



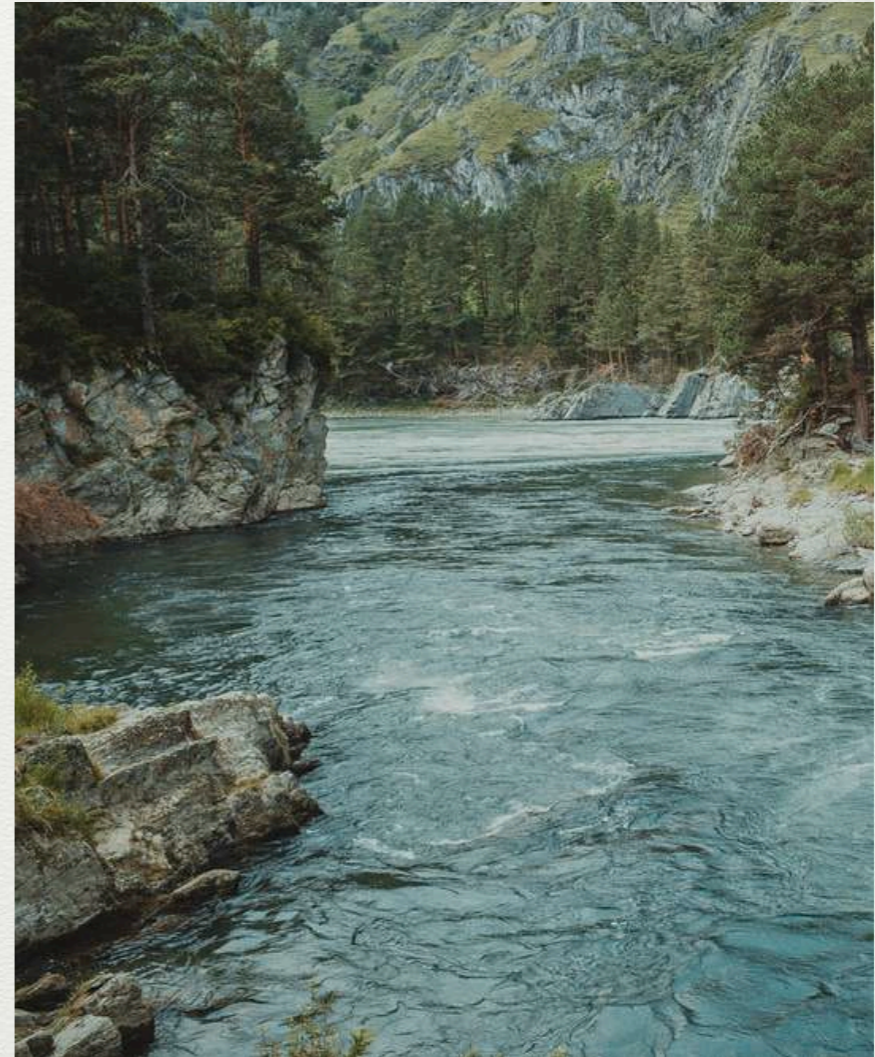
Fluid dynamics in the extreme - The Quark-Gluon Plasma

Jacquelyn Noronha-Hostler
University of Illinois Urbana-Champaign

Theoretical Physics Colloquium via Arizona State University

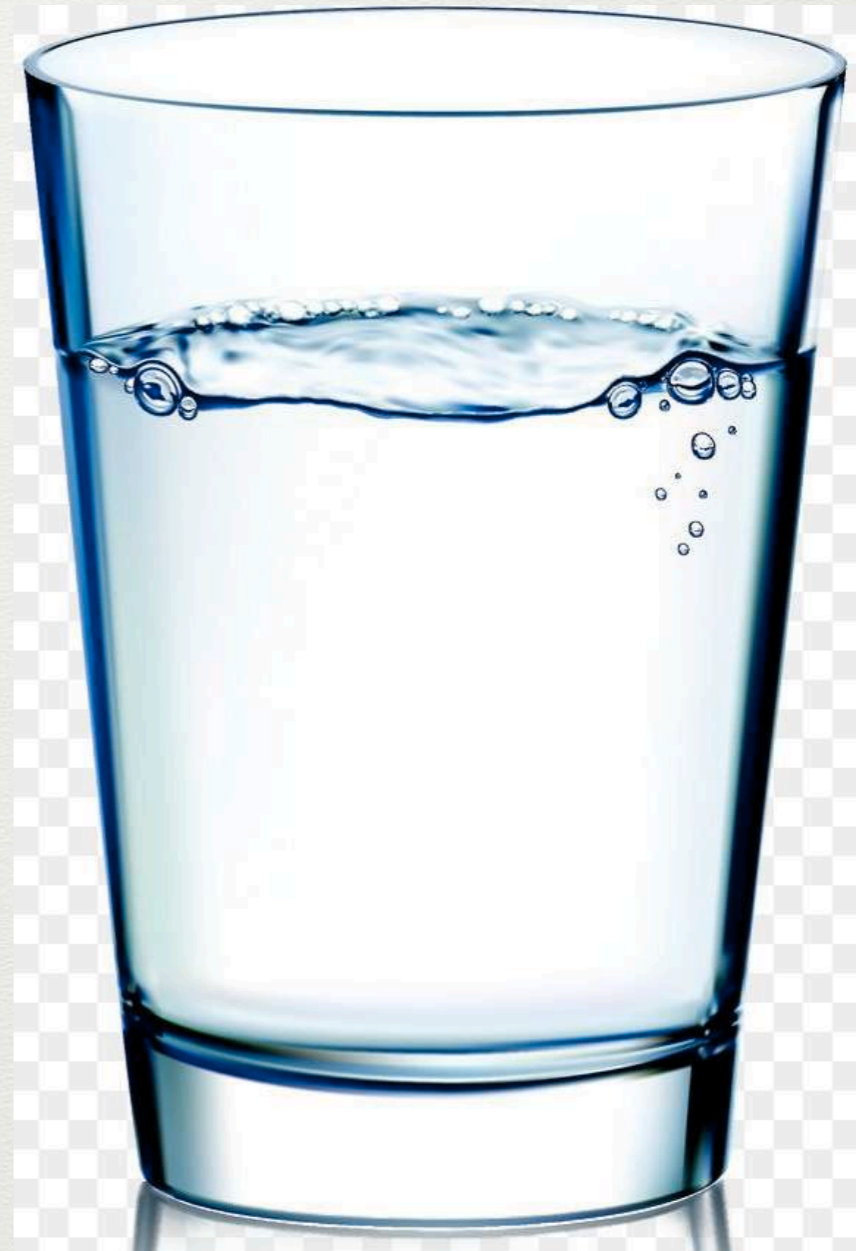
Fluids 101

- It flows (its particles easily move past each other)
- Takes the shape of the container (no permanent shape)
- Cannot resist an outside shearing force
- Variables: $\{\rho, P\}$, not $\{m, F\}$



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When is fluid dynamics applicable?

Large separation of scales i.e. small Knudsen (or inverse Reynolds) number

$$Kn \sim \frac{\text{Small scale}^* (\text{H}_2\text{O molecule})}{\text{Large scale (size of lake)}}$$

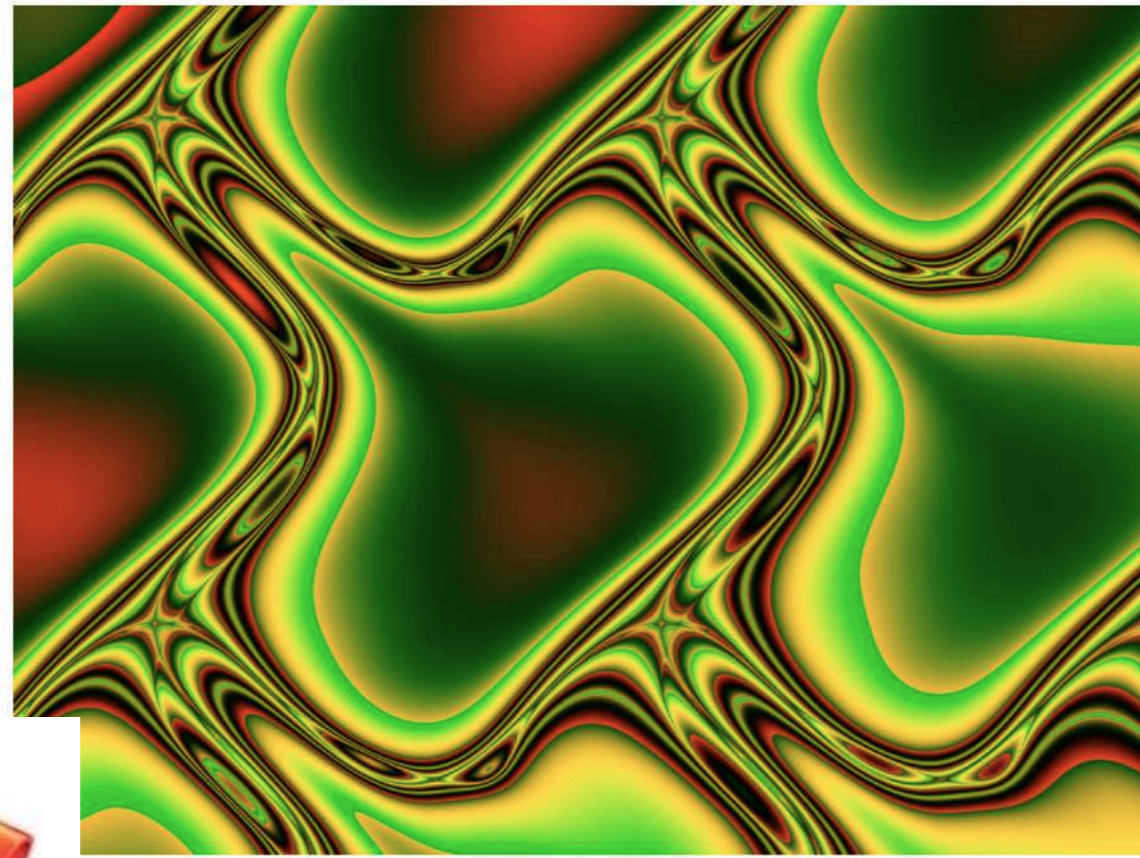
* mean free path i.e. distance before the molecule collides with something else

Question: When can you apply fluid dynamics?

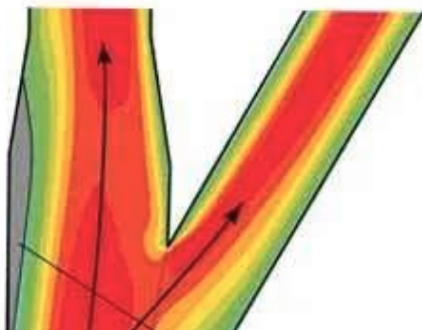
Answer: $Kn \ll 1$

Fluids are everywhere

Traffic Jam



Blood flow



Neutron Star Mergers



Fluids at the **extreme**

What happens when a fluid moves at the speed of light?

New equations of motion (Israel-Stewart) are needed to preserve casualty and stability



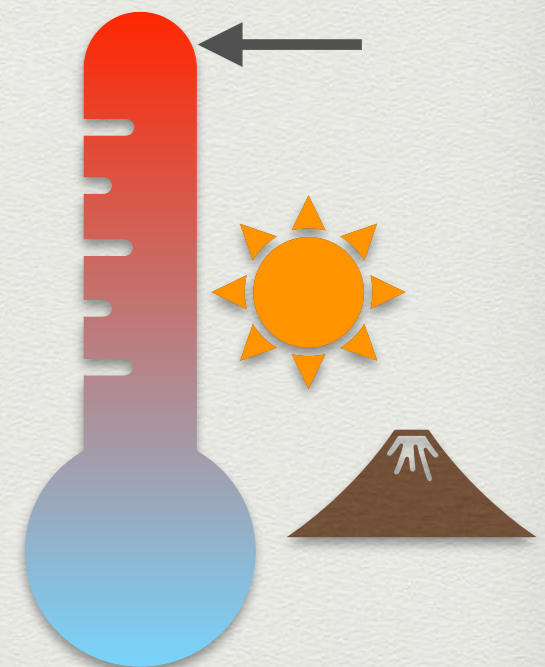
What happens when fluids are heated up to the highest temperatures possible on Earth?

The degrees of freedom are deconfined quarks and gluons

$10^{12} K$

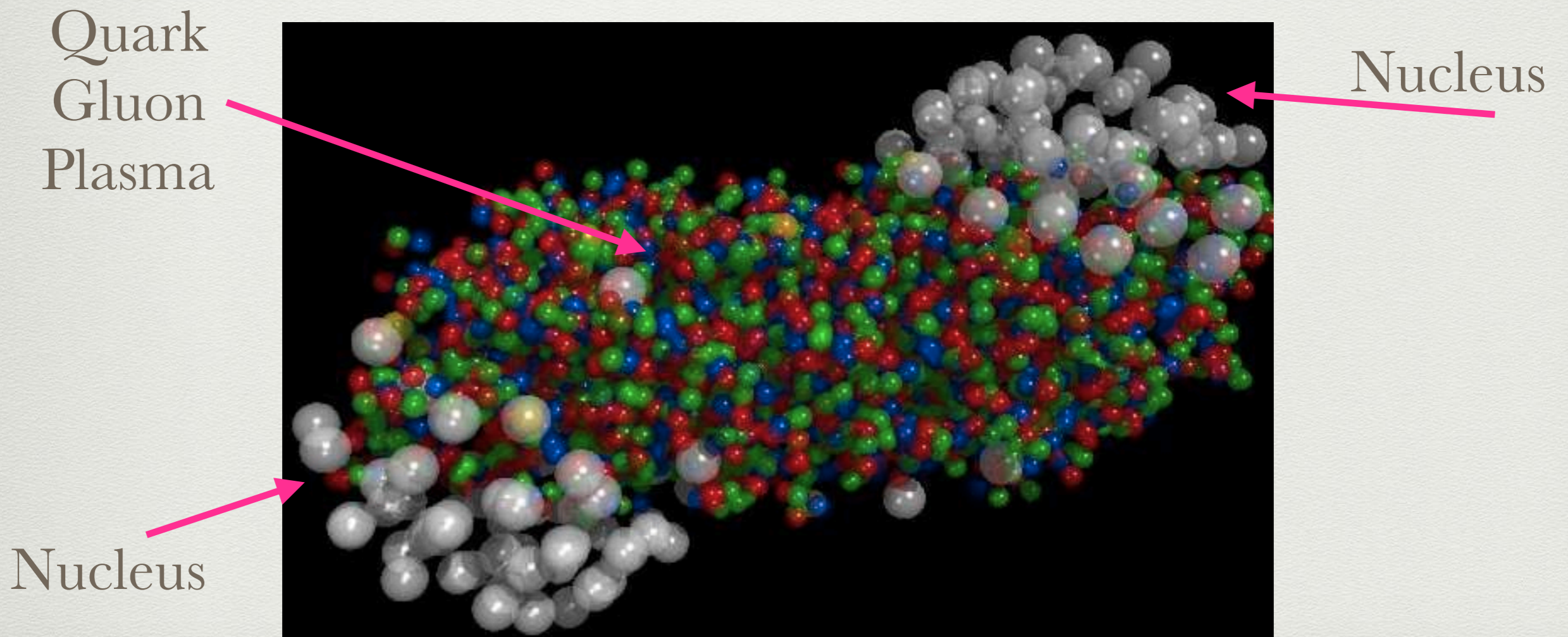
What is the smallest fluid?

We're still figuring that out, but Knudsen numbers get tricky



Why not all three?

The **Quark Gluon Plasma** is created using the highest temperatures on Earth, in the smallest systems possible (colliding nuclei, maybe even colliding protons), and flows at ultra relativistic speeds



The Quark
Gluon
Plasma is
the ...

Hottest



Most Perfect

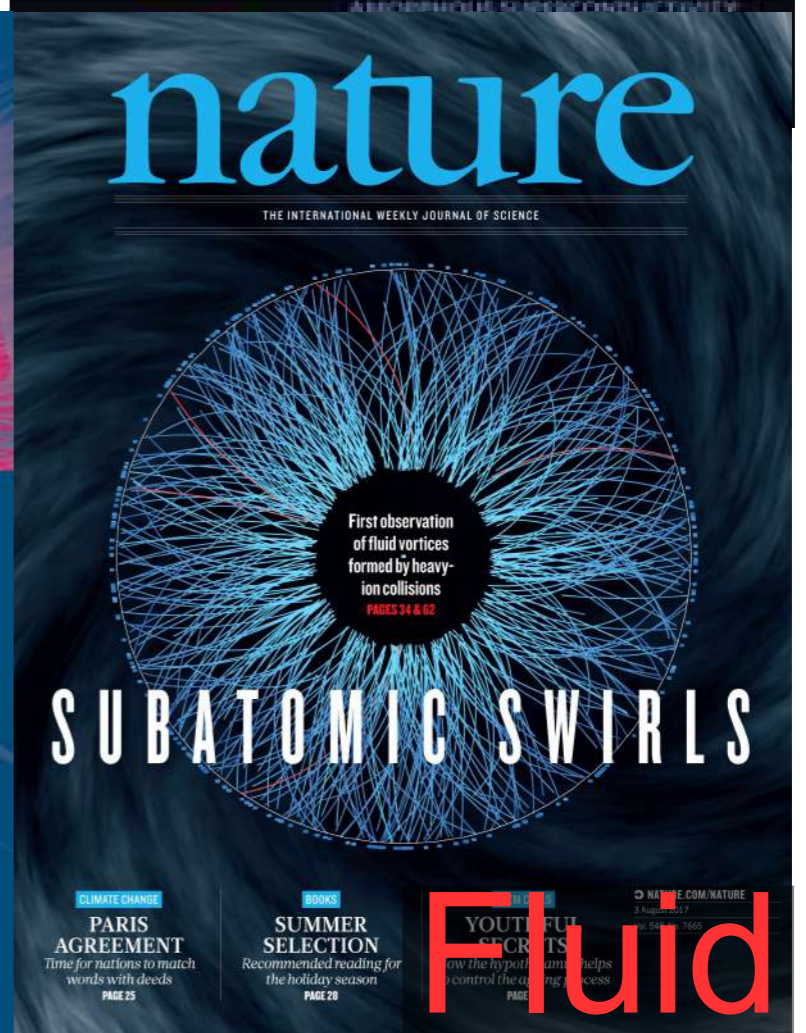
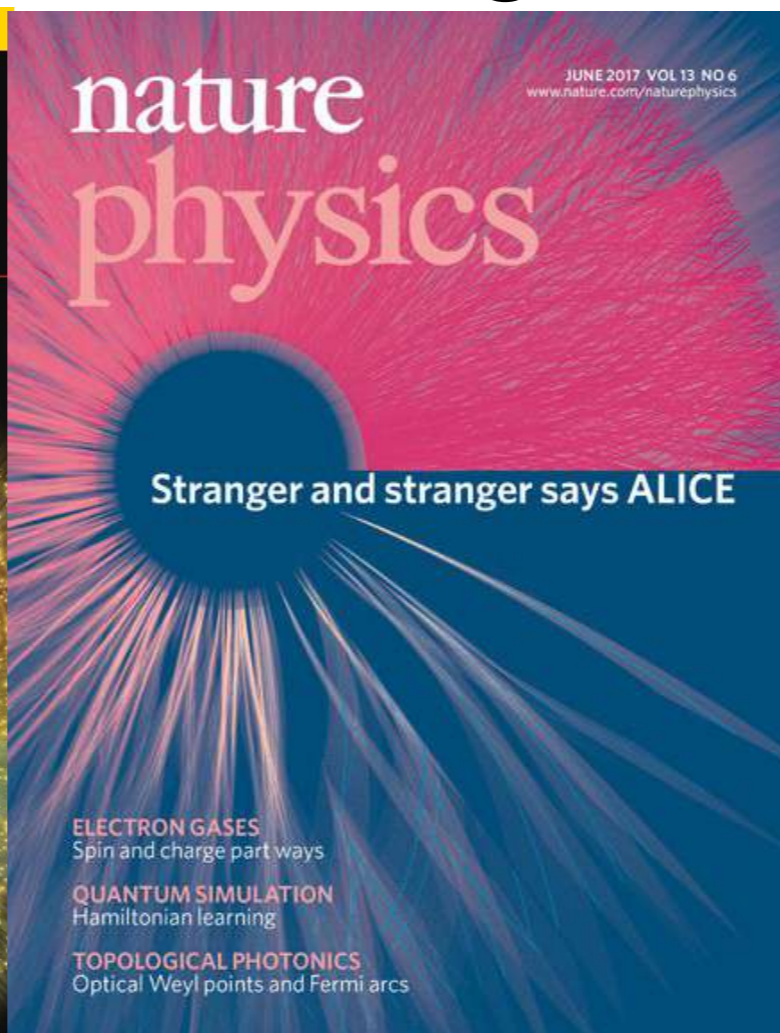
Strange

Smallest

nature
physics

The geometry of a
quark-gluon plasma

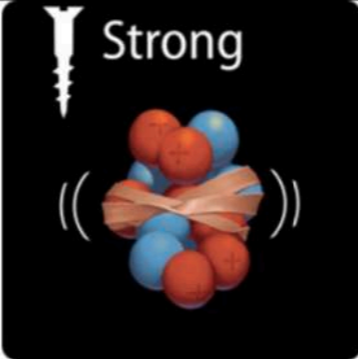
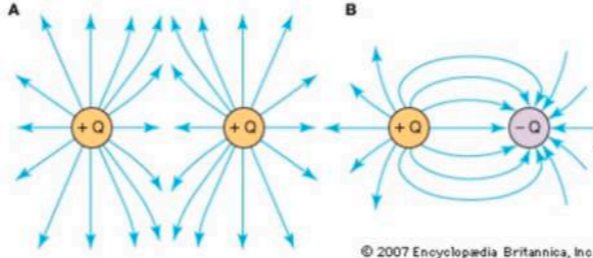
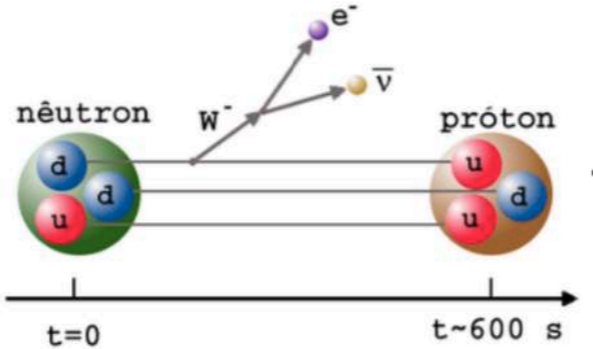
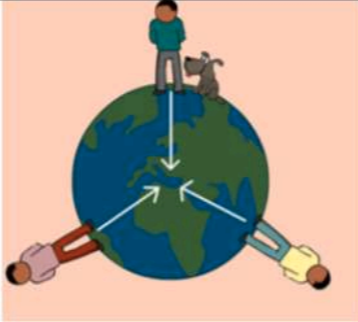
Most Vortical



Fluid

To understand the Quark Gluon Plasma, we
first need to understand the strong force and
Quantum Chromodynamics

Strongest Force

Force		Strength	Range (m)
Strong	 <p>Strong</p>	1	10^{-15} (size of nucleus)
Electromagnetic	 <p>A B</p>	$\frac{1}{137}$	∞
Weak	 <p>nêutron</p> <p>próton</p> <p>W^-</p> <p>e^-</p> <p>$\bar{\nu}$</p> <p>$t=0$</p> <p>$t \sim 600 \text{ s}$</p>	10^{-6}	10^{-18} (0.1% of the proton)
Gravity		$6 * 10^{-39}$	∞

