{+0.007}_{-0.008}$	$2.580^{+0.006}_{-0.007}$	$2.579^{+0.007}_{-0.007}$
Primary spin magnitude χ_1	$0.09^{+0.18}_{-0.08}$	$0.08^{+0.22}_{-0.08}$	$0.31^{+0.52}_{-0.29}$	$0.33^{+0.48}_{-0.29}$
Effective inspiral spin parameter χ_{eff}	$-0.01^{+0.08}_{-0.12}$	$-0.01^{+0.11}_{-0.15}$	$-0.14^{+0.17}_{-0.34}$	$-0.19^{+0.23}_{-0.35}$
Effective precession spin parameter χ_p	$0.07^{+0.15}_{-0.06}$	$0.09^{+0.14}_{-0.07}$	$0.19^{+0.28}_{-0.17}$	$0.21^{+0.30}_{-0.17}$
Luminosity distance D_L/Mpc	280^{+110}_{-110}	280^{+110}_{-110}	310^{+150}_{-110}	300^{+150}_{-100}
Source redshift z	$0.06^{+0.02}_{-0.02}$	$0.06^{+0.02}_{-0.02}$	$0.07^{+0.03}_{-0.02}$	$0.07^{+0.03}_{-0.02}$

no information on matter effects
no significant EM detections

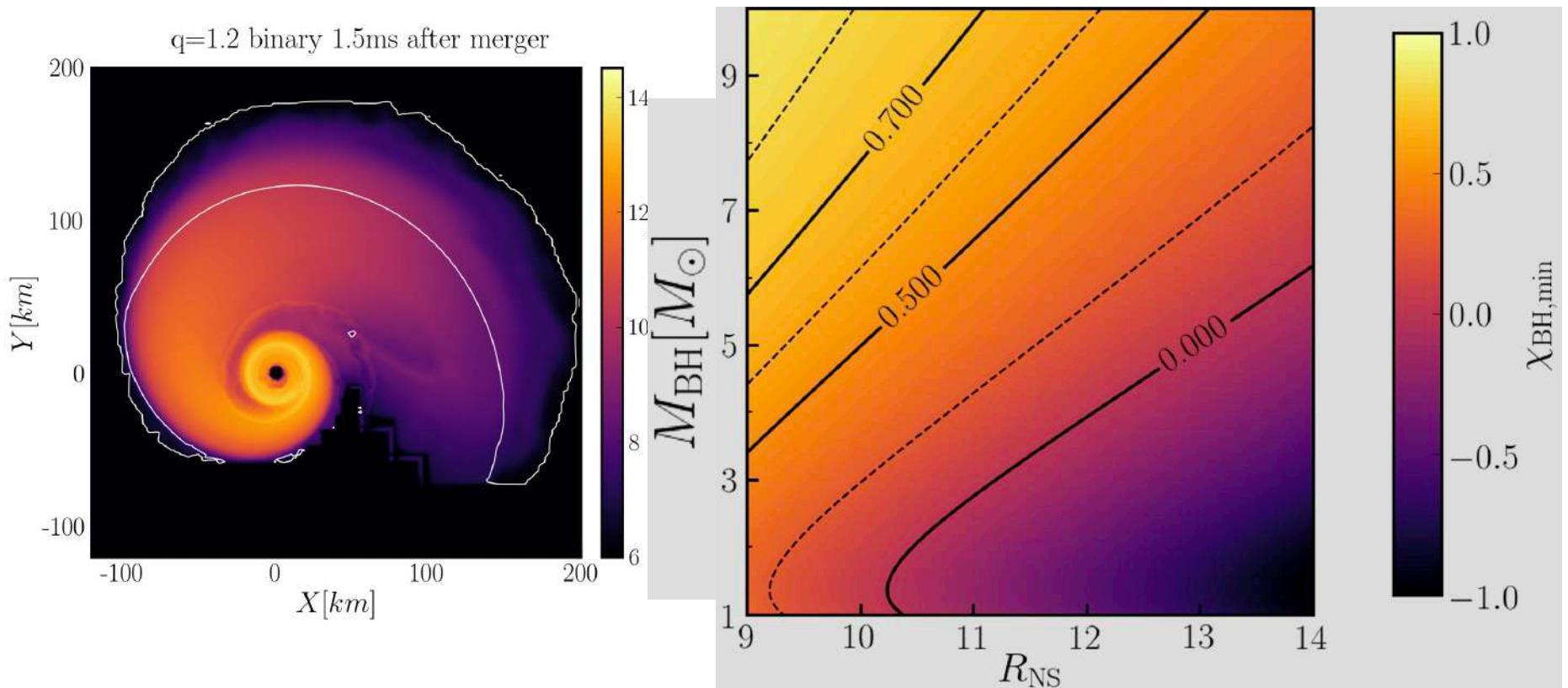
- GW200105: $\sim 1.9 + \sim 9$ solar masses
- GW200115: $\sim 1.5 + \sim 6$ solar masses