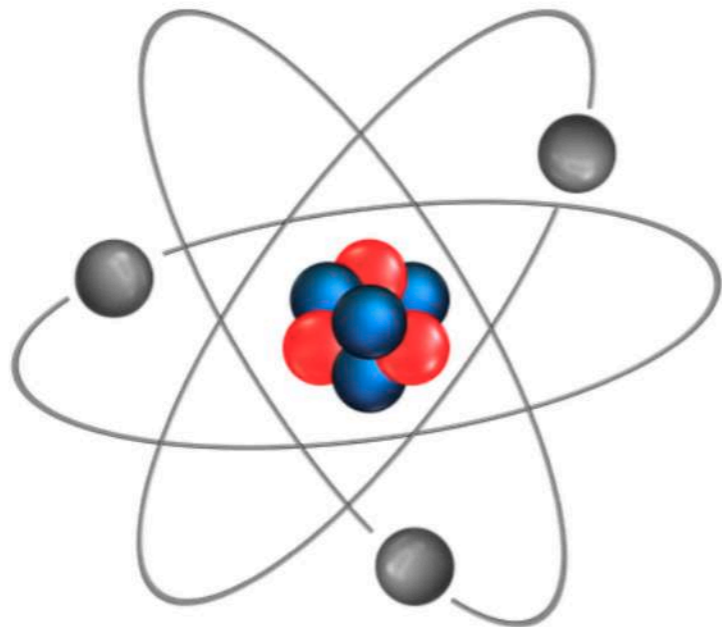
Scales of the universe

Atom

Scale $\sim 10^{-10}$ meter

Discovered by humans
 ~ 1800 's

Experiment Chemistry

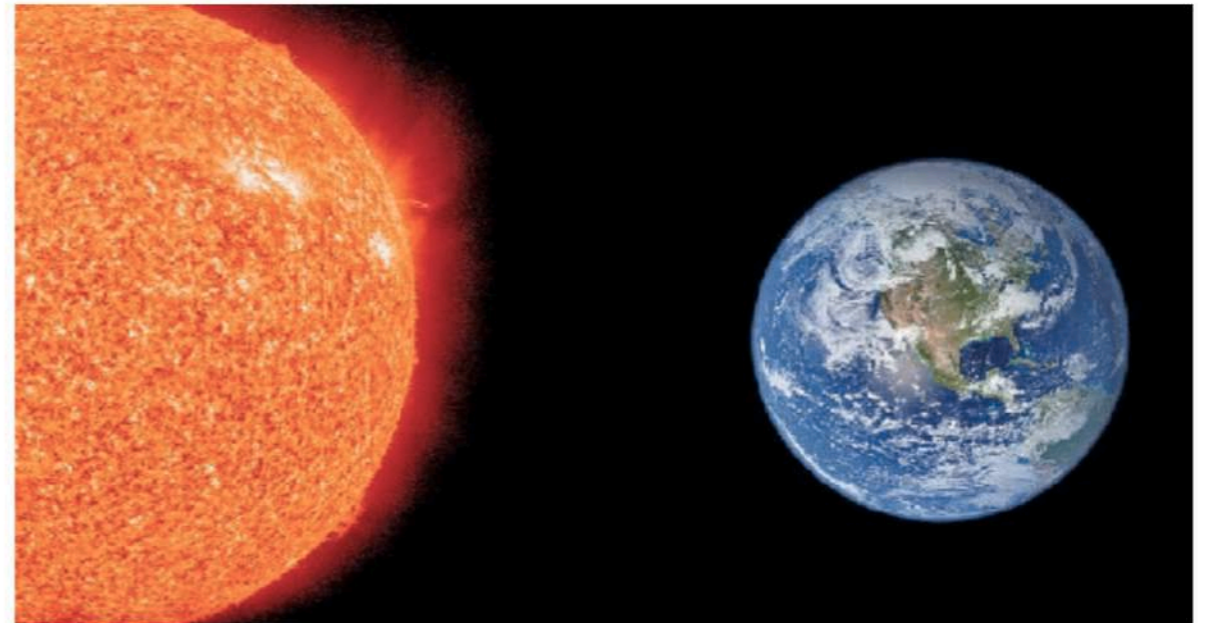


Distance Earth to Sun

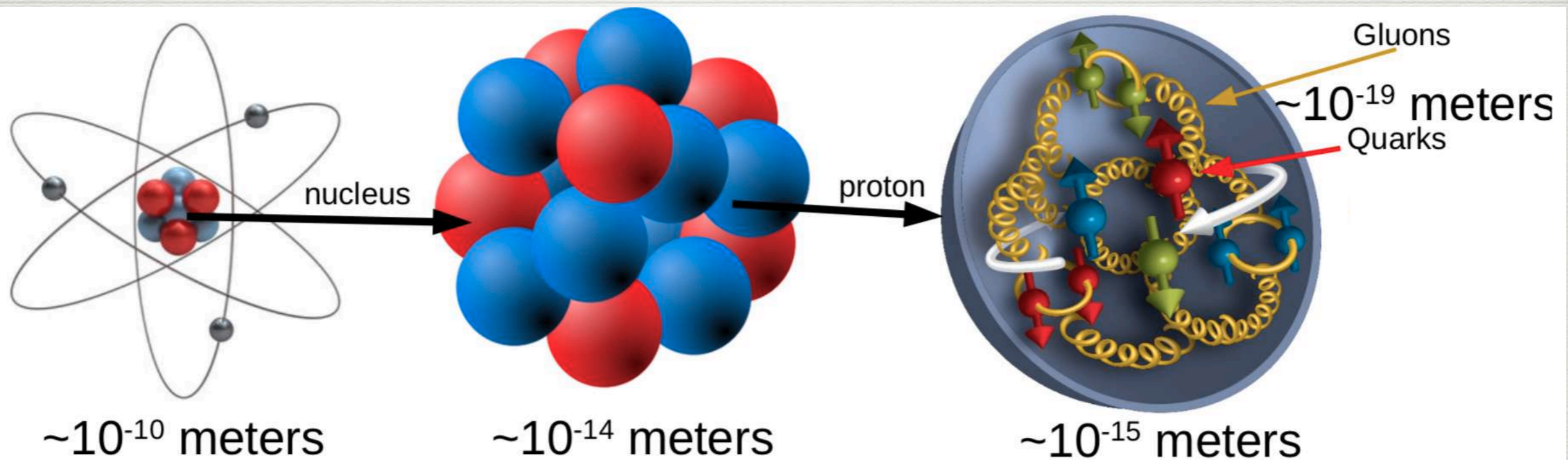
Scale $\sim 10^{11}$ meter

Discovered by humans ~ 250
BC

Experiment various



Scales of the strong force



Distance to nearest star (Alpha Centauri system) $\sim 10^{16}$ m

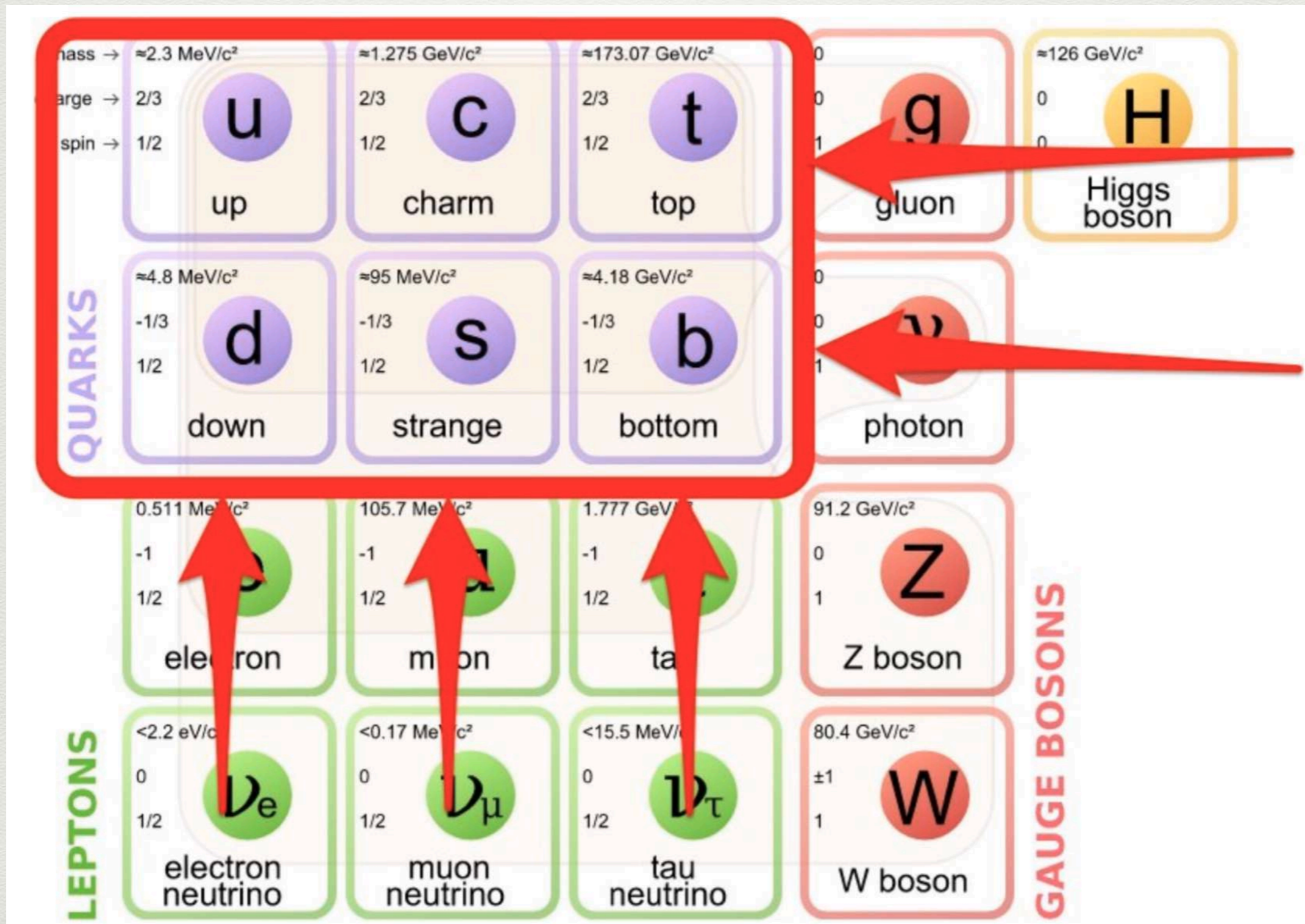
Natural Units: $\hbar = c = 1$ i.e. $E = Mc^2 \rightarrow E = M$

Length 1 Femtometer [fm] = 10^{-15} [m]

Temperature 1000 Megaelectron-volt [MeV] = 11 billion [K]

Mass 1 [MeV] = $1.79 * 10^{-30}$ [kg]

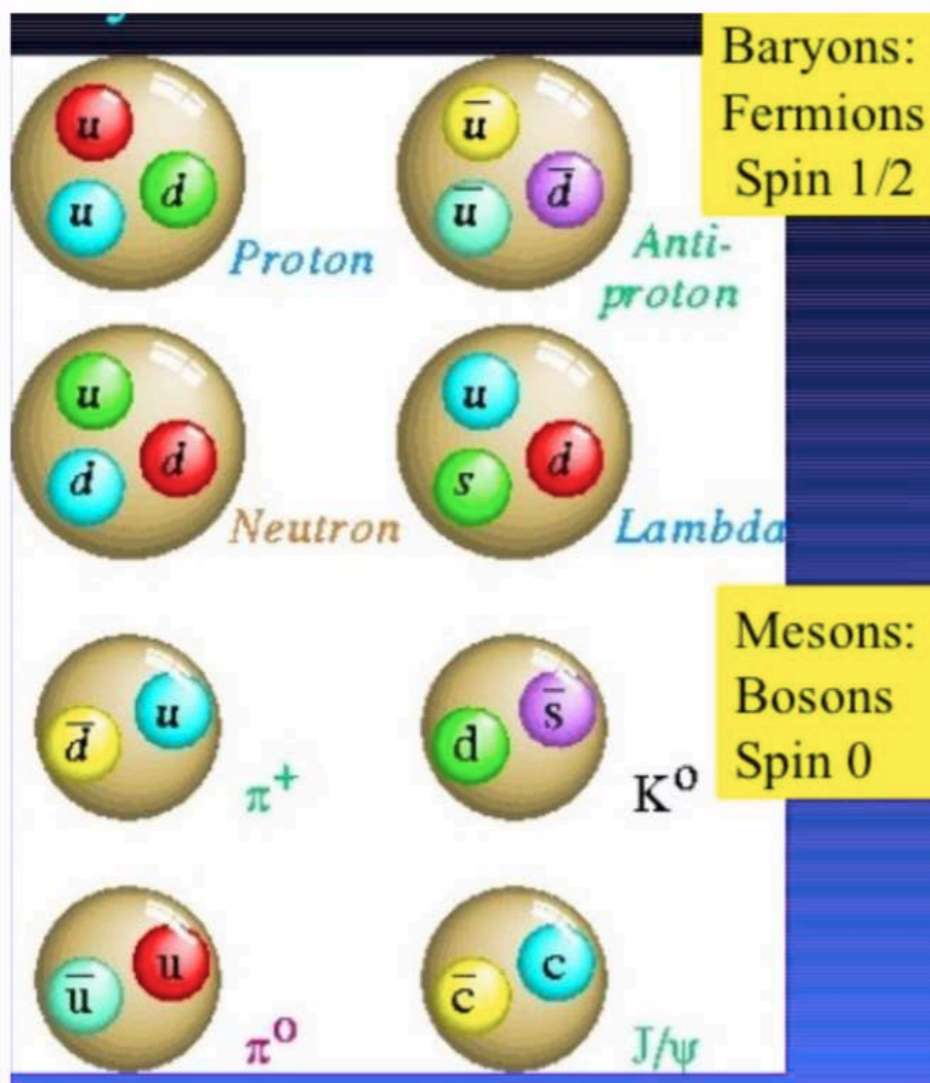
Standard Model



Theory of the strong force: Quantum Chromodynamics (QCD)

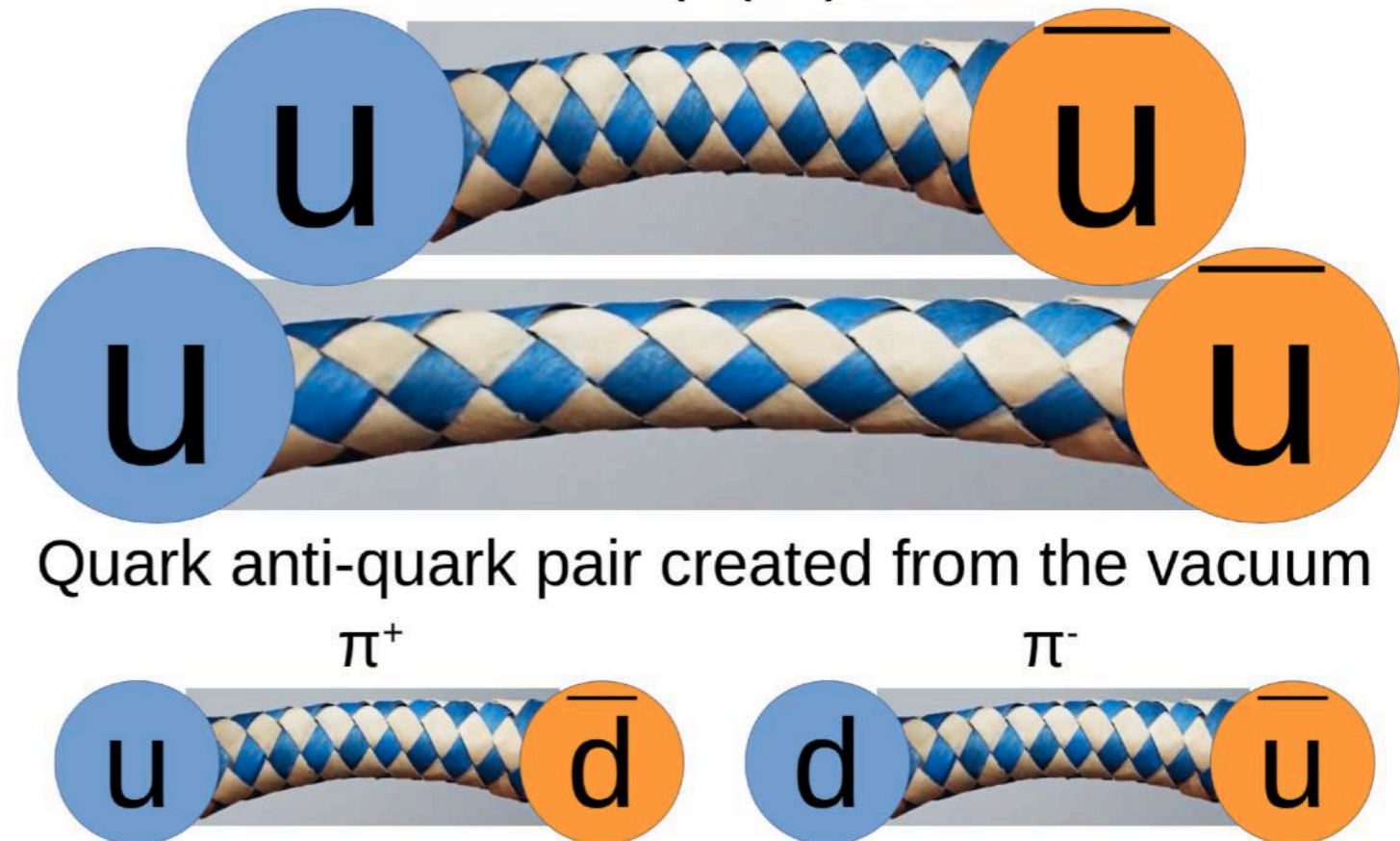
Quantum Chromodynamics (QCD)- David Politzer, Frank Wilczek and David Gross 1973 (2004 Nobel Prize for Physics)

Bound states=Hadrons



Confinement- no free quarks

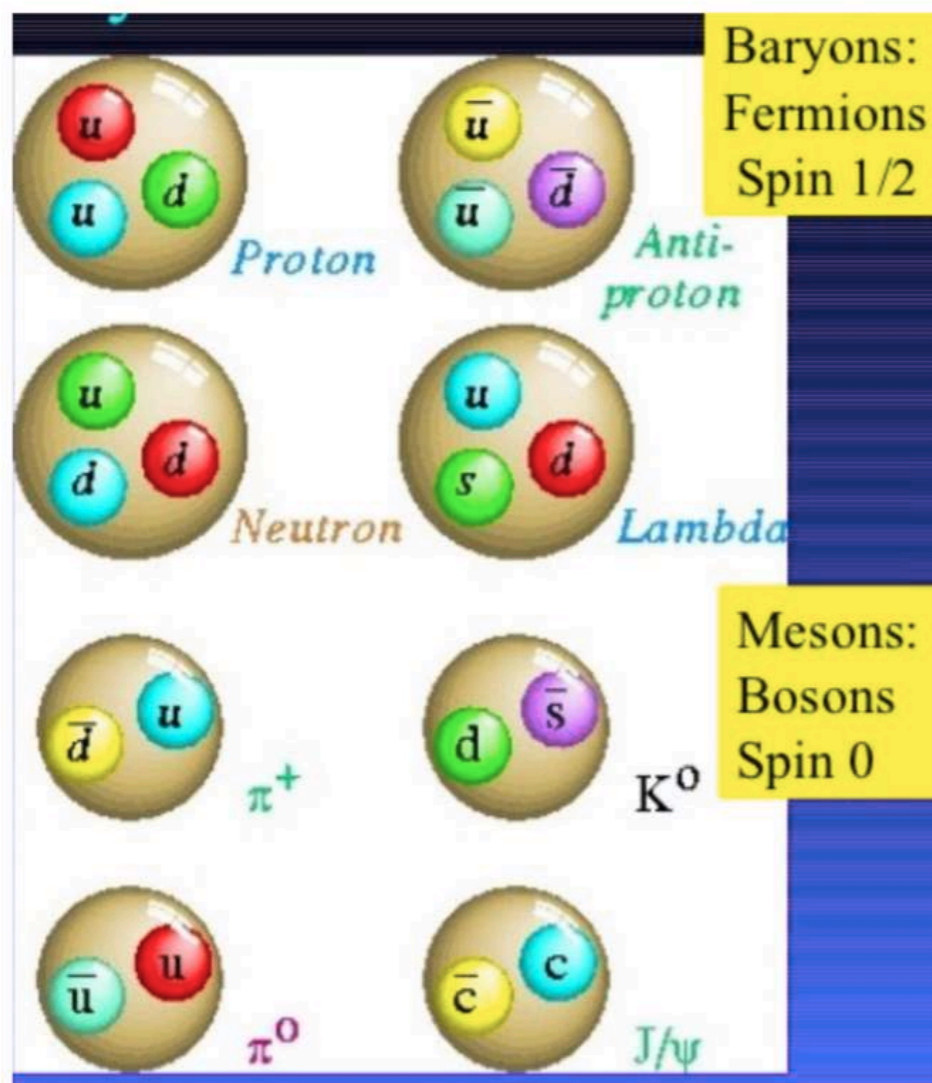
Start with a $\rho^0(u\bar{u})$ meson



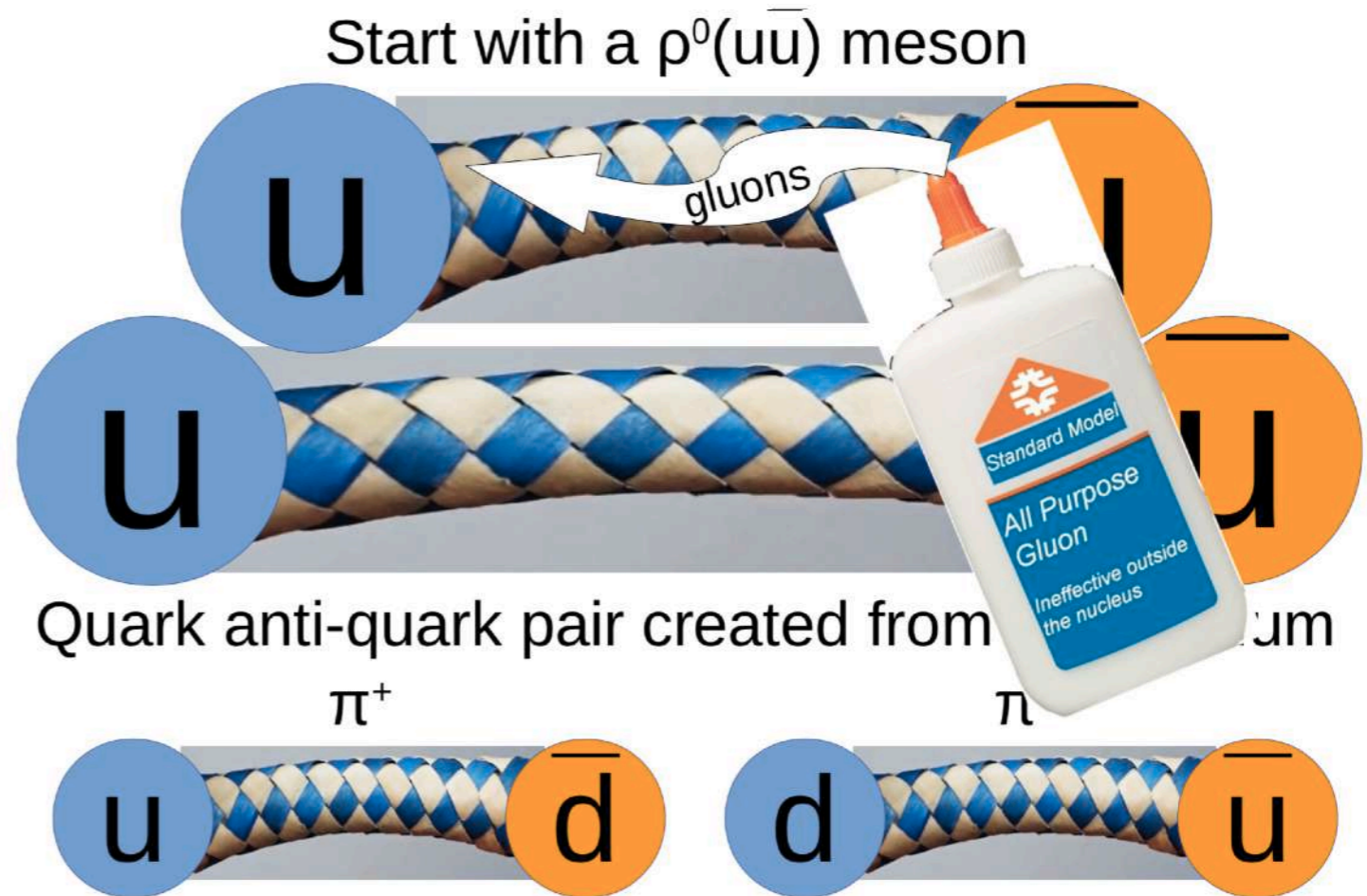
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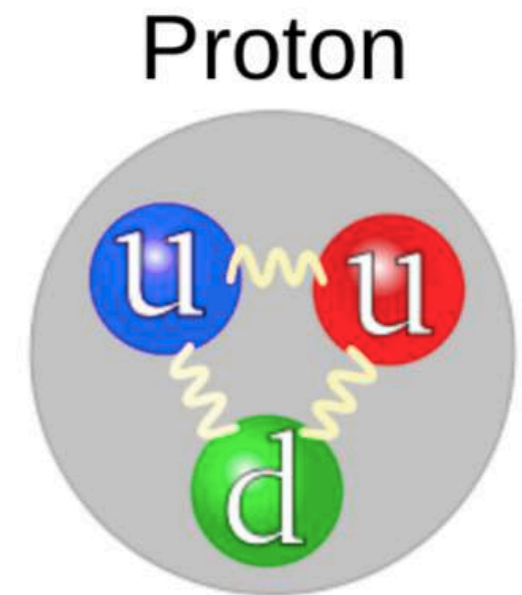
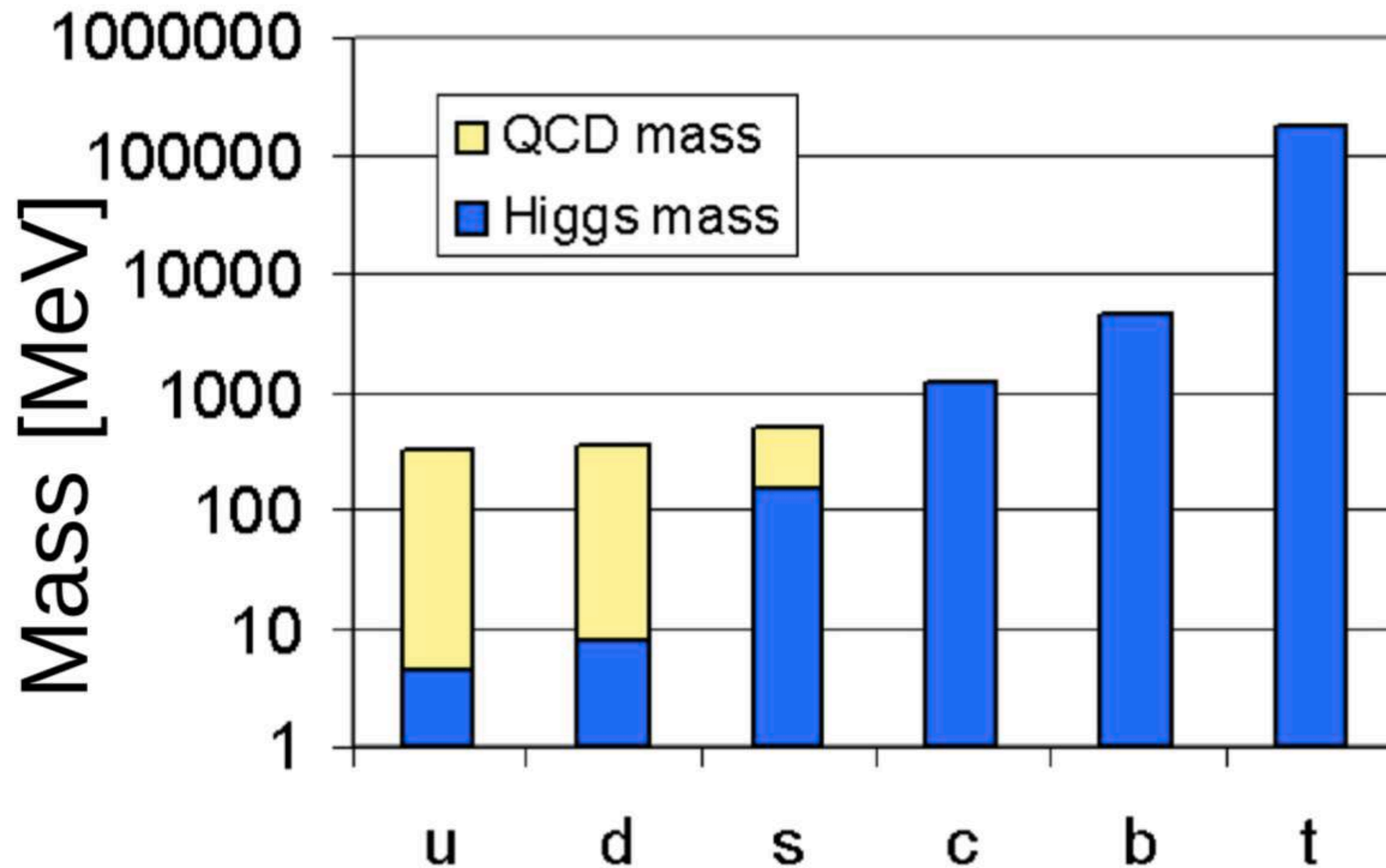


Confinement- no free quarks



Visible Matter

QCD is $\sim 95\%$ of visible (baryonic) matter

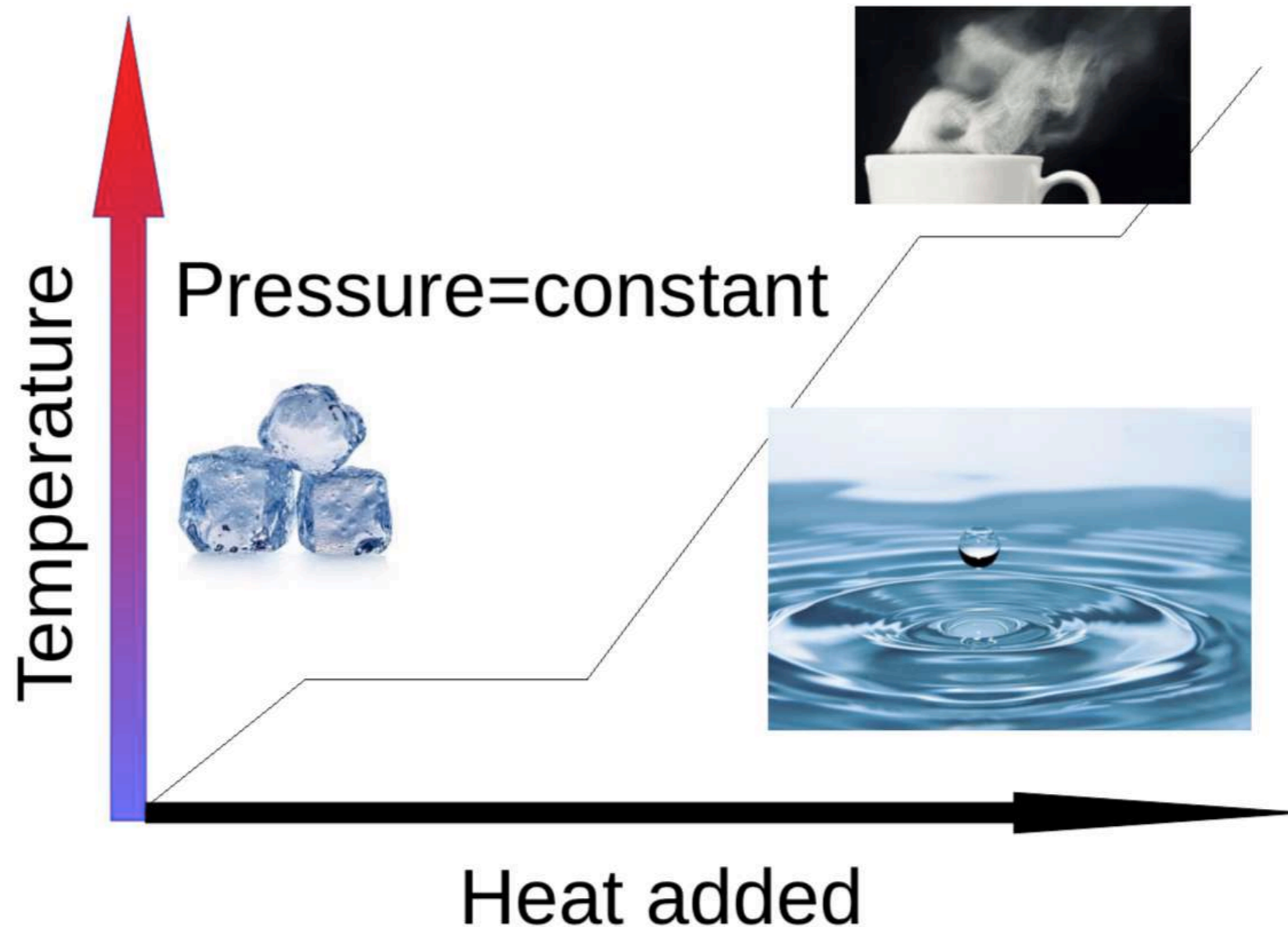


$$M_{\text{quark}} = 2 \cdot 2.3 + 4.8 \text{ [MeV]} \\ = 9.4 \text{ [MeV]}$$

$$M_{\text{proton}} = 938 \text{ [MeV]}$$

$$M_{\text{QCD}} = 928.6 \text{ [MeV]}$$

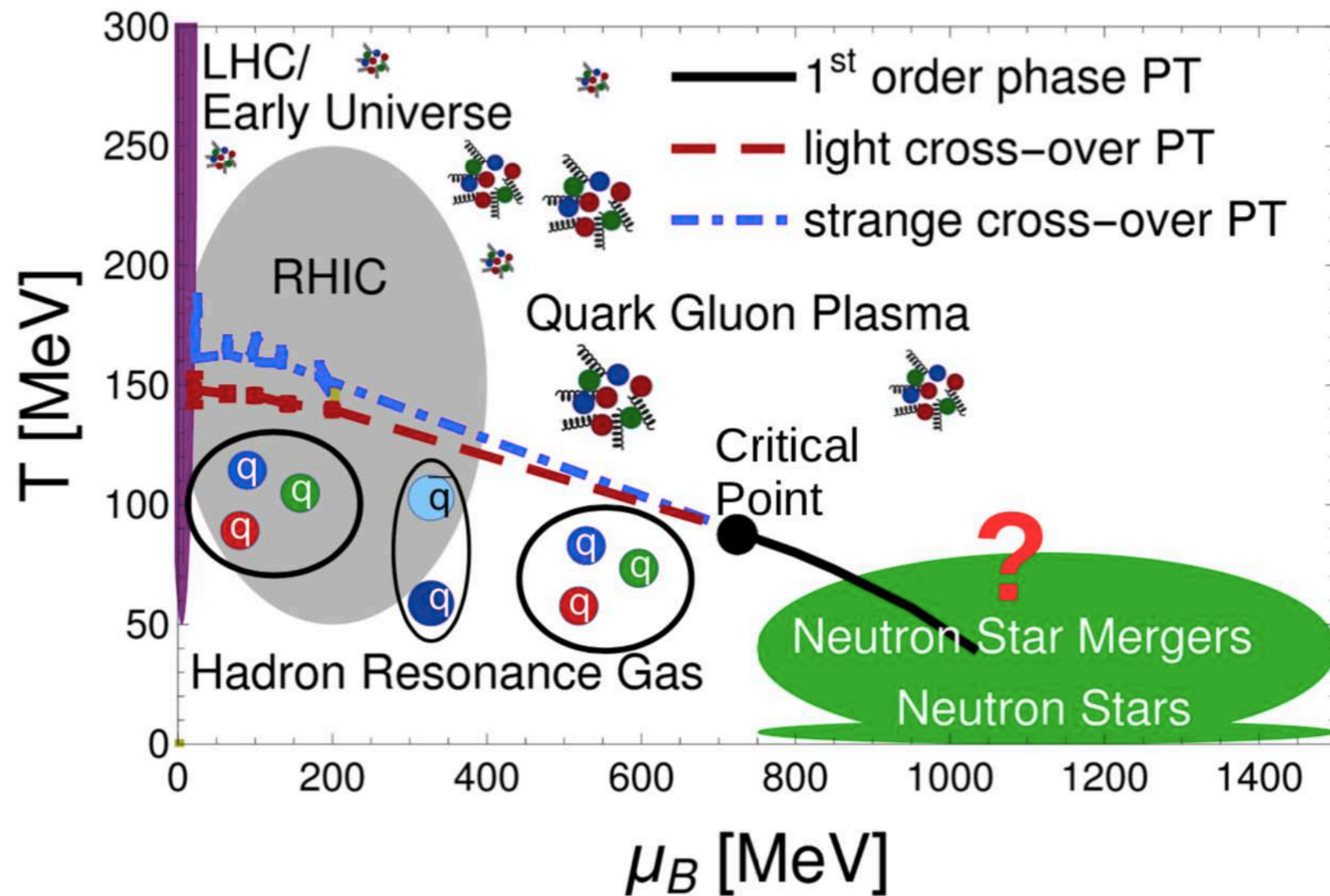
Phase Transitions of Water



Studying strong interactions/Quarks and Gluons

Can we \uparrow T enough to deconfine Quarks and Gluons?

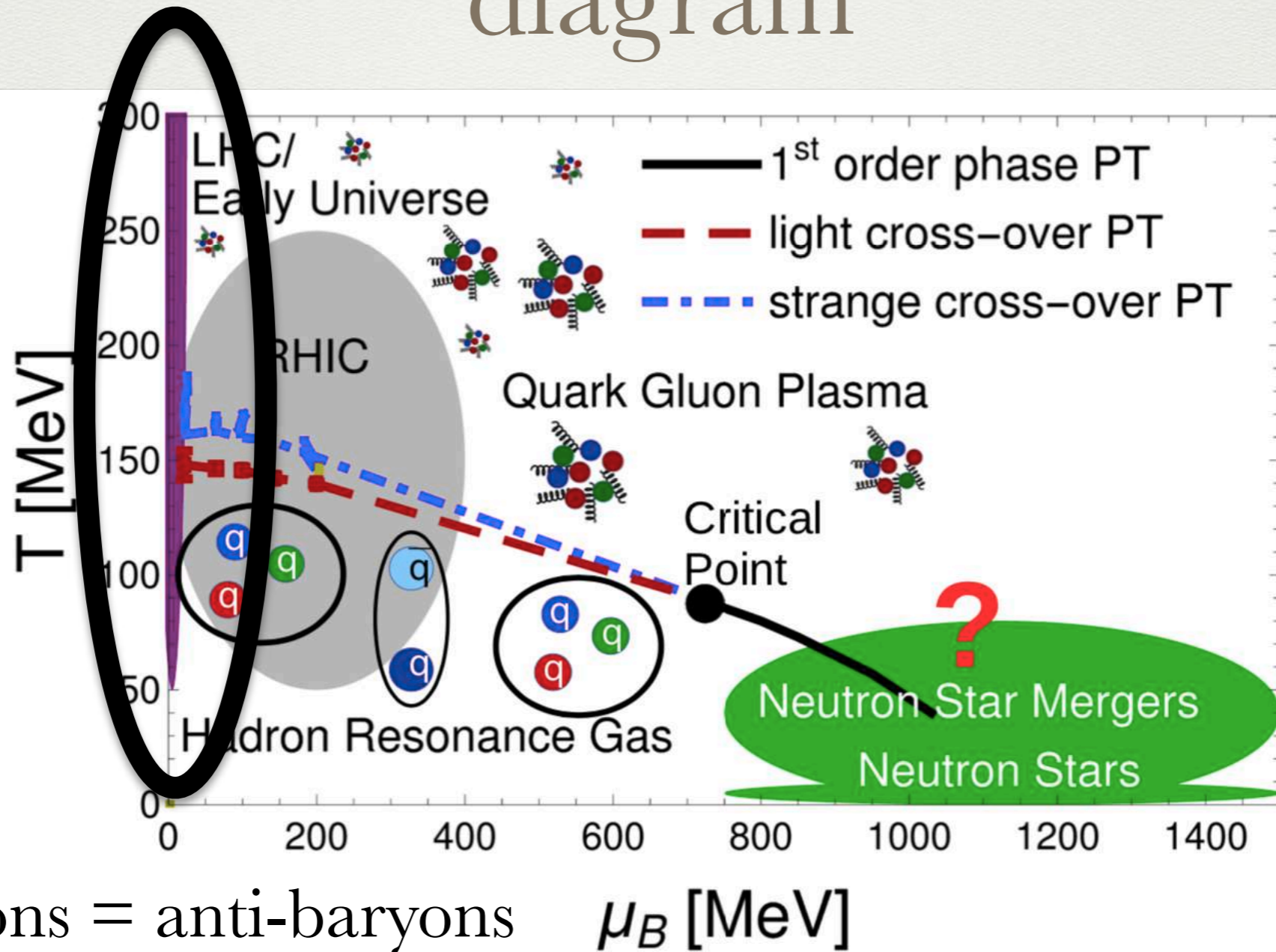
Current Cartoon of the QCD phase diagram



References

Light transition Phys.Lett. B738 (2014) 305-310; **Strange Transition** JNH and Ratti arxiv 1804.10661 ; **Neutron Star (mergers)** V. Dexheimer arXiv:1708.08342; **Holography** Critelli, JNH et al, Phys.Rev. D96 (2017) no.9, 096026

Current Cartoon of the QCD phase diagram



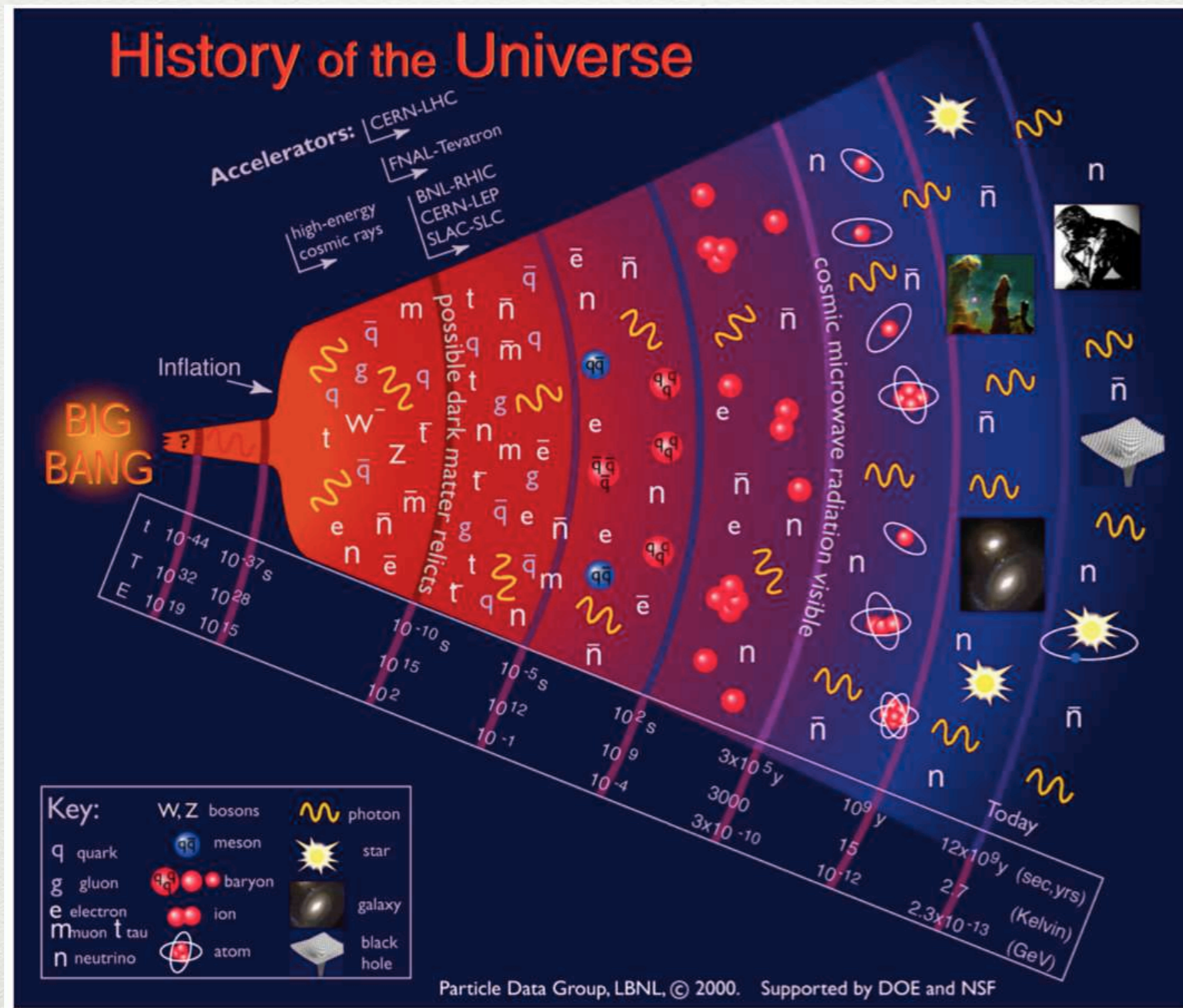
References

Light transition Phys.Lett. B738 (2014) 305-310; **Strange Transition** JNH and Ratti arxiv 1804.10661 ; **Neutron Star (mergers)** V. Dexheimer arXiv:1708.08342; **Holography** Critelli, JNH et al, Phys.Rev. D96 (2017) no.9, 096026

Deconfined Quarks and Gluons in the Early Universe

$\sim 10^{-6}$ s after the Big Bang \rightarrow Quark Gluon Plasma

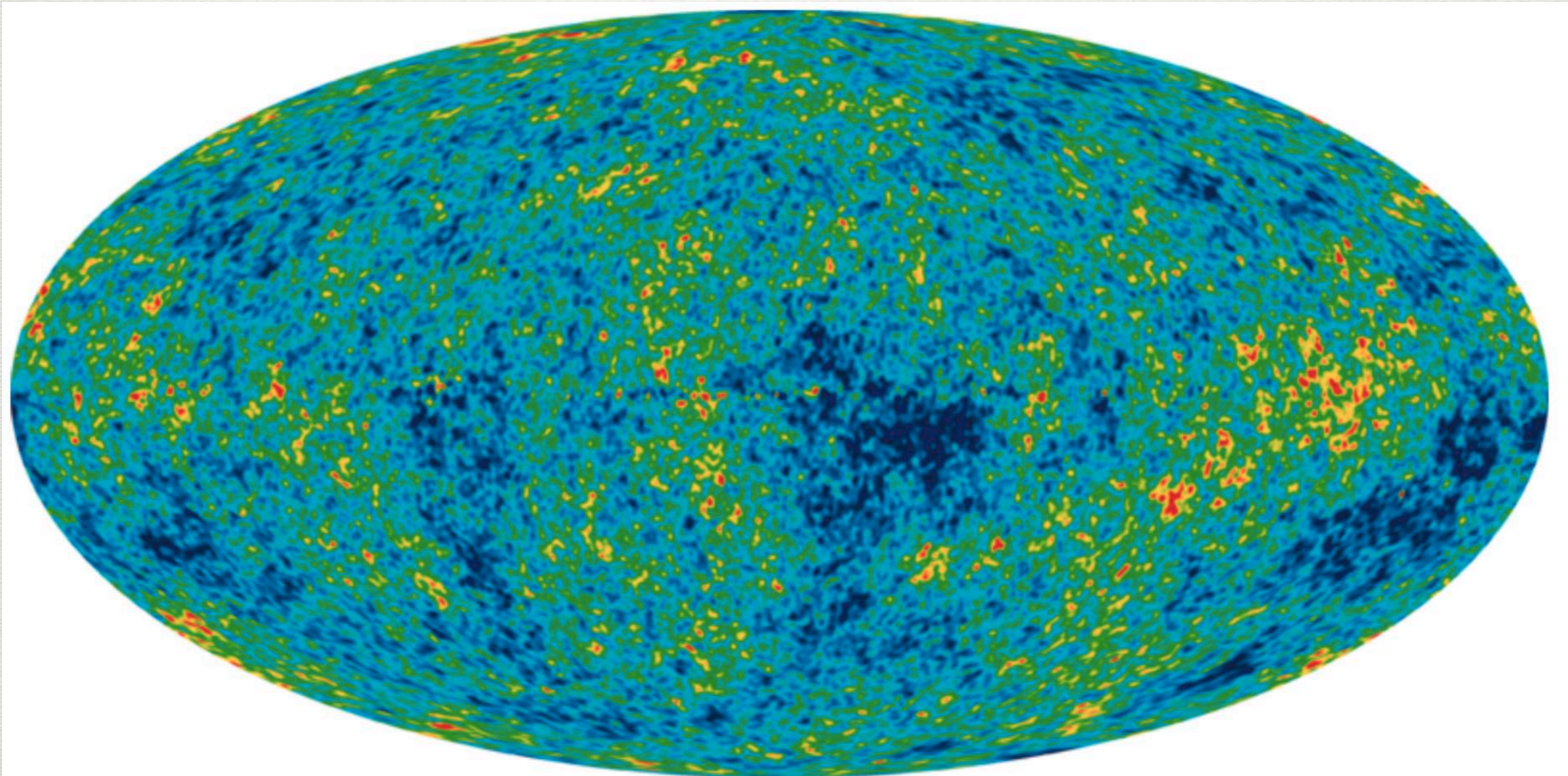
(1975 Collins and Perry)



How far back in time can we see?

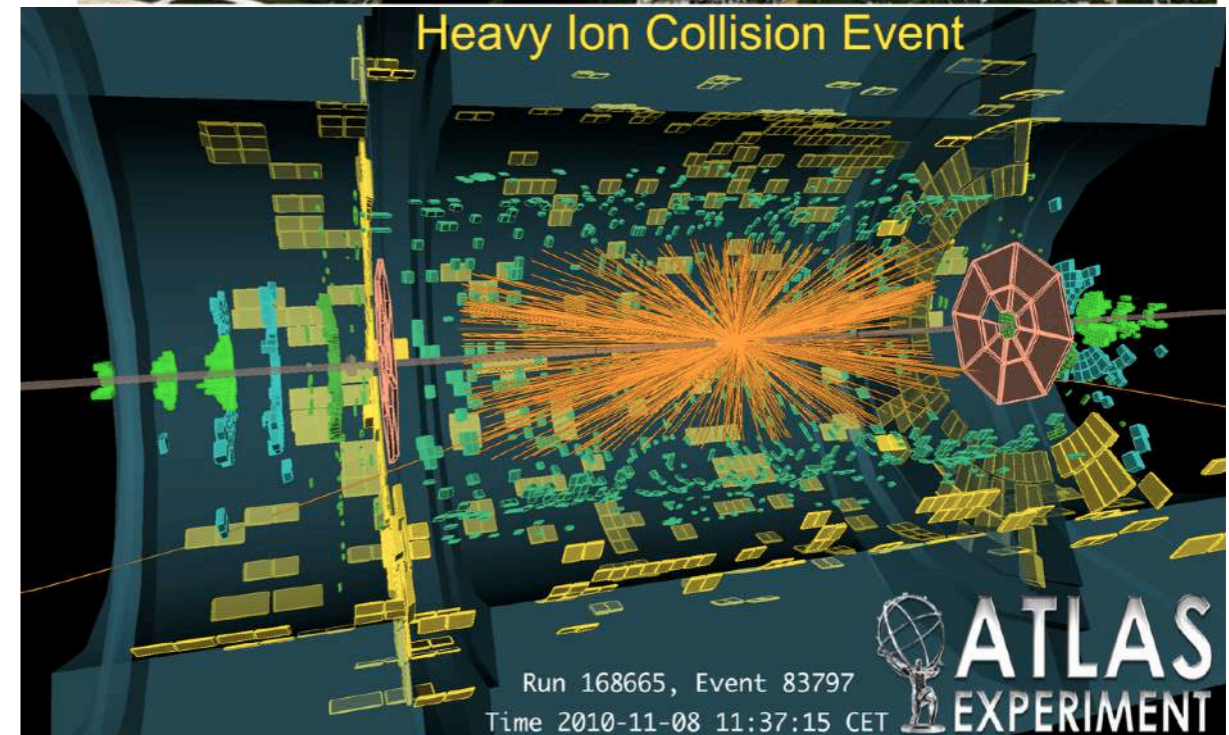
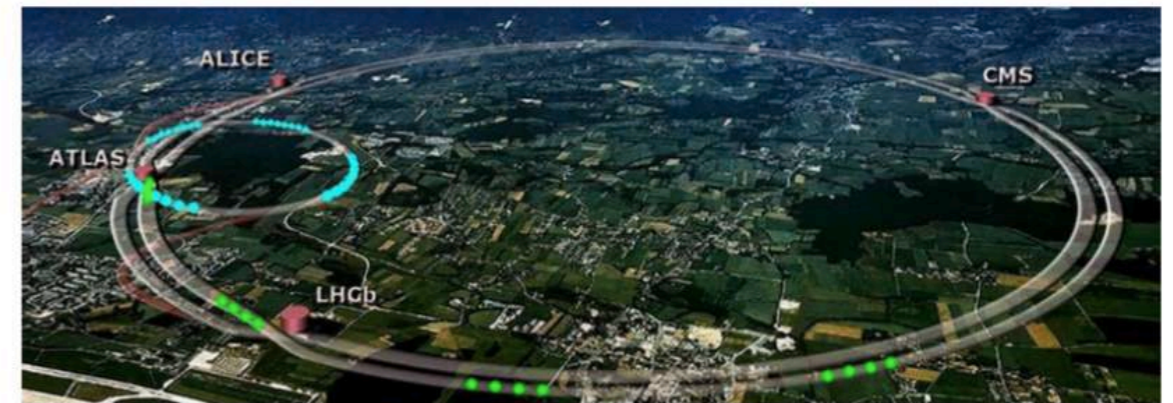
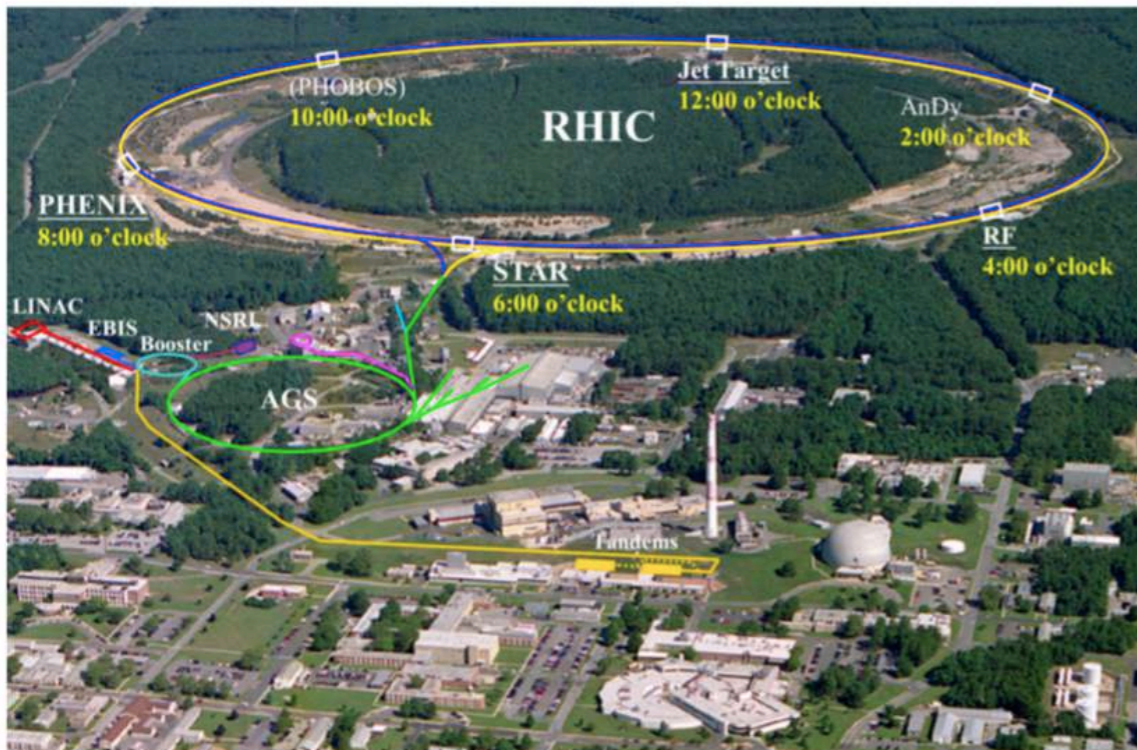
Cosmic Microwave Background $\sim 10^5$ **years** after Big Bang