see events of GWTC-3: arXiv:2111.03606

Outcome of a NSBH merger

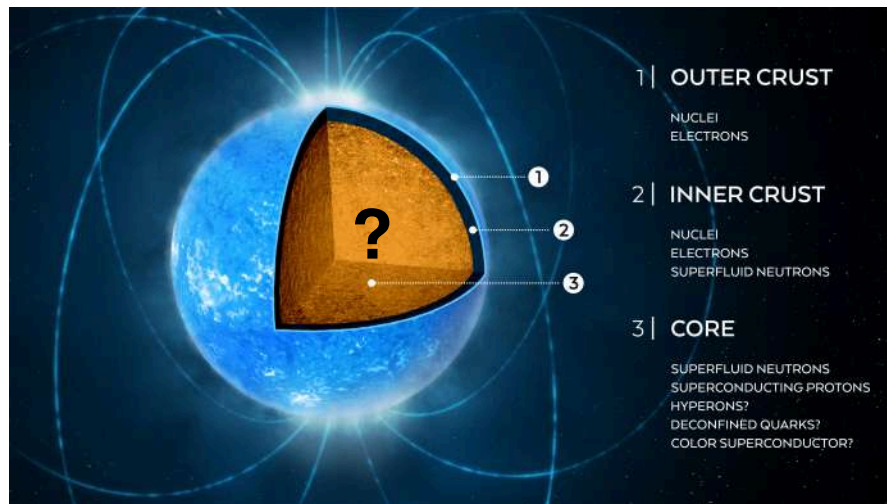
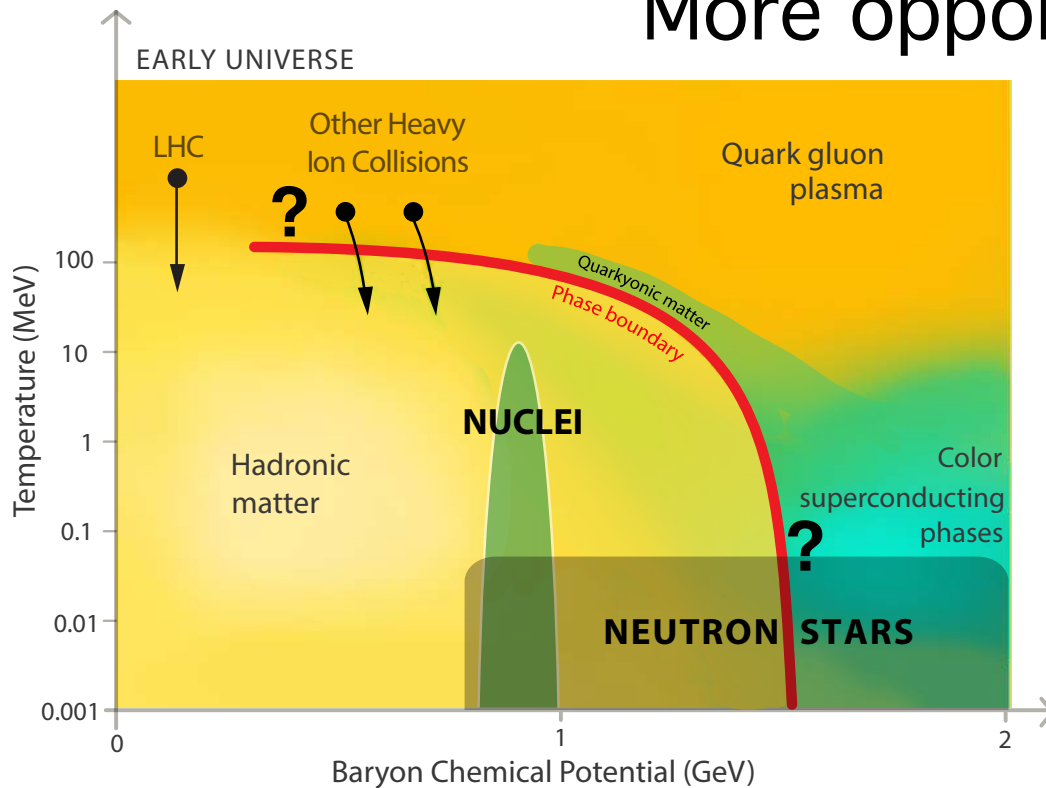


Foucart et al. (2018)



- NS is either tidally disrupted or plunges into the BH - mass ratio, spin, EoS
- radius determines if tides are **measurable** & if **EM** signals can be produced

More opportunities



- probing dense matter in NSs
- cooling of NS 1987A - neutrino emissivity, stellar superfluids **(nuclear theory, condensed matter)** ^{new!}
 - merger evolution and astro/GW signals - out-of-equilibrium physics; composition details **(simulation, nucleosynthesis)**
 - next Galactic supernova? **(neutrino physics)**
 - asteroseismology **(hydrodynamics, GR, nucl-th)**
 - ...and more - add your own!

Rev. Mod. Phys.
88, 021001 (2016)

THANK YOU!

Q & A