Quark Gluon Plasma existed $\sim 10^{-6}$ **seconds**



Little Bangs in the Lab

The Large Hadron Collider and RHIC create "little bangs":
deconfined quarks and gluons in the lab



T=4 trillion degrees Celsius

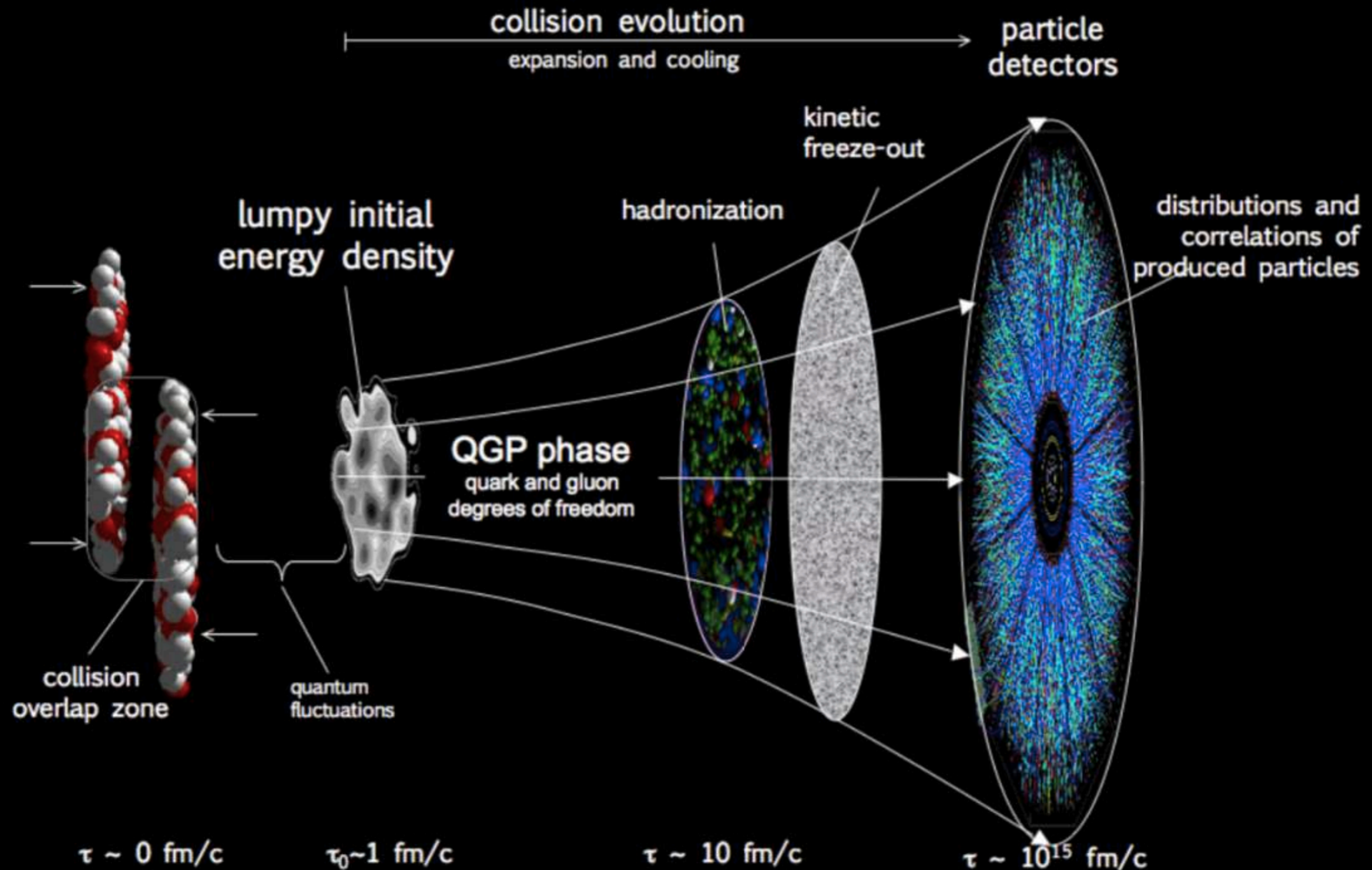


10^8 times hotter than the sun!

Evolution of a heavy-ion collision

Smashing two gold ions at the speed of light

Nuclear collisions and the QGP expansion

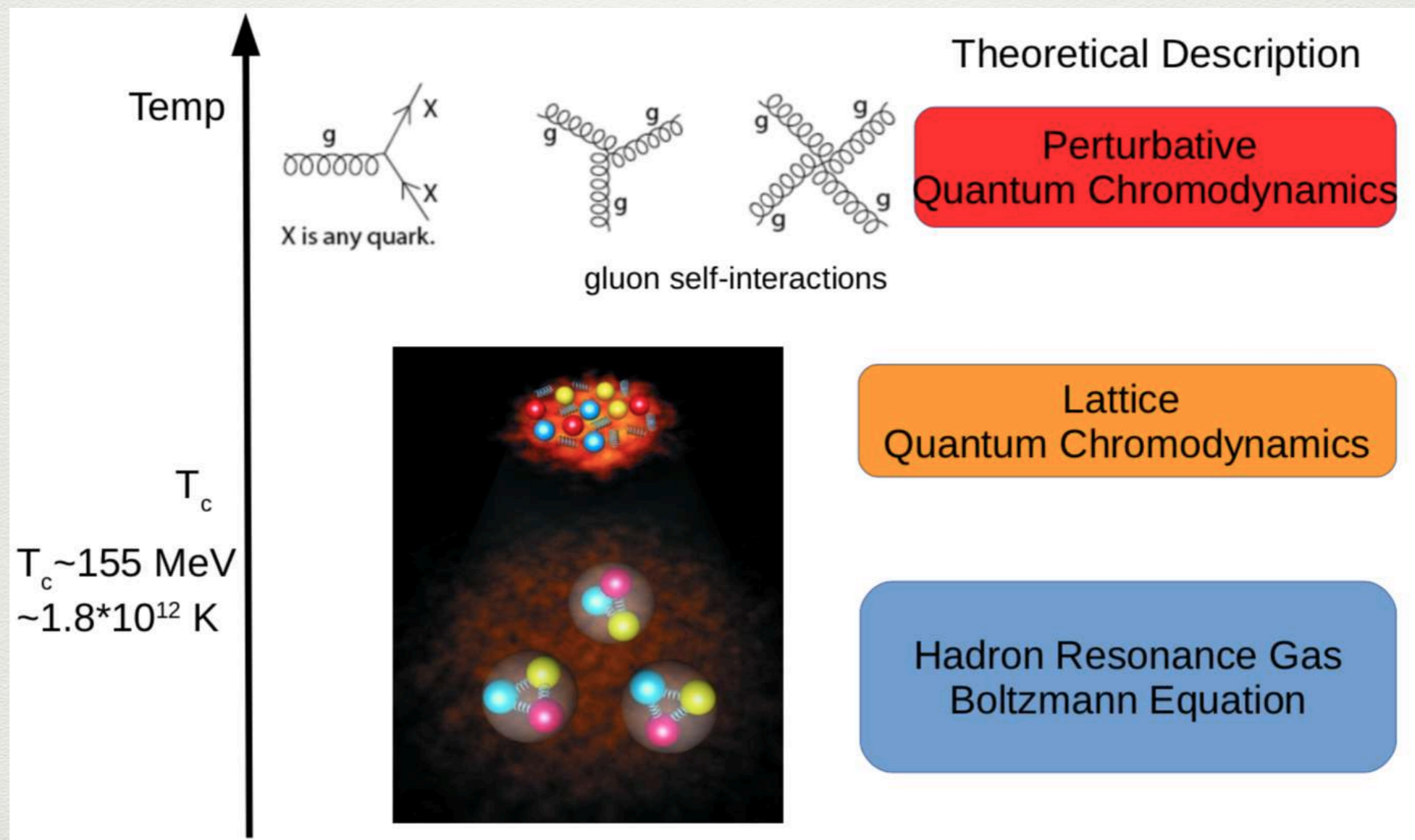


Big Bang vs. Heavy-Ion Collisions

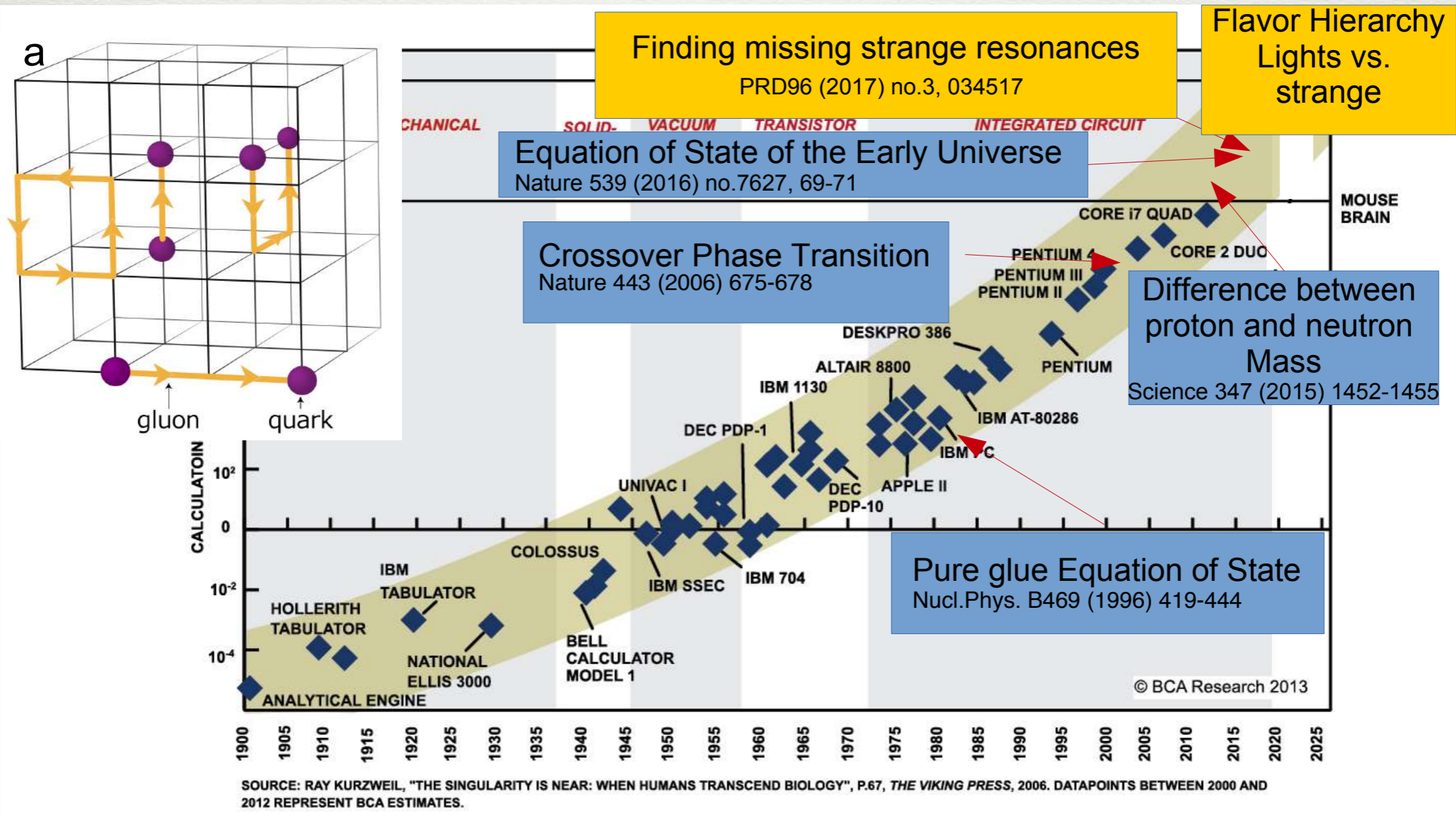
	Quark Epoch	Little Bang
Similarities	Pressure, entropy, energy density	
	Quark/gluons vs. Hadrons	
	Temperature when hadrons are formed	
	Strong Force	
	Nearly perfect fluid	
Differences	System Size	
	Entire Universe	10^{-14} m
		Finite volume effects
	Equilibrated	Out-of-Equilibrium
		Viscosity matters!
		Expansion rate large
	1 data point	Billions of events
		Initial Conditions (many different shapes)

Solving Quantum Chromodynamics

$$L_{QCD} = \underbrace{-\frac{1}{4} F_{\mu\nu}^a F_a^{\mu\nu}}_{\text{Gluon Interactions}} + \underbrace{\bar{\psi}^q \left(i\gamma_\mu D^\mu - m_q \right) \psi^q}_{\text{Quark Interactions}} \quad \text{where } D^\mu = \partial^\mu - ig \underbrace{A^\mu(x)}_{\text{Gluons}}$$



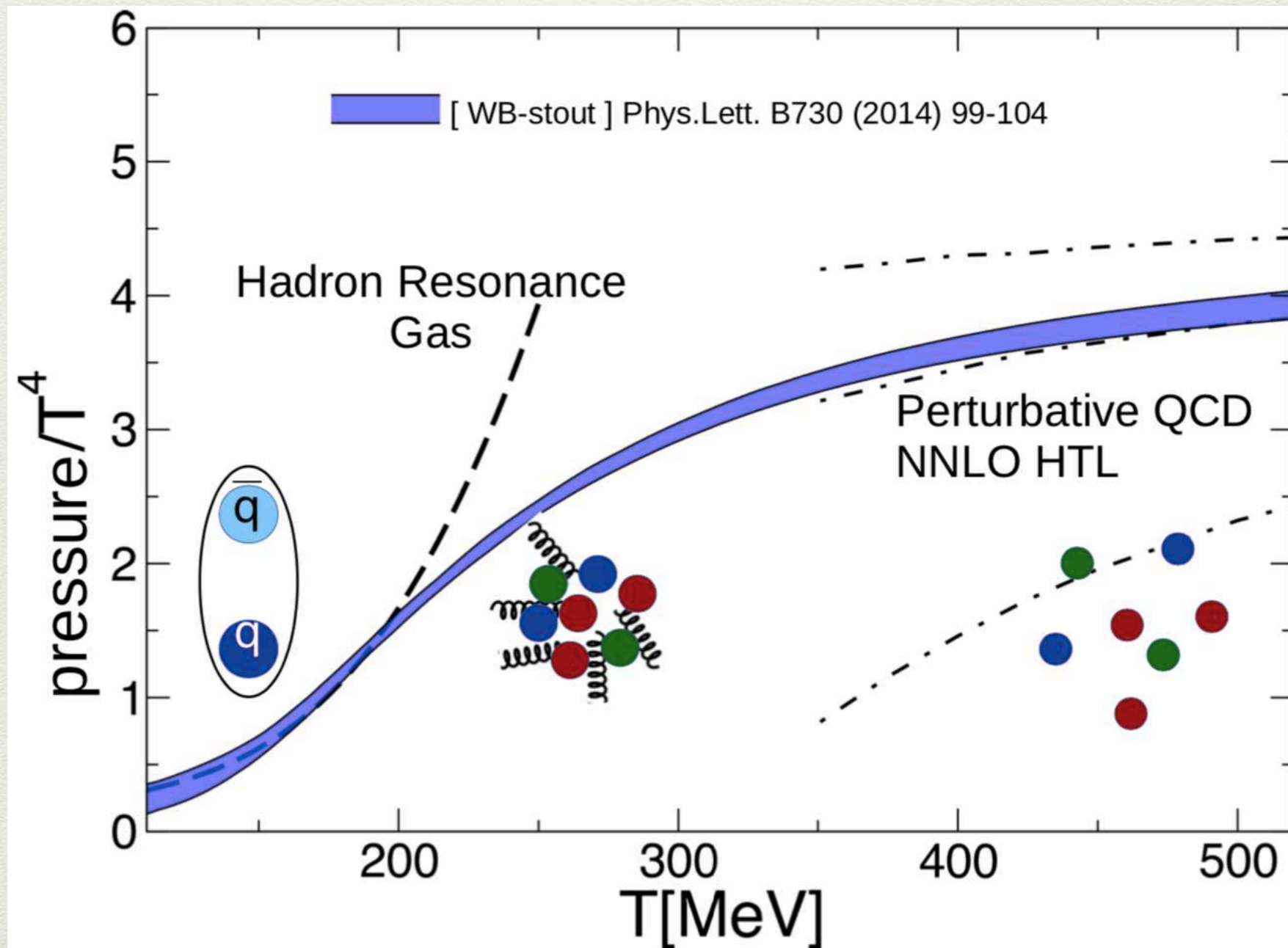
Lattice QCD: Solving Quantum Chromodynamics



Moore's Law: number of transistors per square inch on integrated circuits had doubled every year since their invention

Lattice QCD: Phase Transition

Cross-over phase transition $T \sim 155 \text{ MeV}$

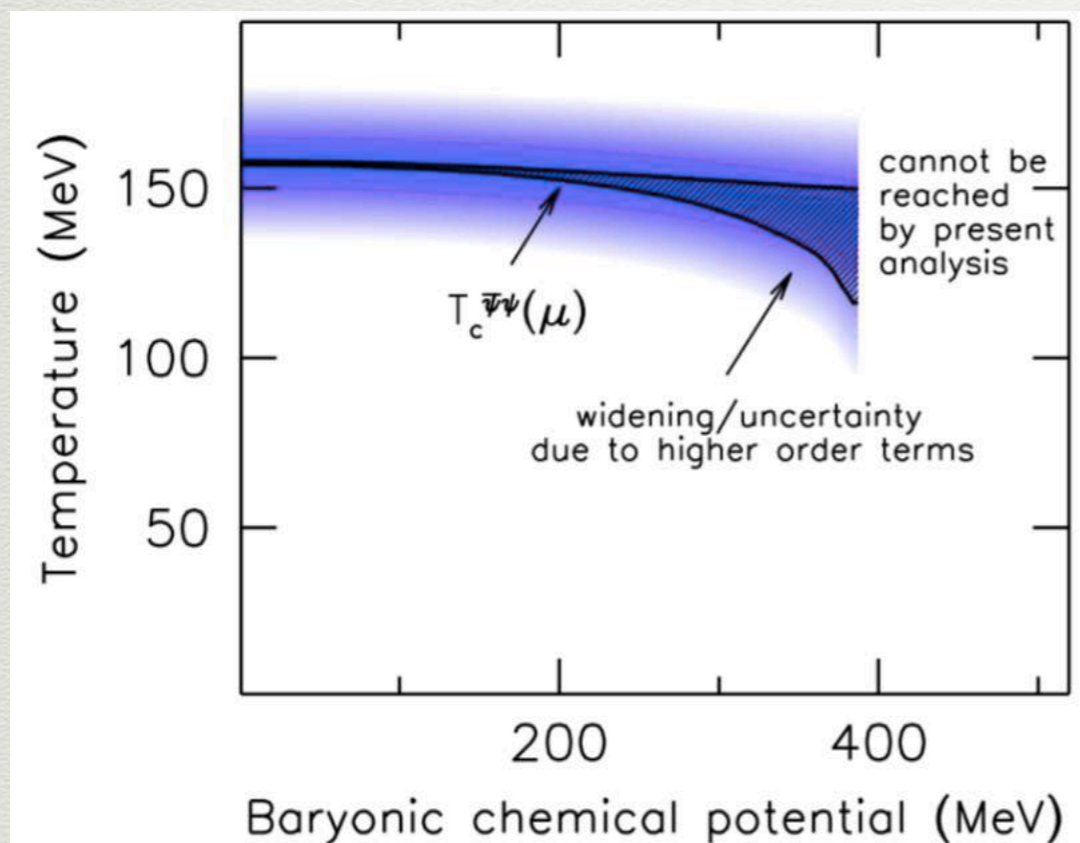


Limitations of Lattice QCD

Equilibrium Properties

Fermi Sign Problem

Baryons $>$ # anti-baryons



[WB] R. Bellwied et al., Phys.Lett. B751 (2015) 559-564

Work around - Taylor expansion

Out-of-Equilibrium

- Transport coefficients
- Dynamical description of the Quark Gluon Plasma
- Effective Models:
The Quark Gluon Plasma can be described by **relativistic viscous hydrodynamics** with a hadronic afterburner

What is a good (or “perfect”) fluid?

Good fluid



Bad fluid



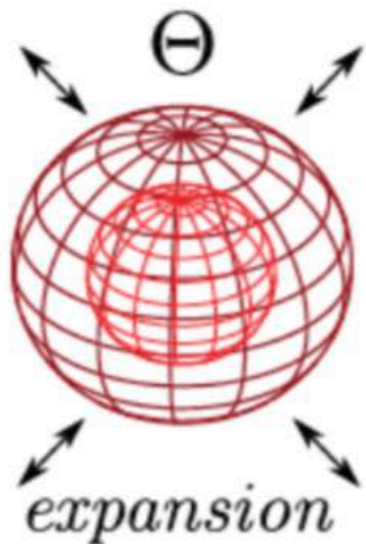
Best fluids are the closest to an ideal fluid i.e vanishing viscosity

Transport coefficients/viscosities

Transport coefficient: Perturb the fluid from equilibrium- how quickly does it return to equilibrium?

Bulk

$$\Pi \sim -\zeta \Theta$$



Shear

$$\pi^{\mu\nu} \sim 2\eta\sigma^{\mu\nu}$$

$$\sigma_{\mu\nu}$$

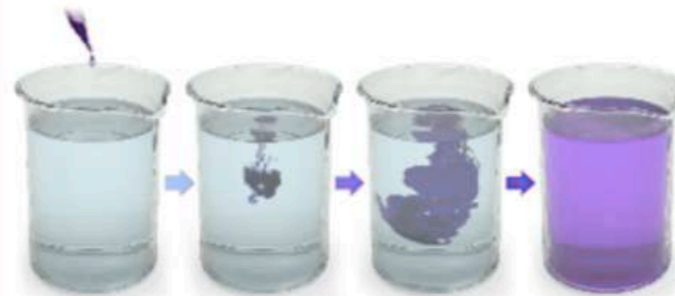


shear

Diffusion

$$q_{\perp}^{\mu} \sim \kappa \nabla_{\perp}^{\mu} (\mu/T)$$

QCD conserved charges
(B,S,Q)



Diffusion

Vorticity

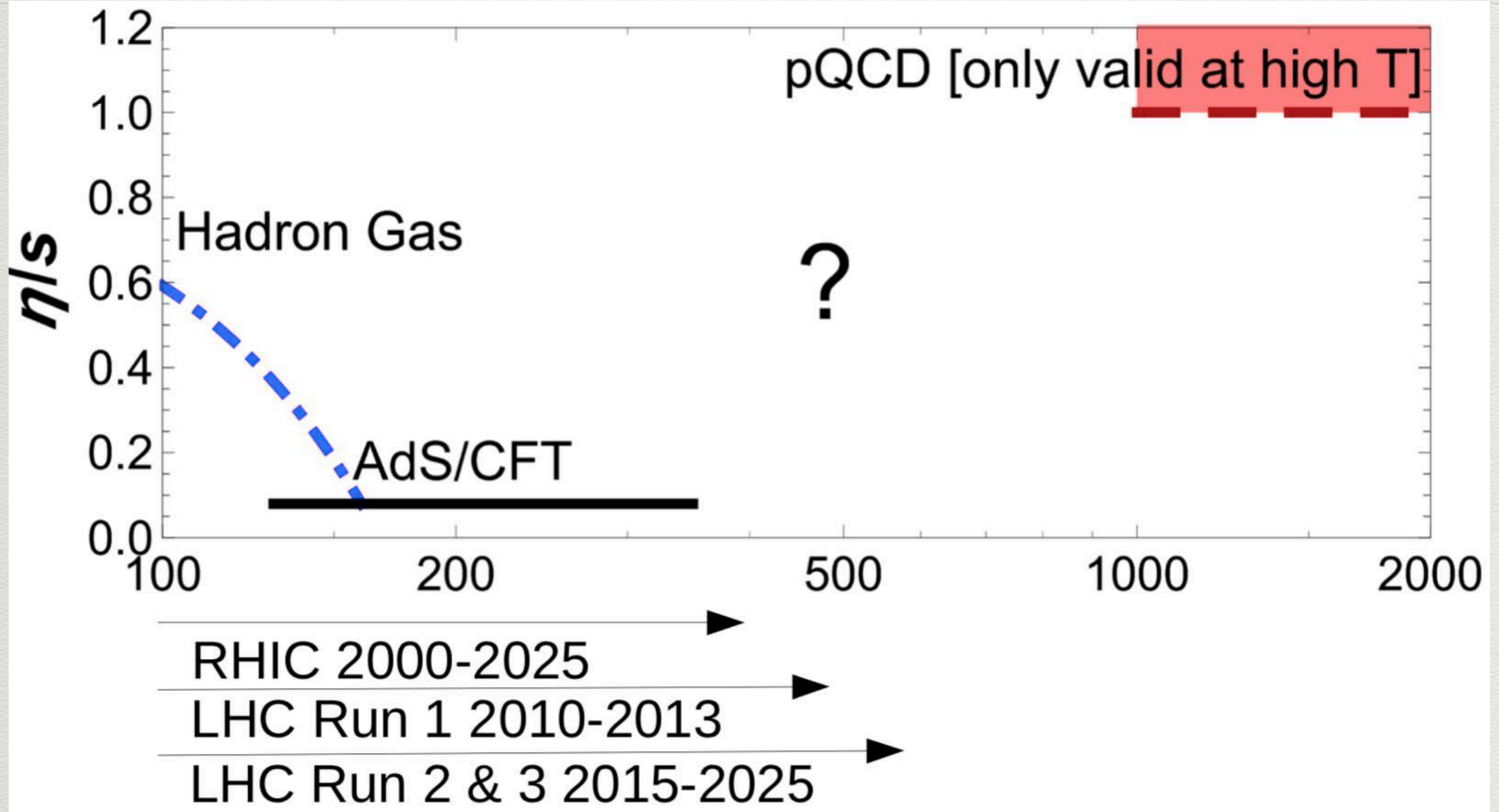
$$\omega_{\mu}$$



vorticity

Viscosity - resistance to deformation or “thickness” of liquid

Experimental probes of $\eta/s(T)$

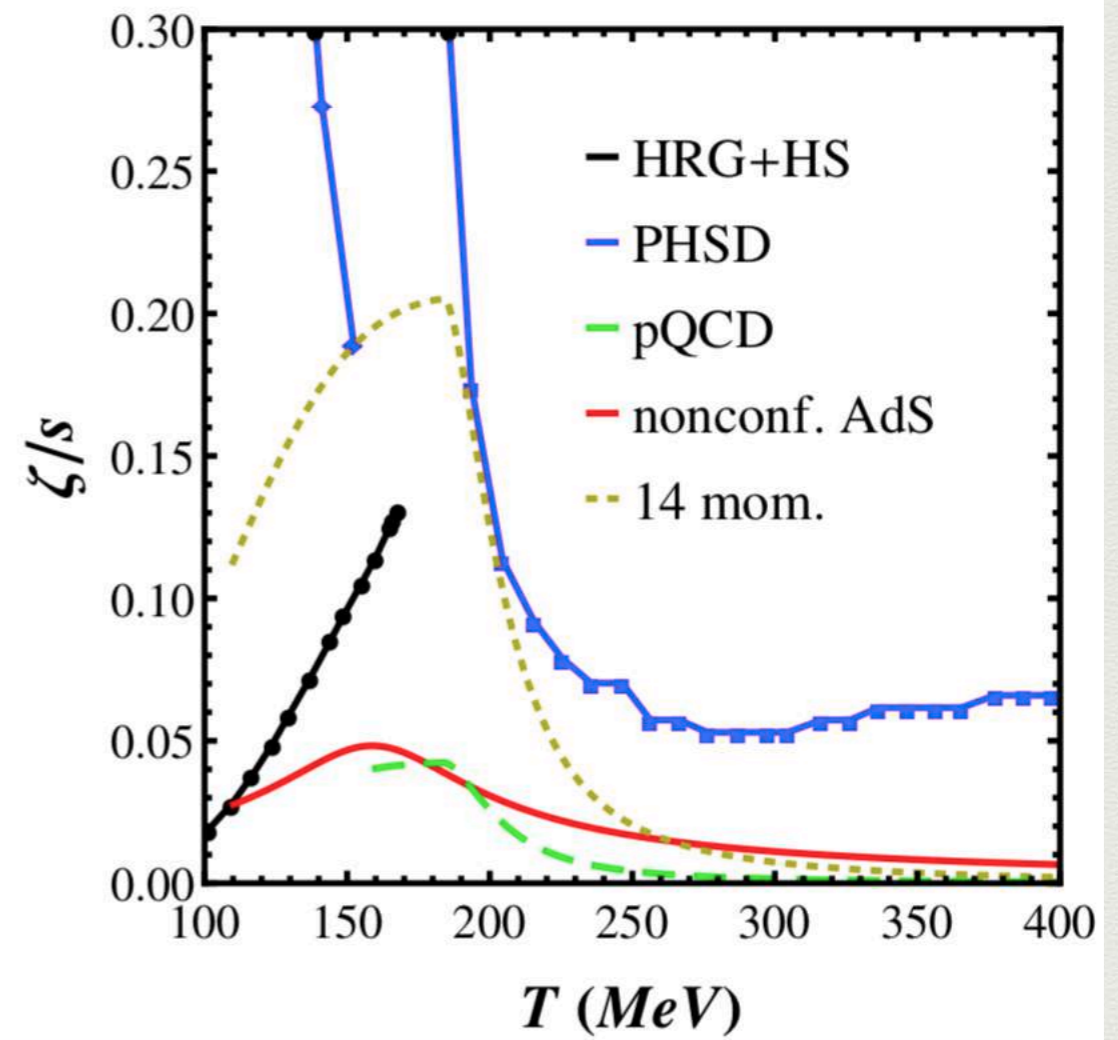
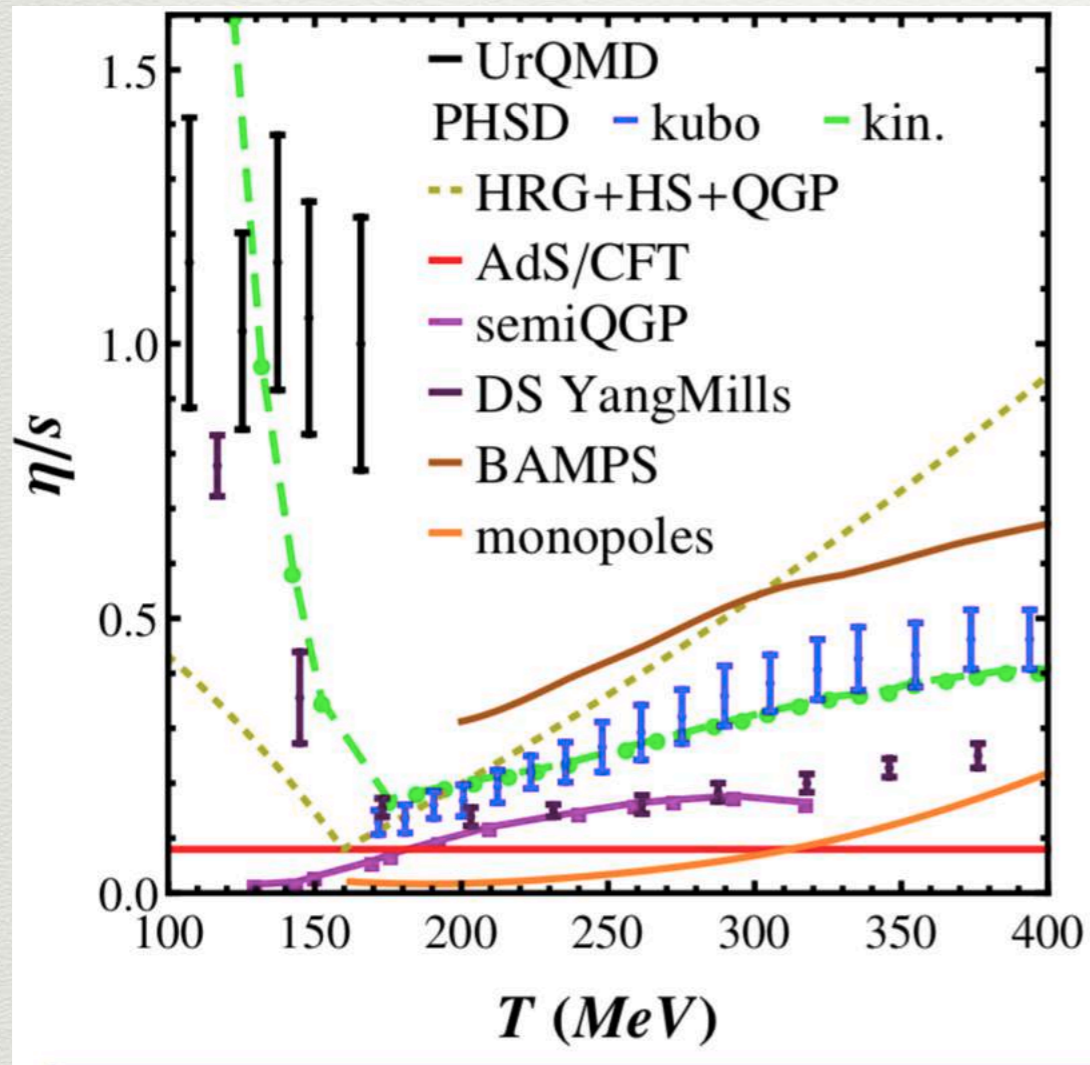


Hadron Gas: JNH et al, PRL103(2009)172302; PRC86(2012)024913

AdS/CFT: Kovtun, Son, Starinets PRL94(2005)111601

pQCD: Arnold, Moore, Yaffe JHEP 0011(2000)001 ; JHEP0305(2003)051

Theoretical calculations of viscosity

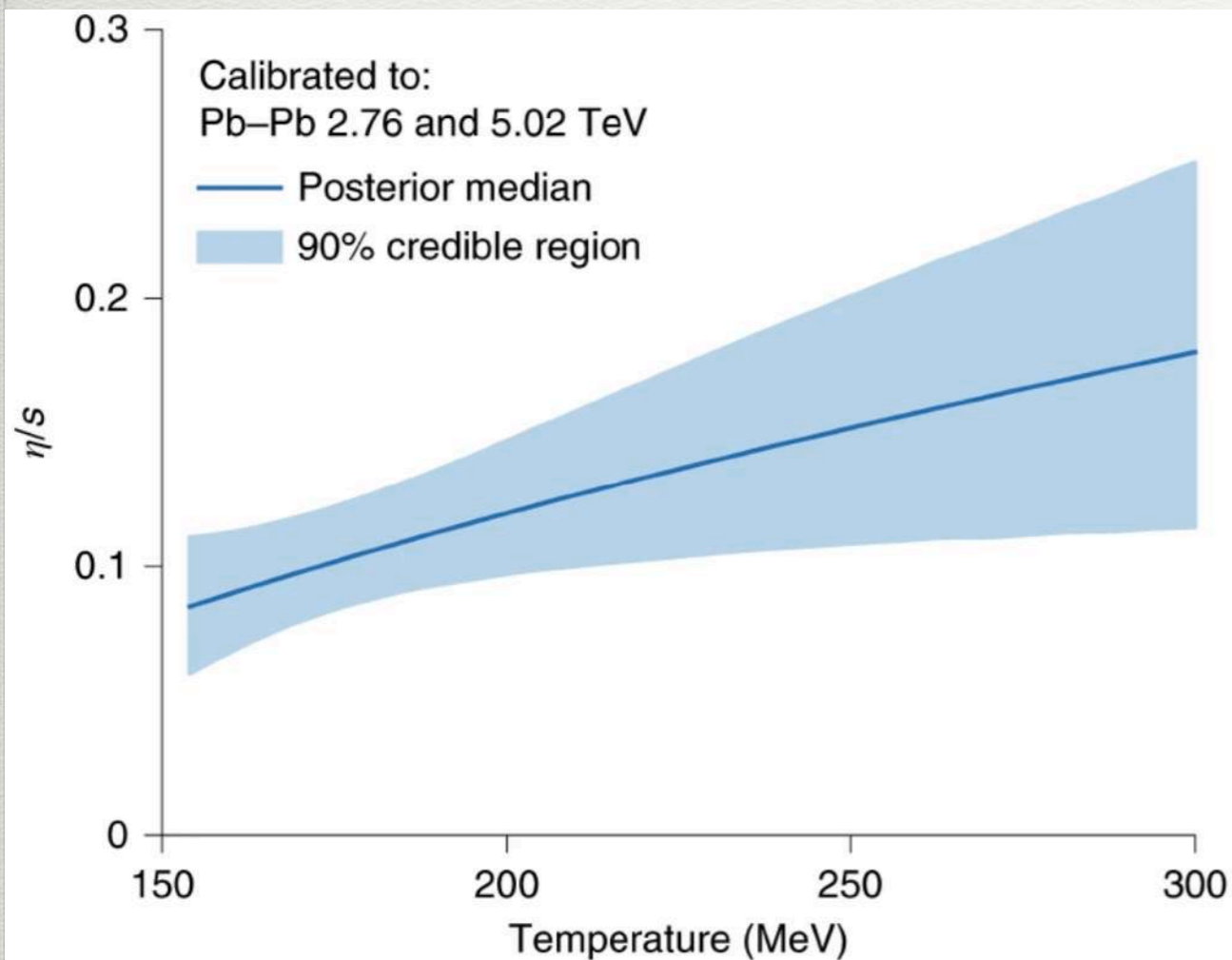


See references in *JNH arXiv:1512.06315*

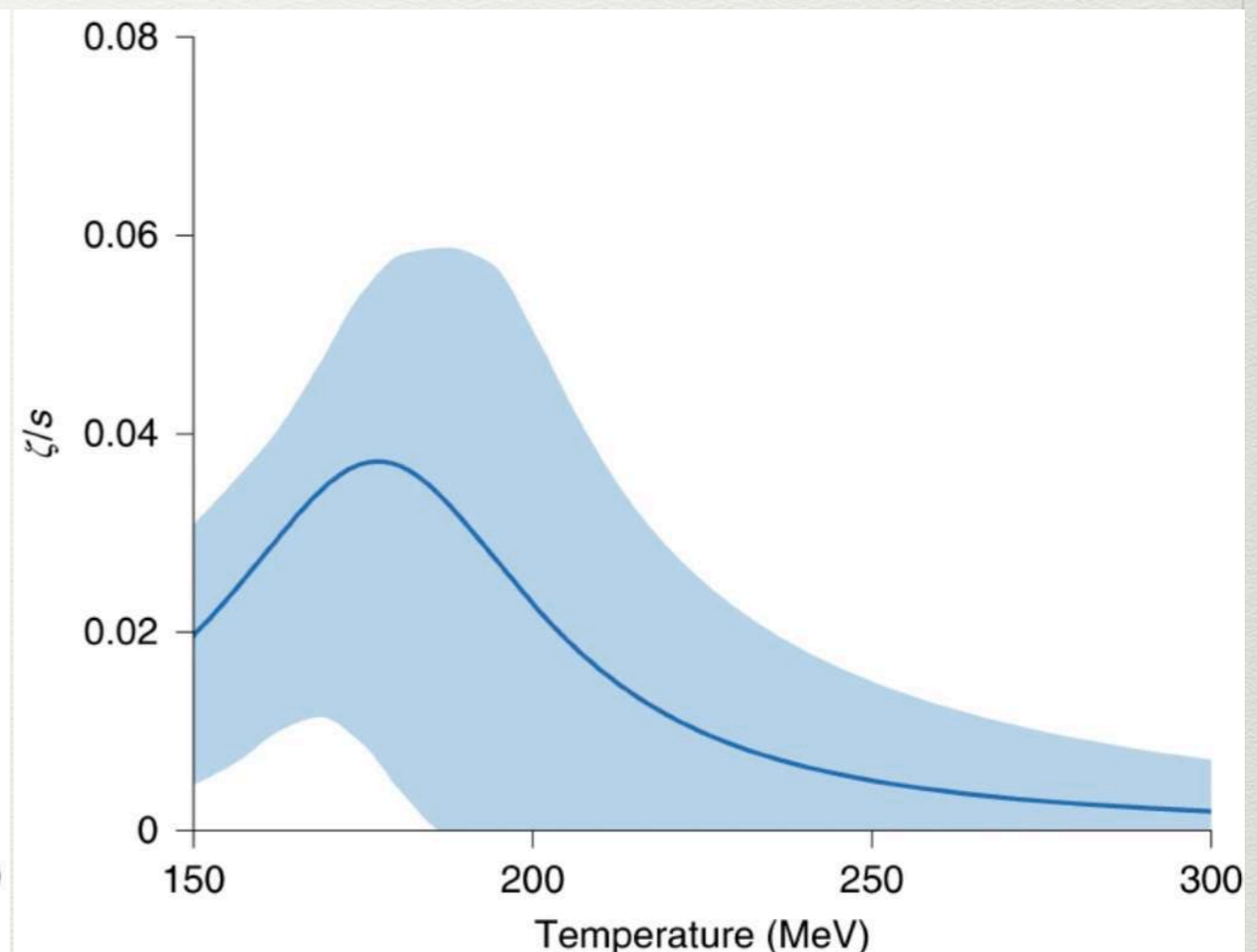
Dip expected: *Phys.Rev.Lett.* 97 (2006) 152303, *Nucl.Phys.* A769 (2006) 71-94, *Phys.Rev.Lett.* 103 (2009) 172302

Bayesian analysis (agnostic η/s & ζ/s)

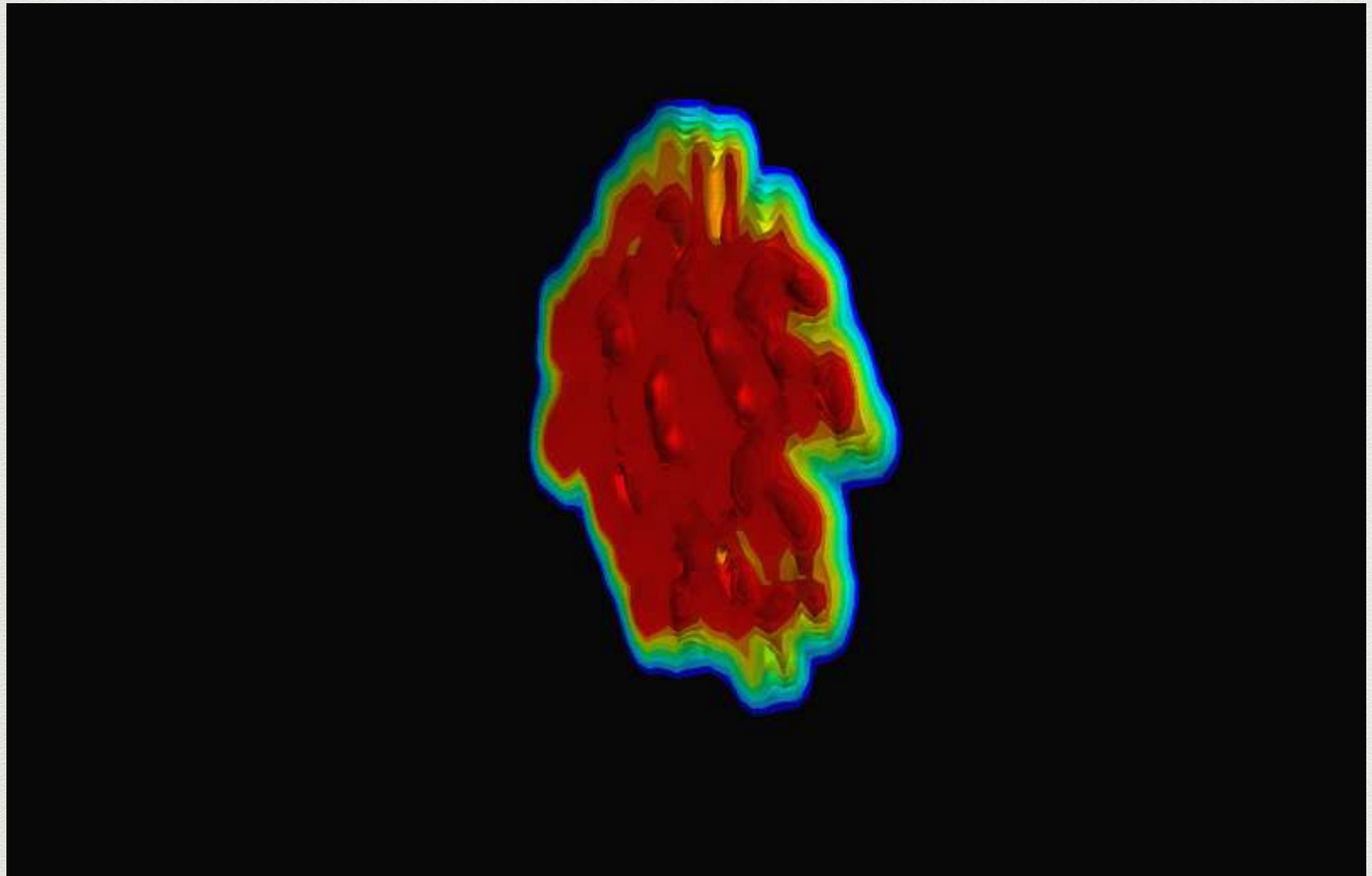
Shear viscosity



Bulk viscosity

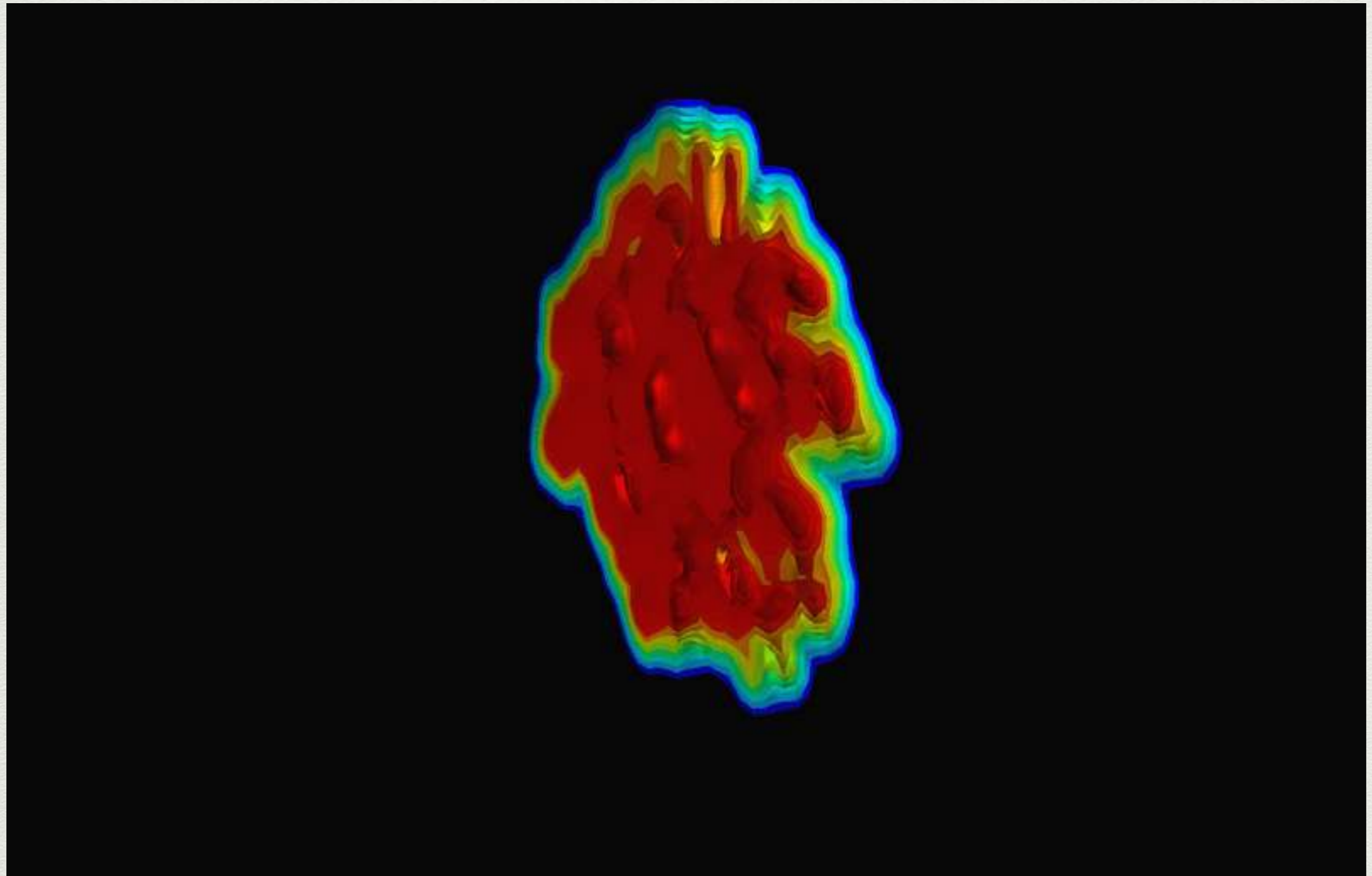


Relativistic fluids



Schenke IP-Glasma+MUSIC

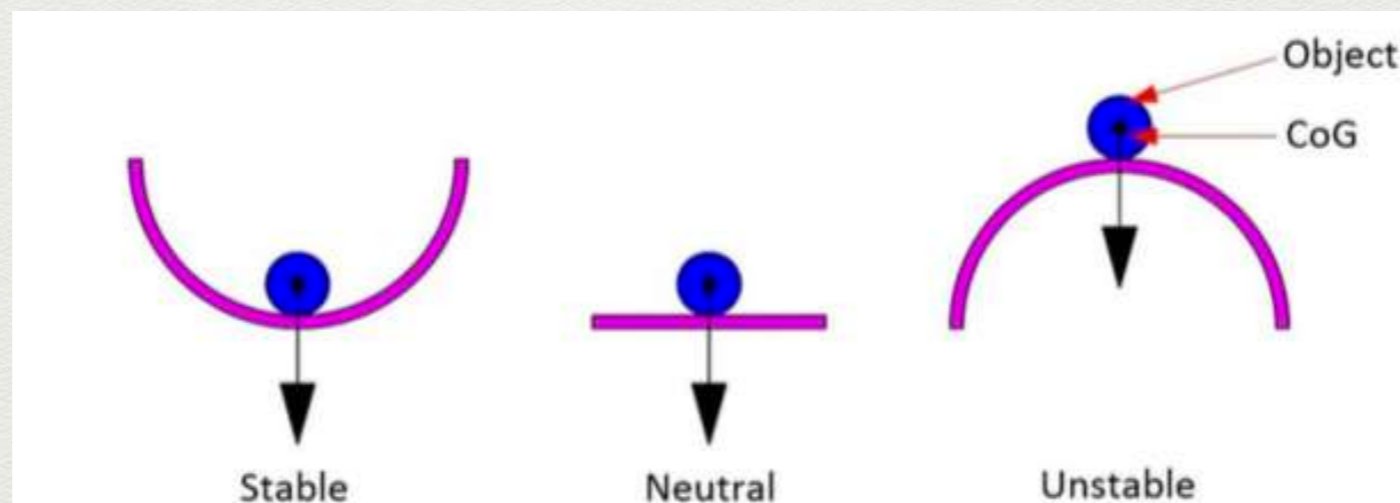
Relativistic fluids



Schenke IP-Glasma+MUSIC

Relativistic viscous fluid dynamics

- Navier stokes equations are used for non-relativistic systems with viscosity
- At relativistic velocities, Navier Stokes equations become **acausal** and **unstable**.



- Israel-Stewart equations incorporate a relaxation time (a finite time for the system to return to equilibrium)

Israel-Stewart Equations of Motion

Annals Phys. 118 (1979) 341-372

Conservation of Energy and Momentum

$$\partial_{\mu} T^{\mu\nu} = 0 \text{ and } \partial_{\mu} N^{\mu} = 0$$

The energy-moment tensor contains a bulk dissipative term Π and the shear stress tensor $\pi^{\mu\nu}$ is

$$T^{\mu\nu} = \varepsilon u^{\nu} u^{\nu} - (p + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$
$$\tau_{\pi} \left(\Delta_{\mu\nu\alpha\beta} D \pi^{\alpha\beta} + \frac{4}{3} \pi_{\mu\nu} \theta \right) = 2\eta \sigma_{\mu\nu} - \pi_{\mu\nu}$$
$$+ \Pi, N^{\mu} \dots$$

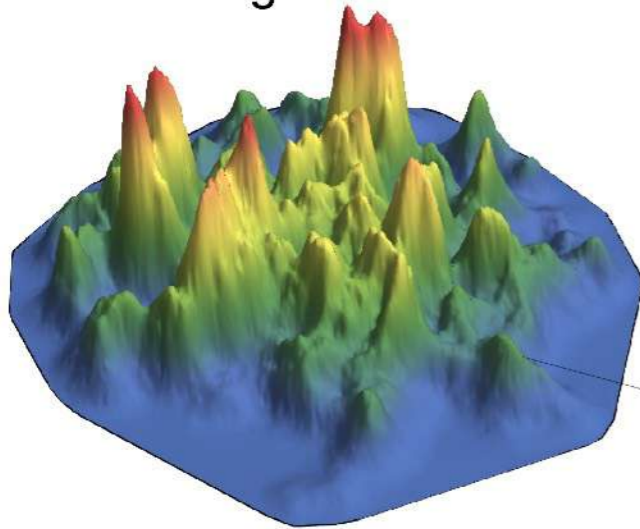
Coordinate System: $x^{\mu} = (\tau, x, y, \eta)$ where $\tau = \sqrt{t^2 - z^2}$ and

$$\eta = 0.5 \ln \left(\frac{t + z}{t - z} \right)$$

“Standard Model” of the Quark Gluon Plasma

Initial Conditions

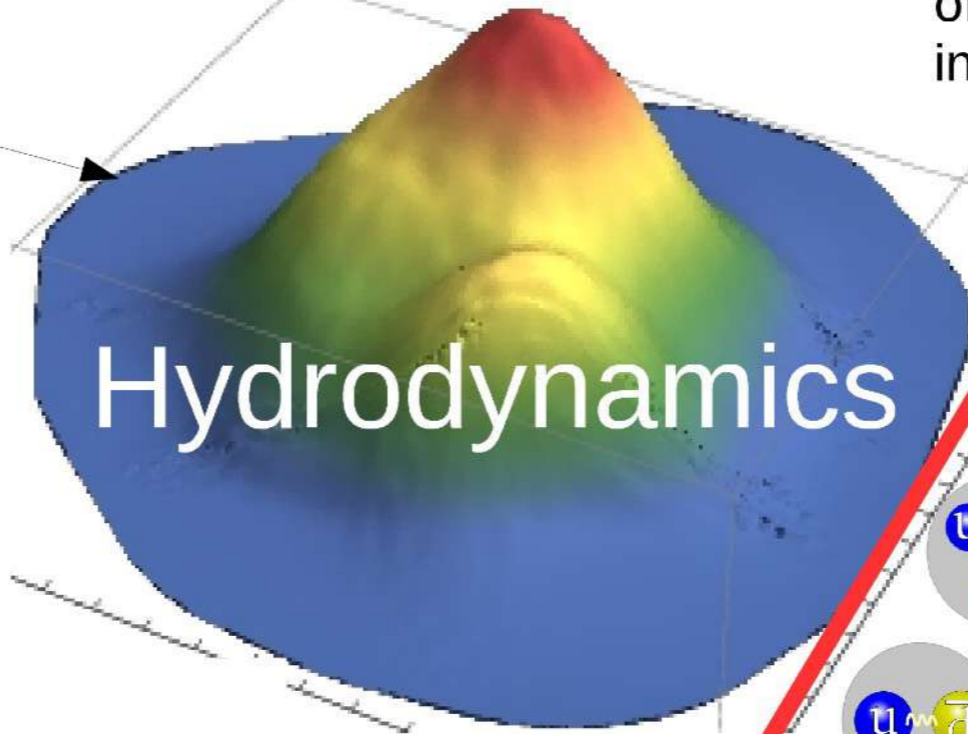
Quantum fluctuations in the position of protons, neutrons, quarks, and gluons



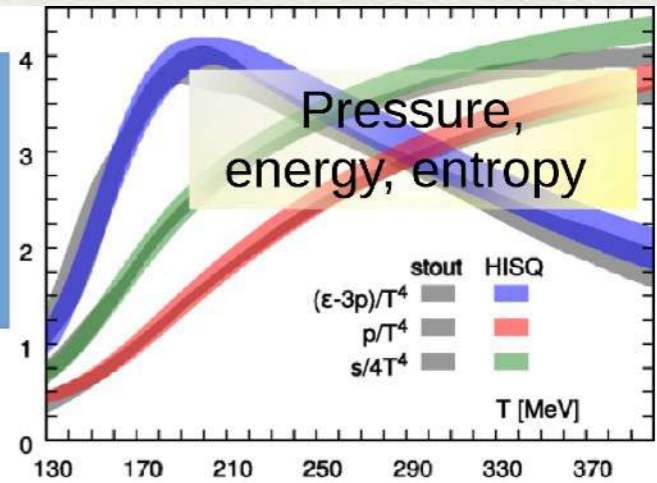
τ_0 initial time to switch on hydro

Hydrodynamics
(for heavy-ions collisions)
in a nutshell

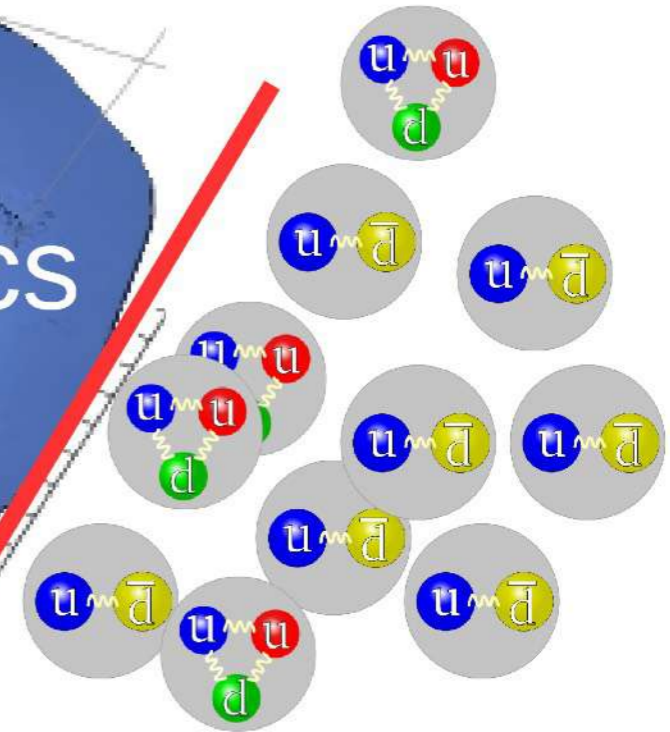
Hydrodynamics viscosity and thermodynamics



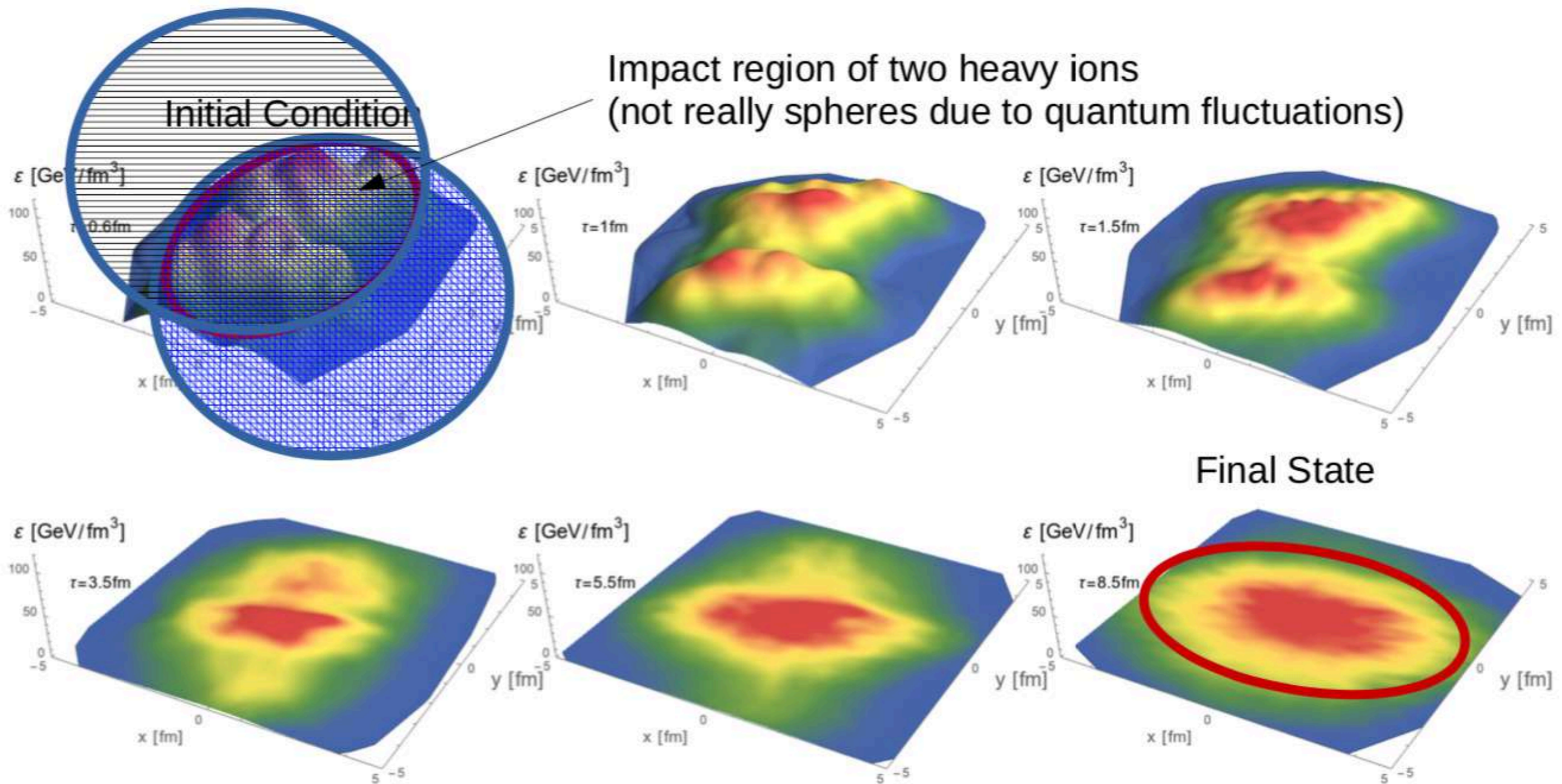
T_{sw} temperature at which the Quark Gluon Plasma switches to hadrons



Hadron Gas: number of hadrons, decays, interactions etc

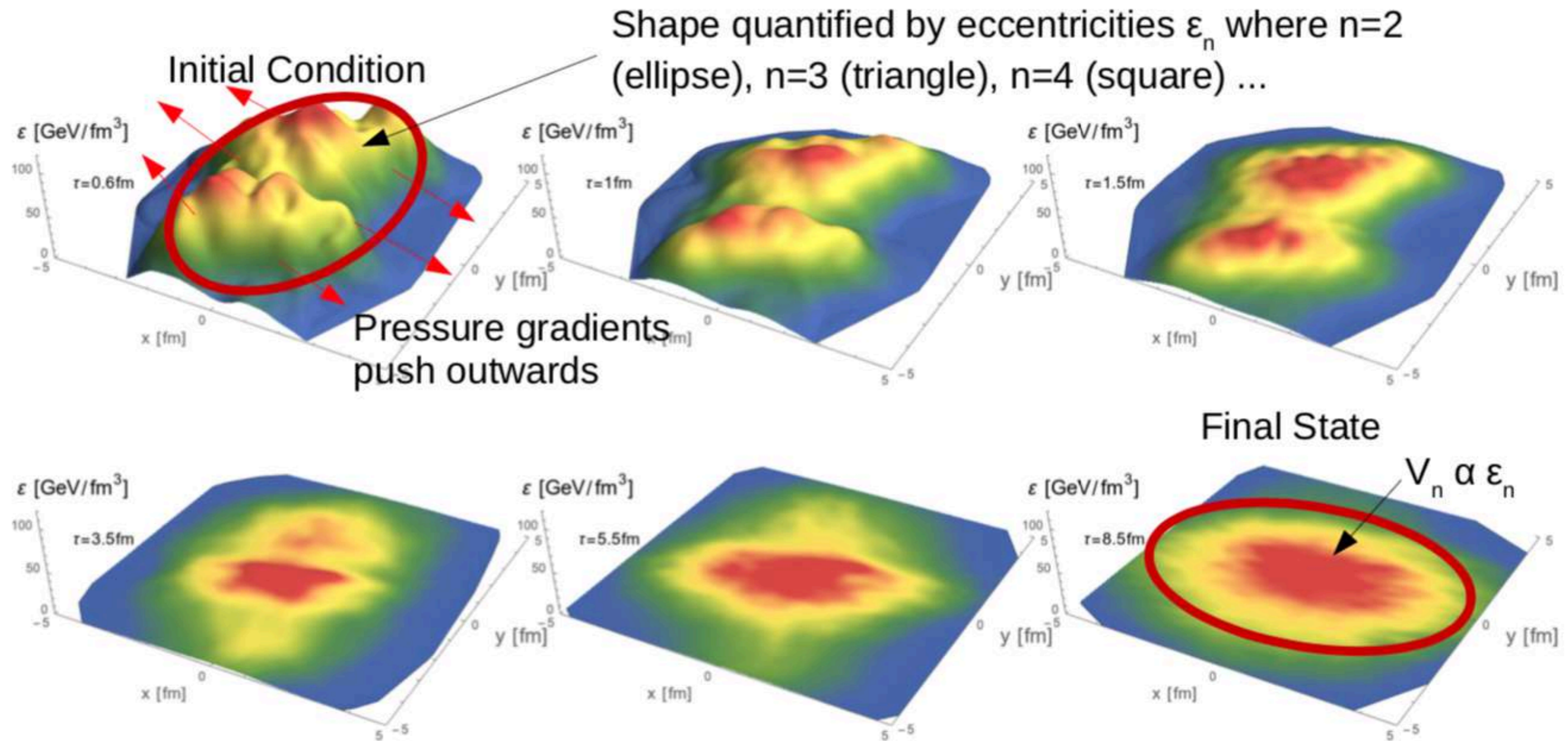


Initial conditions



Eccentricities ϵ_2 's are directly related to the final measured flow observables v_n 's

Initial conditions

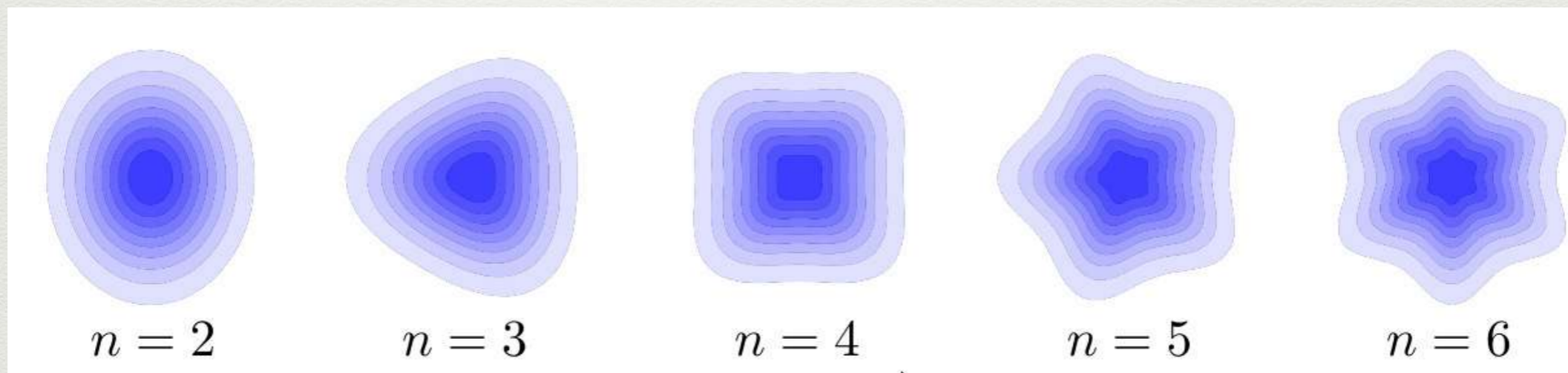


Eccentricities ε_2 's are directly related to the final measured flow observables v_n 's

Quantifying flow

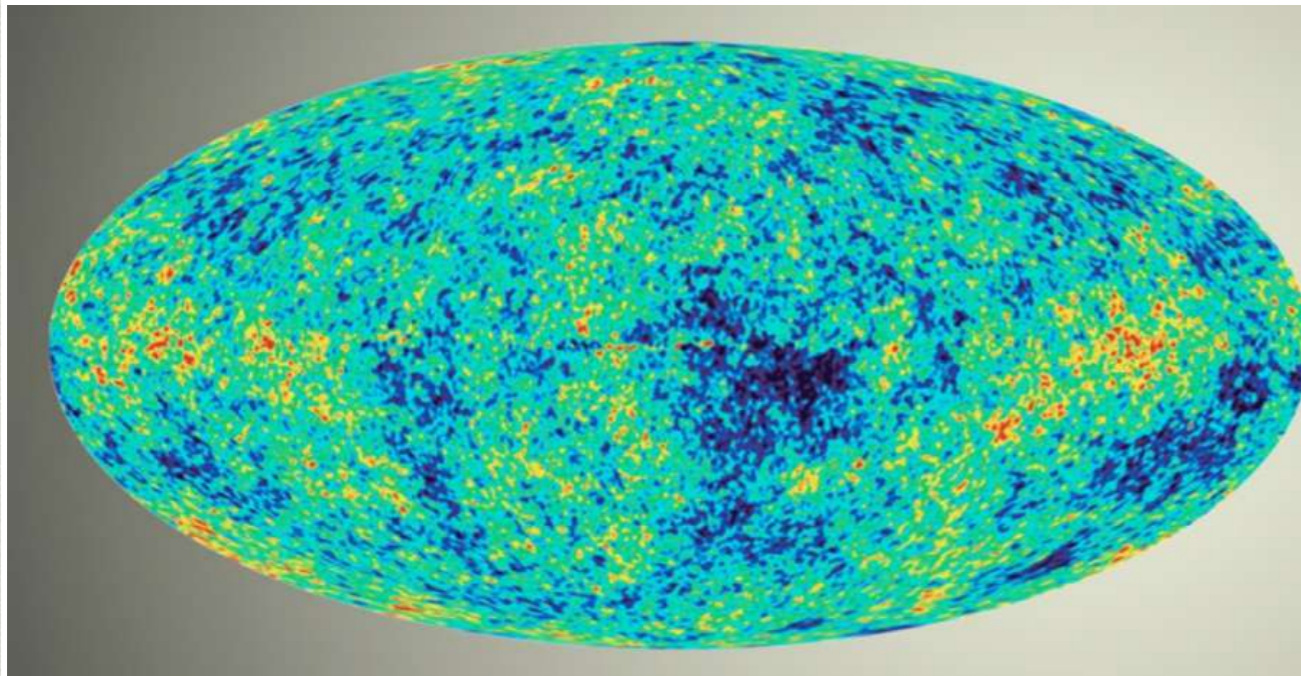
The distribution of particles can be written as a Fourier series

$$E \frac{d^3 N}{d^3 p} = \frac{1}{2\pi} \frac{d^2 N}{p_T dp_T dy} \left[1 + \sum_n 2v_n \cos \left[n (\phi - \psi_n) \right] \right]$$

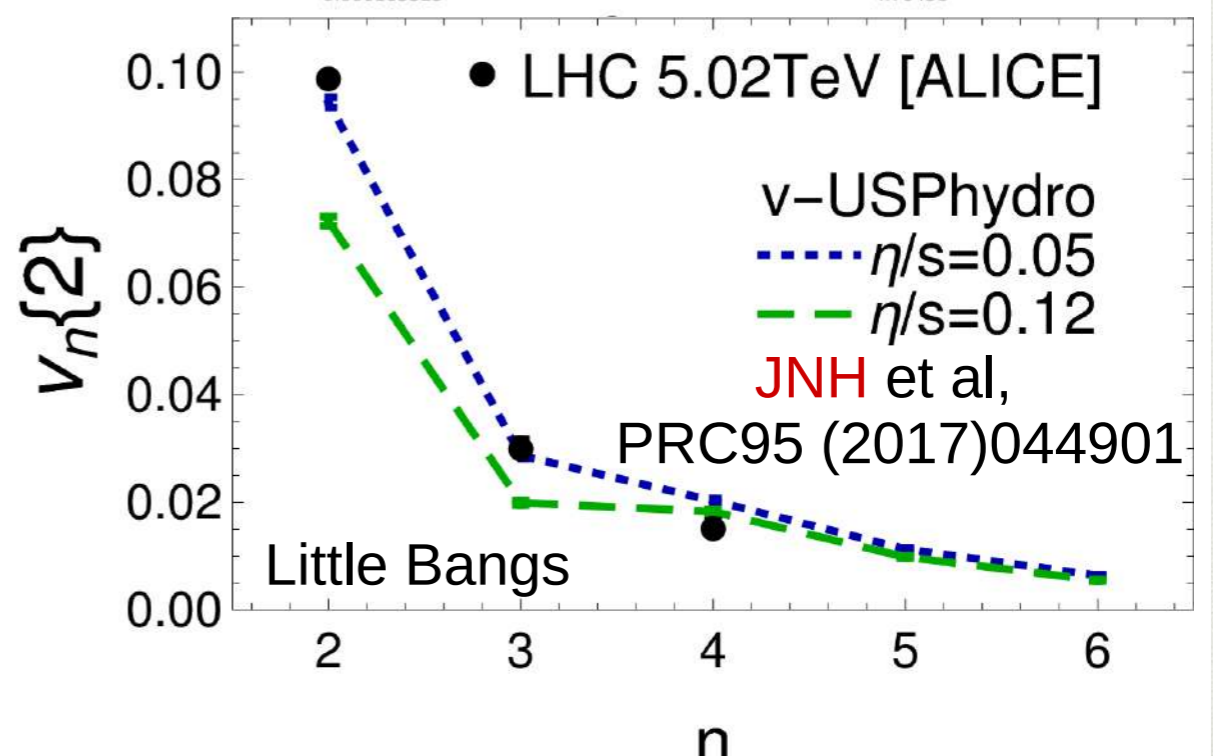
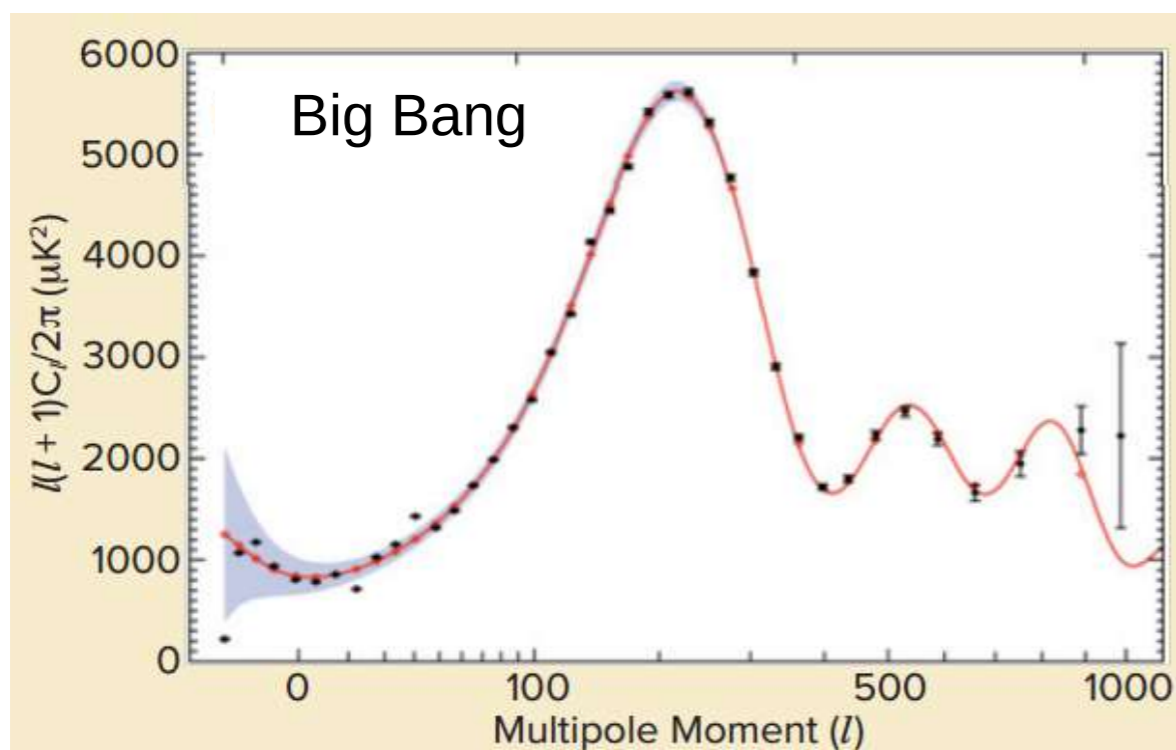
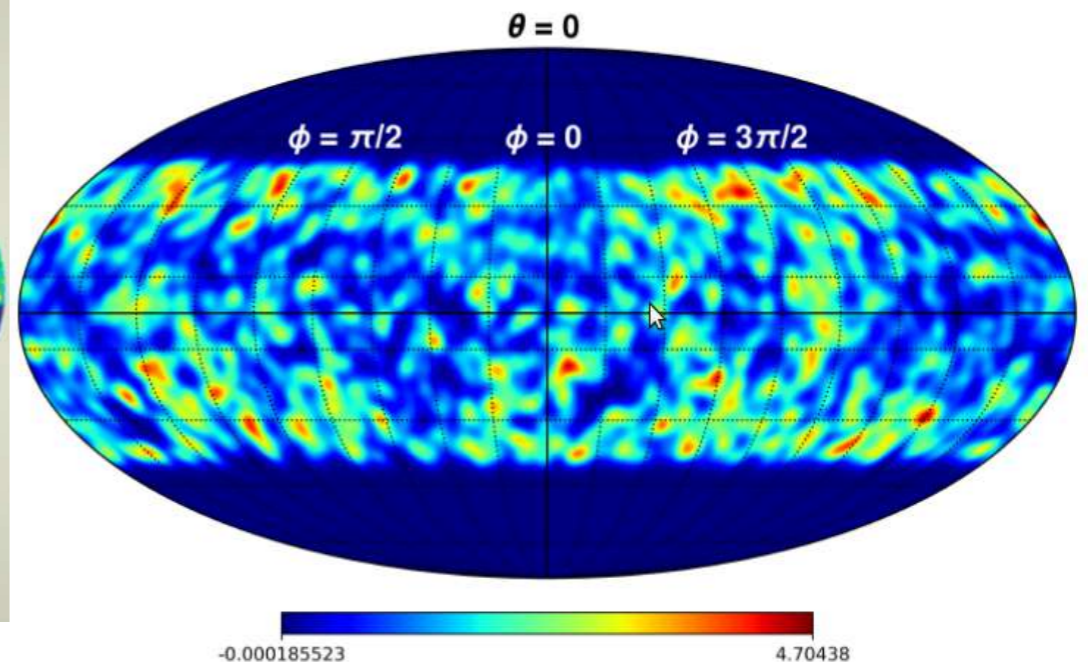


Collective flow: Flow harmonics, $v_n\{m\}$, are calculated by correlating $m=2$ to 8 particles \rightarrow collective behavior

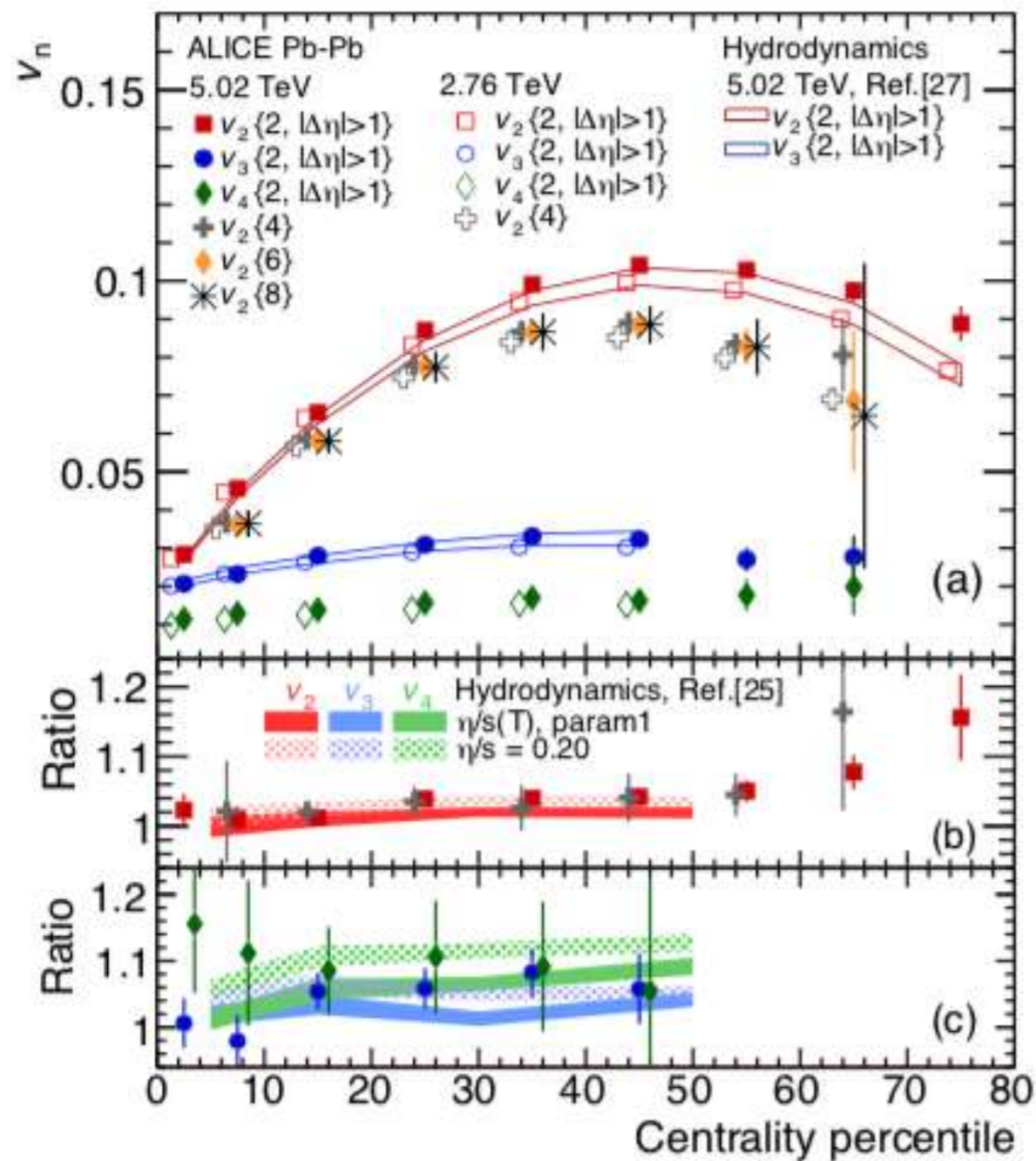
CMB vs. Heavy Ion Collisions



Vieira, Machado et al, *Phys.Rev. C*99 (2019) no.5, 054910



Precise predictions with hydrodynamics



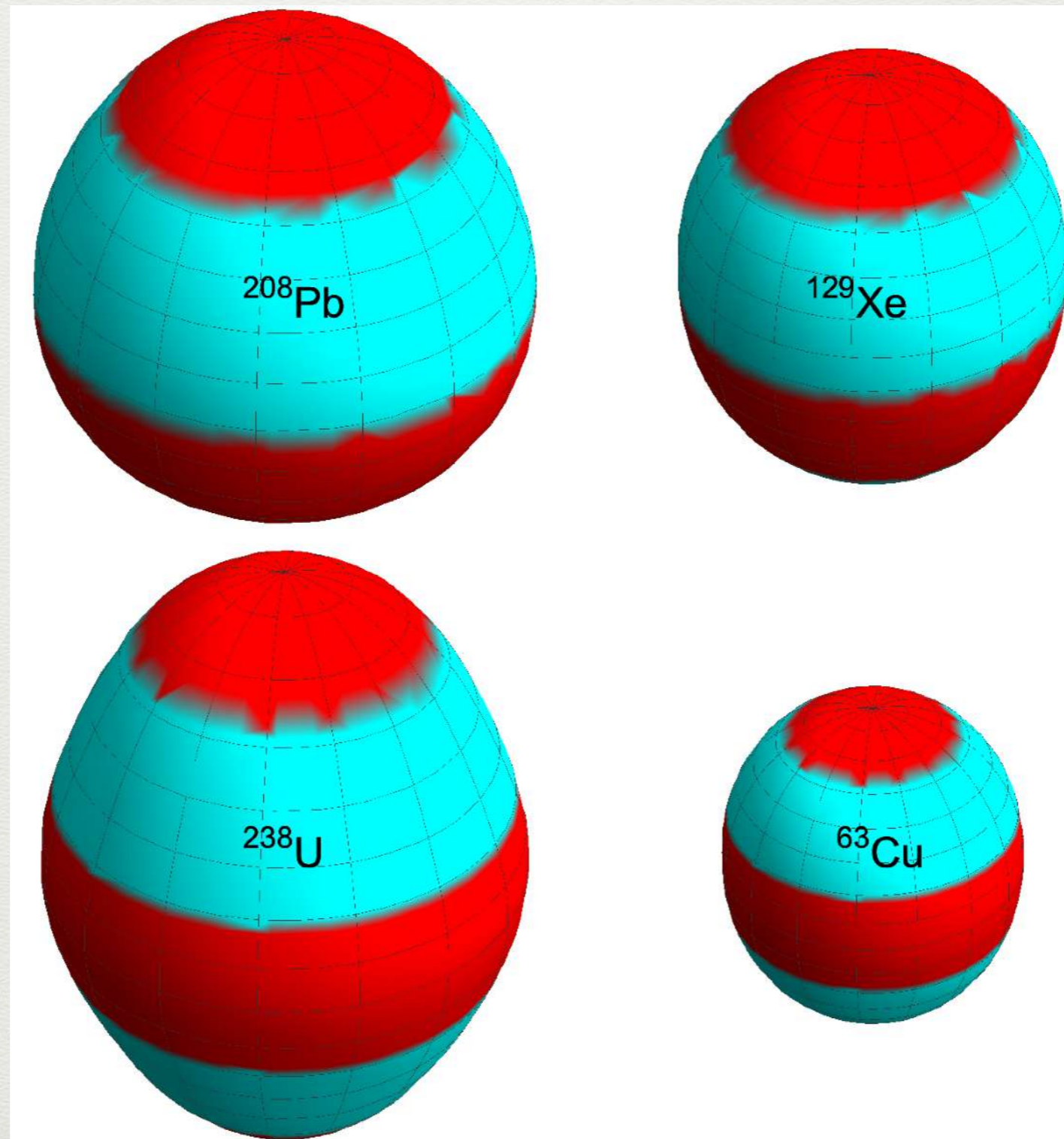
Hydrodynamic models can successfully make predictions at the $\sim 1\%$ level.

ALICE Phys.Rev.Lett. 116 (2016) no.13, 132302

v-USPhydro predictions: JNH et al, Phys.Rev. C93 (2016) no.3, 034912

EKRT predictions: Niemi et al, Phys. Rev. C 93, 014912 (2016)

Influence of different nuclei

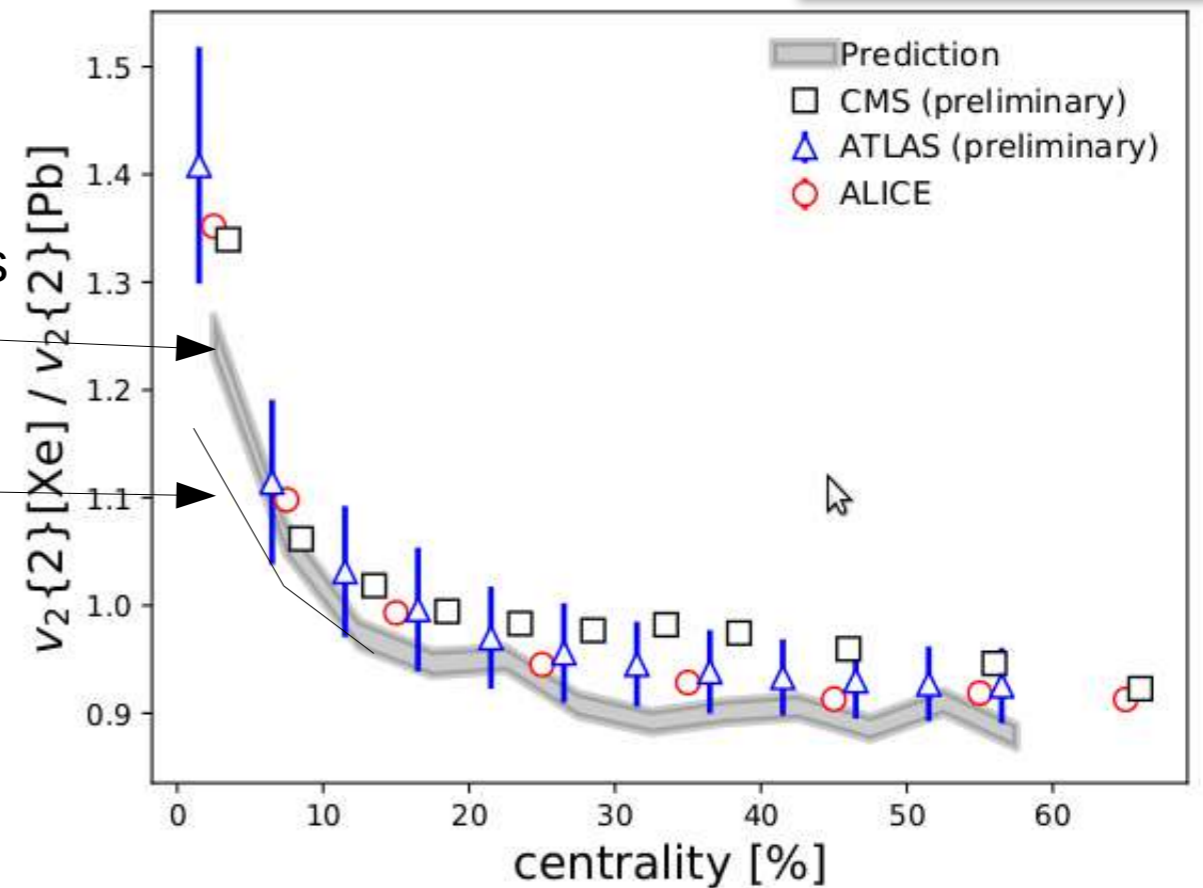
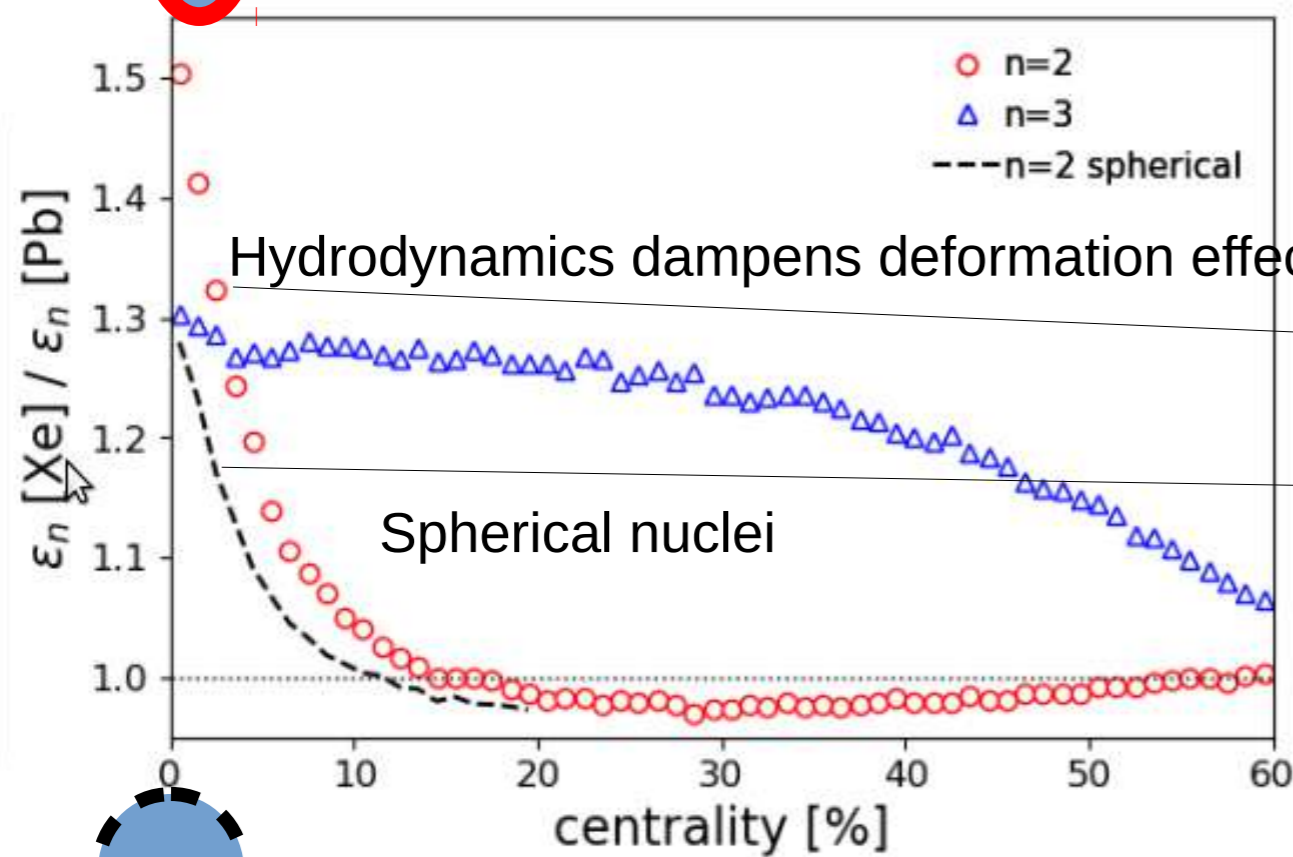


Finding a deformed nucleus in ^{129}Xe



Giuliano Giacalone
PhD Saclay

v-USPhydro sensitive to deformed nucleus
Giacalone, **JNH** et al. Phys. Rev. C 97,034904 (2018)

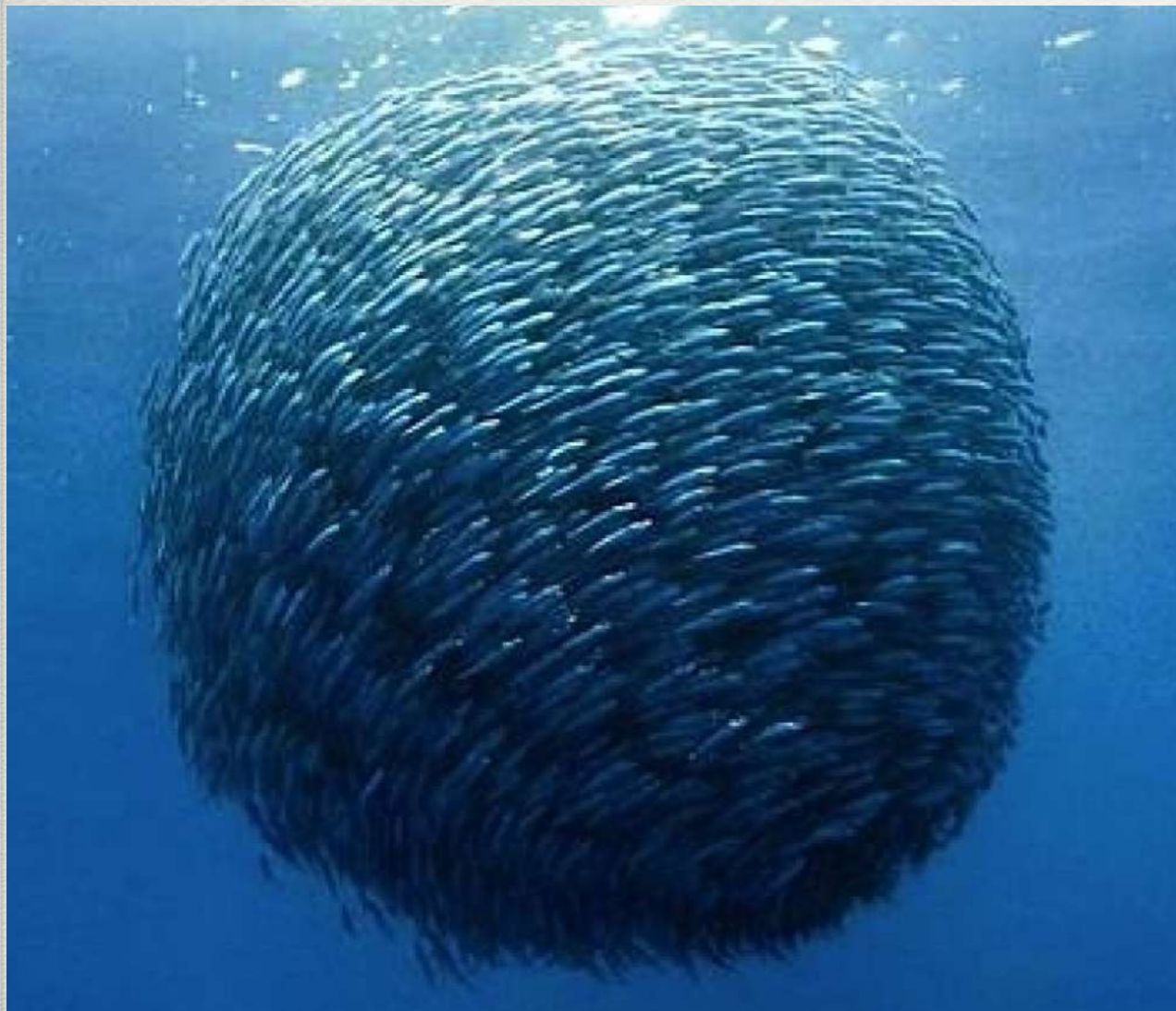


Deviation from experimental data: possible constraints on nuclear structure?

Limits on the smallest fluid

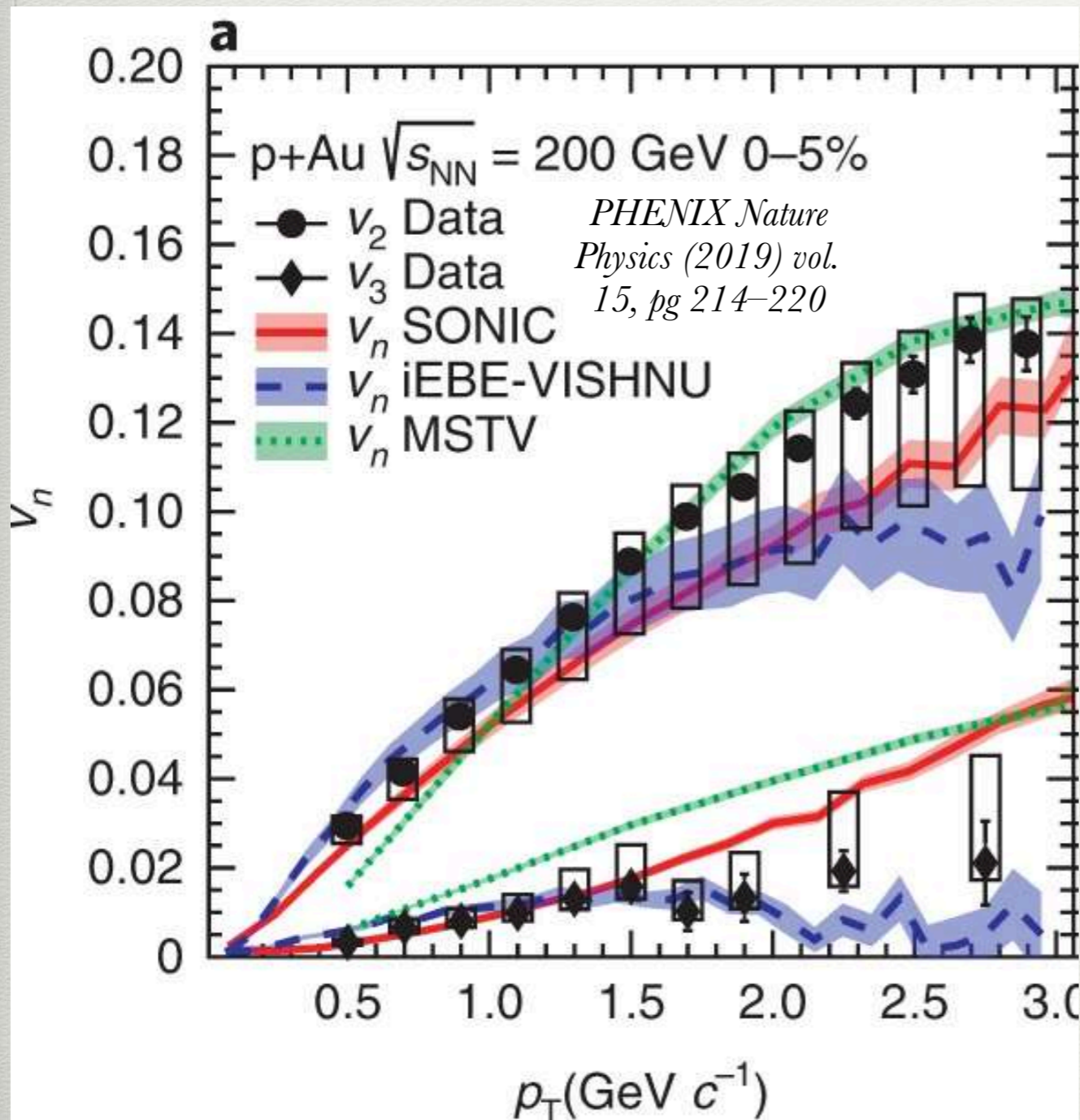
When do you have too few particles to use hydrodynamics?

$$Kn \sim \frac{\text{Small scale}}{\text{Large scale}} \sim 1$$

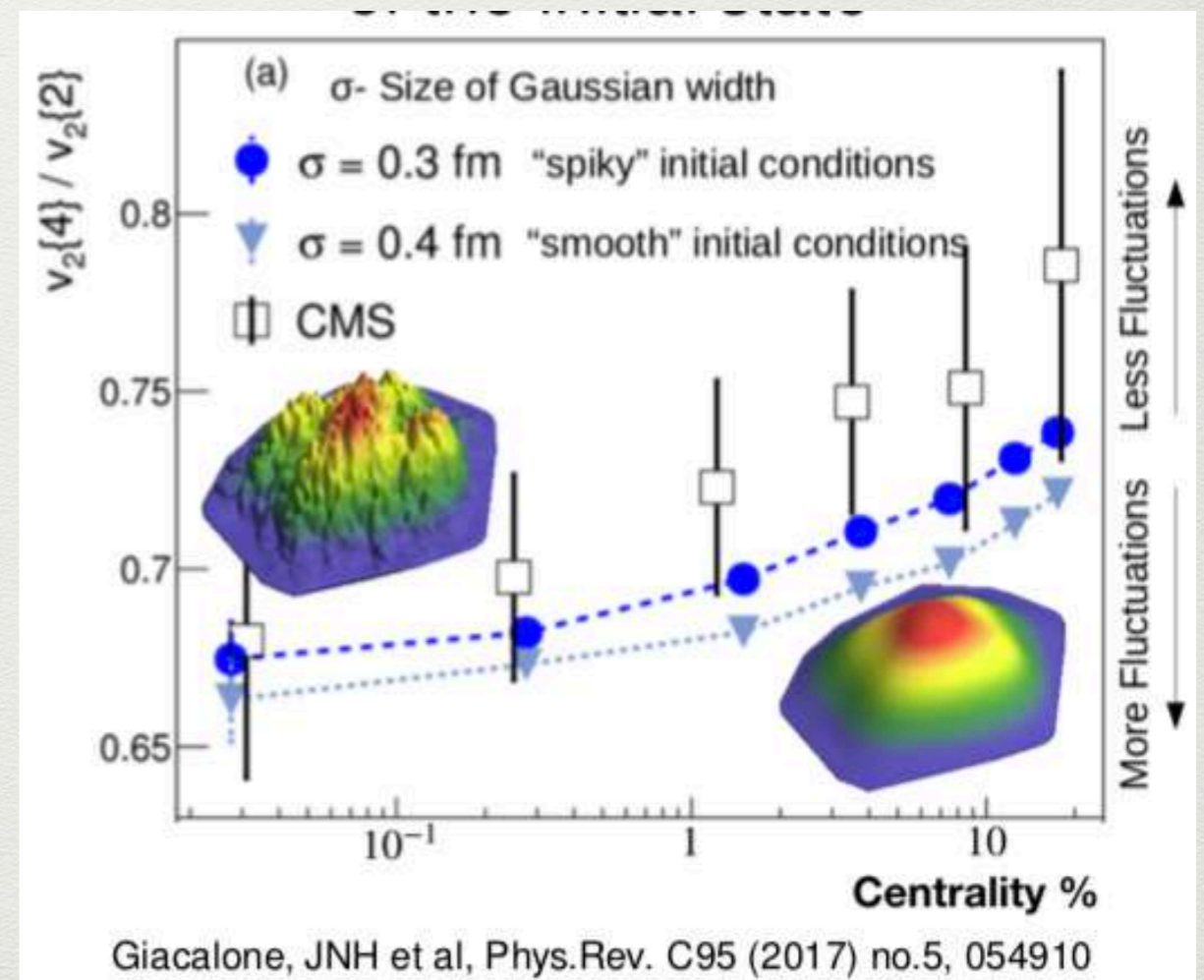


Experiment versus theory in small systems

Hydro matches data well



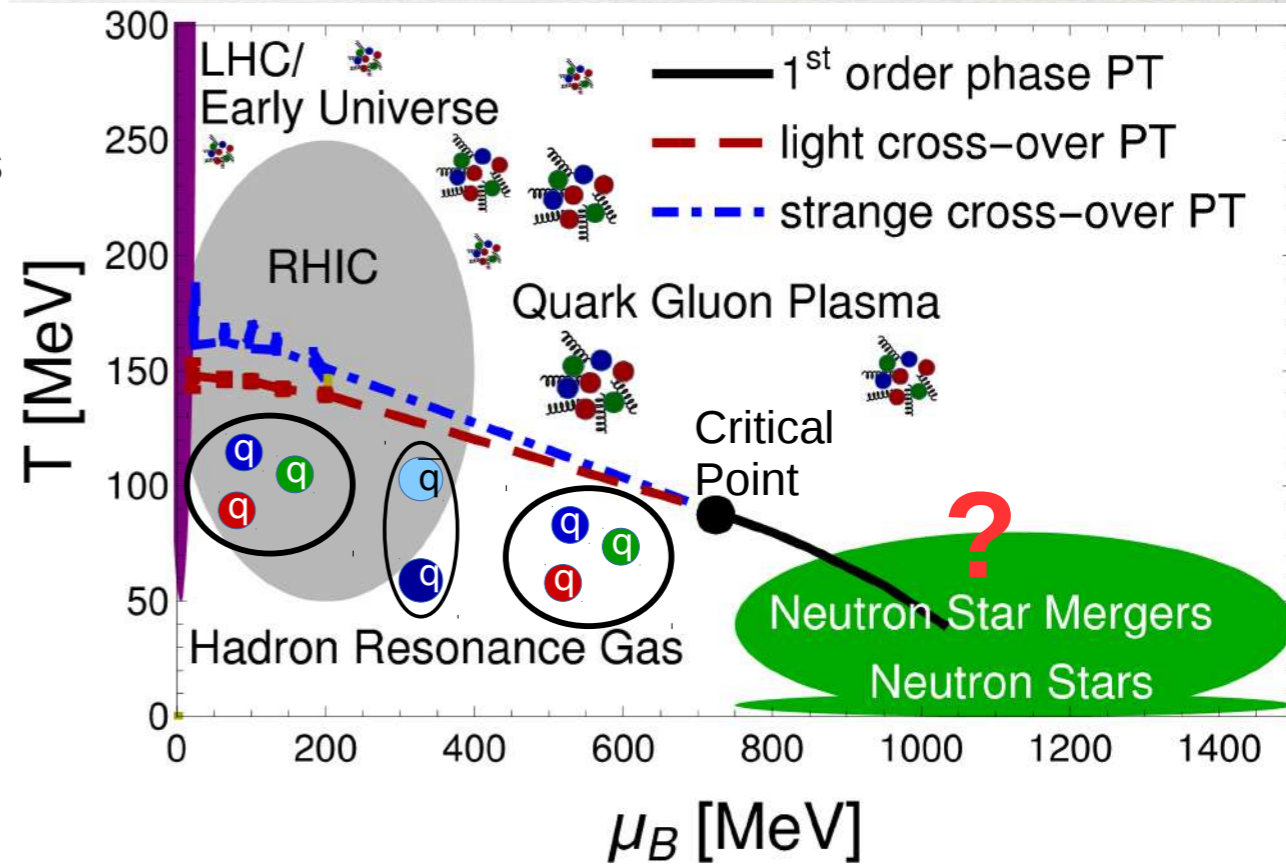
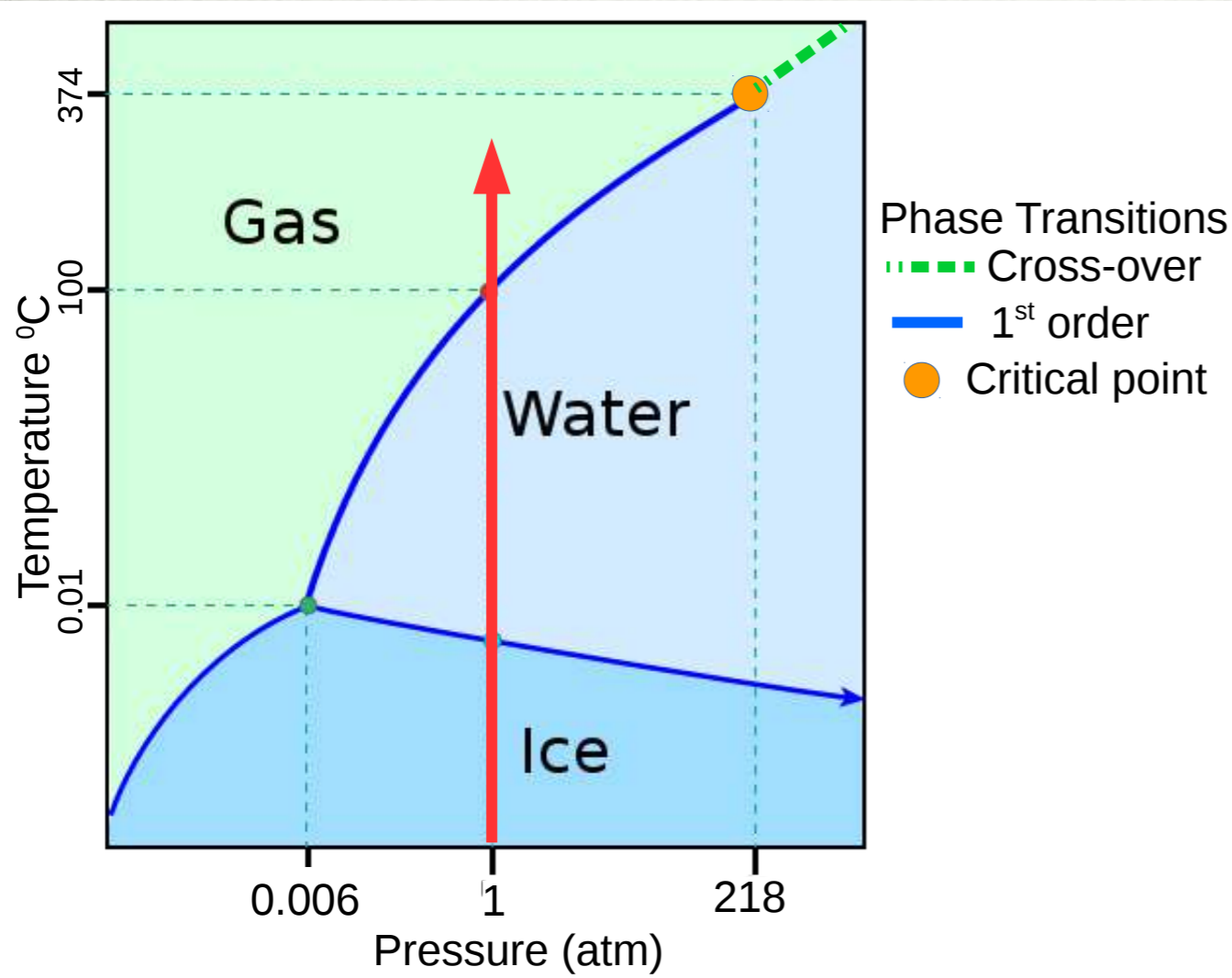
Questions remain on the initial conditions, applicability of hydrodynamics, and certain missing signals



Next frontiers of relativistic hydrodynamics

- Magnetohydrodynamics/Chiral Magnetic Effect
- **Conserved charges of QCD- baryon number, strangeness, and electric charge**
 - Each quark carries multiple charges!
- Source term in hydrodynamics for jets
- Critical fluctuations

QCD critical point



Search underway for the QCD critical point at the Beam Energy Scan II

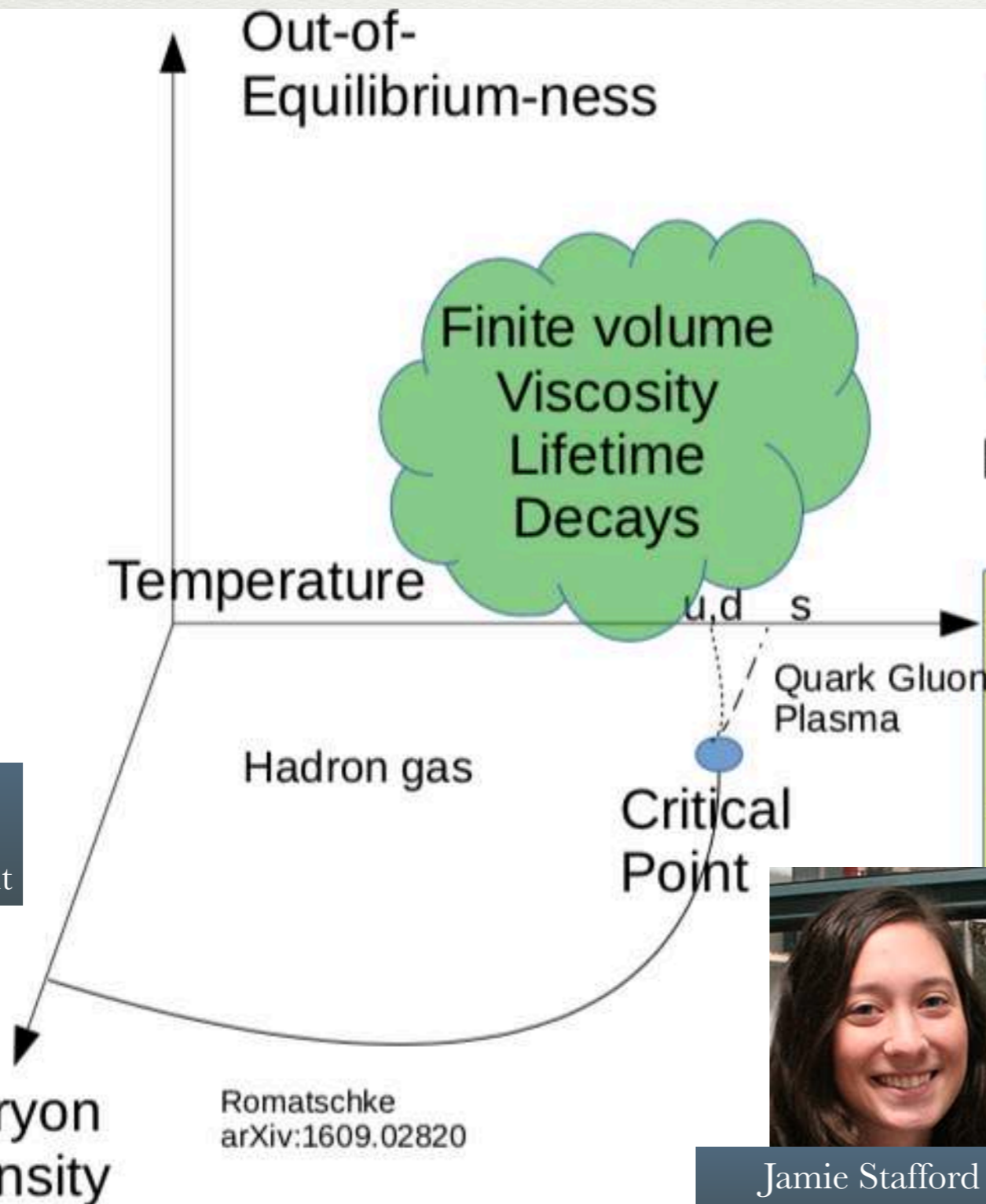
Out-of-equilibrium search for the CP



Paolo Parotto
PD Wuppertal



Debora Mroczek
REU Houston
UIUC PhD Student



Beam Energy Scan Theory
Collaboration

- EOS+3D Ising CP
Parotto, JNH, et al, Phys.Rev. C101 (2020) no.3, 034901
- Next, dynamical model with critical fluctuations and out-of-equilibrium effects+Lattice EOS



Jamie Stafford
PhD student Houston

BSQ EOS

Phys.Rev. C100 (2019) no.6, 064910
Phys.Rev. C100 (2019) no.2, 024907



Matt Sievert
Postdoc UIUC

M. Martinez
PD NCSU

D. Wertepny
PD Ben Gurion

Relativistic hydrodynamics with conserved charges (BSQ)

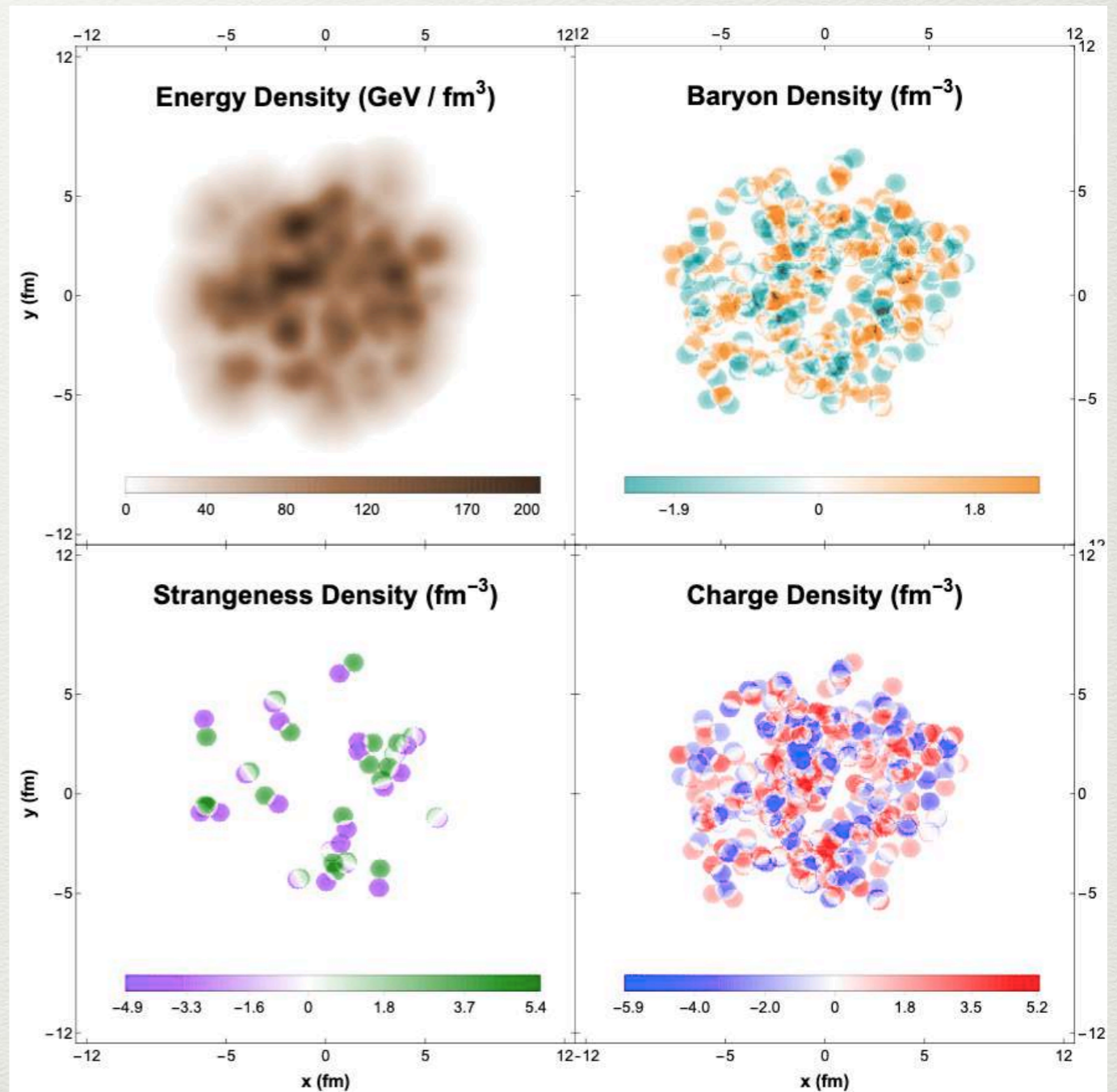
Example Initial Condition

ICCING - Initial conditions
with conserved charges (BSQ)

*Sievert, Martinez, Wertepny, JNH arXiv:1911.10272;
arXiv:1911.12454*

Previous initial
conditions assumed only
gluons \rightarrow ICCING
initializes quarks as well!

Next, incorporate Parton
Distribution Functions
(future Electron Ion
Collider)

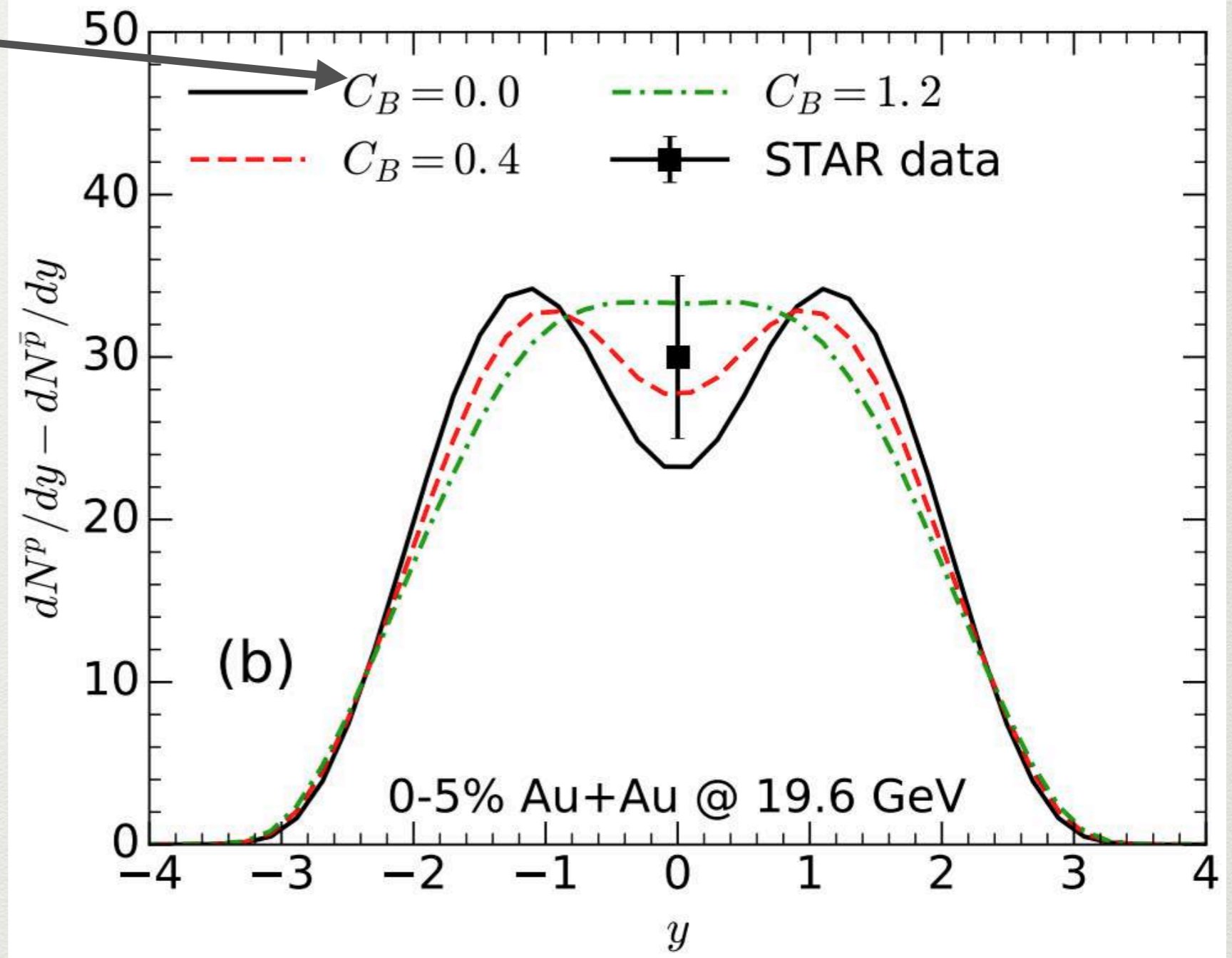


Larger diffusion - net-protons larger in center

Baryon diffusion parameter

BSQ diffusion the next stage

Fotakis et al, [arXiv:1912.09103](https://arxiv.org/abs/1912.09103)
Rose et al, [arXiv:2001.10606](https://arxiv.org/abs/2001.10606)



Denicol et al, *Phys. Rev. C* 98, 034916 (2018)



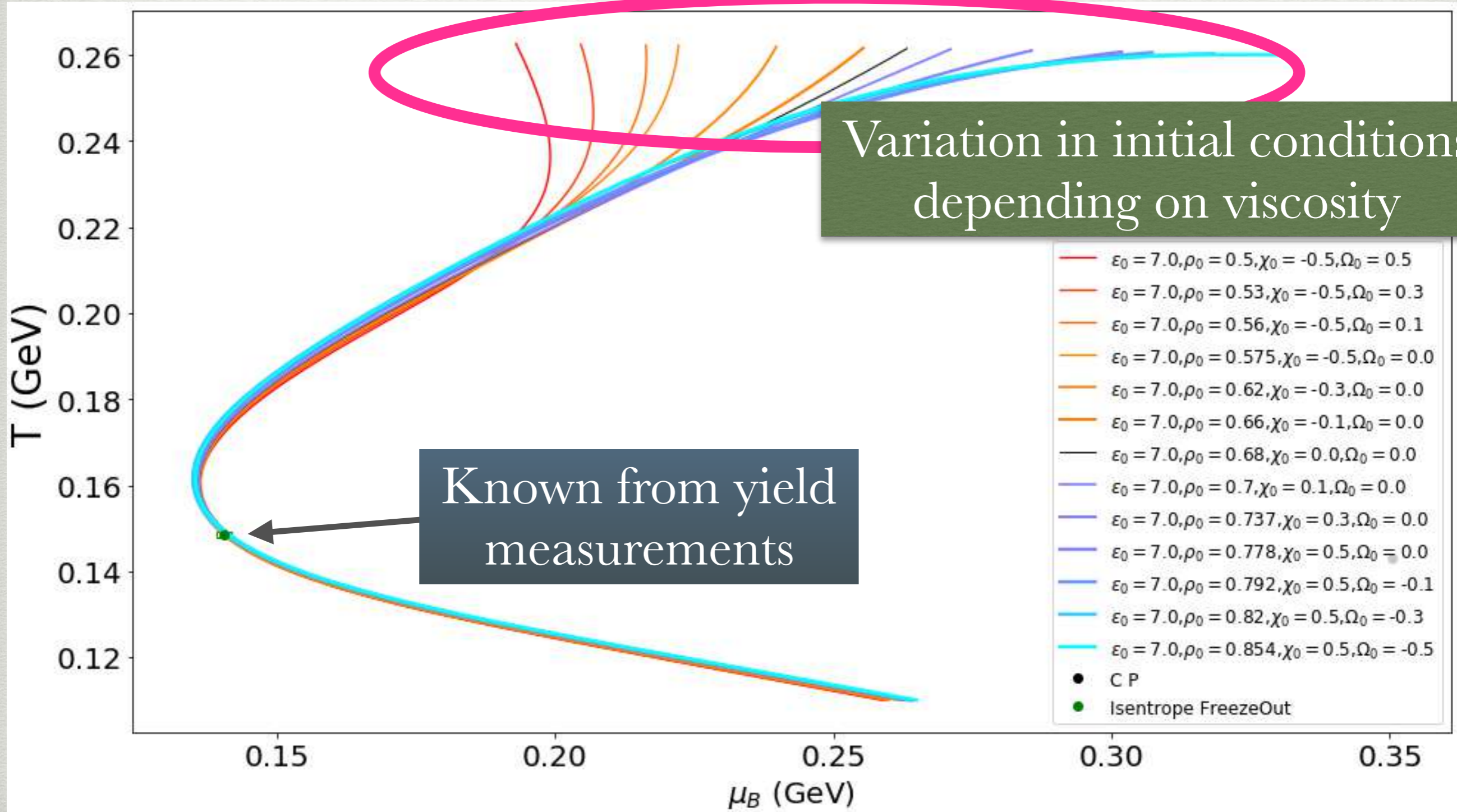
Travis Dore
PhD student



Emma McLaughlin
REU student

B(SQ) Hydrodynamics

Dore, McLaughlin, JNH, to appear soon



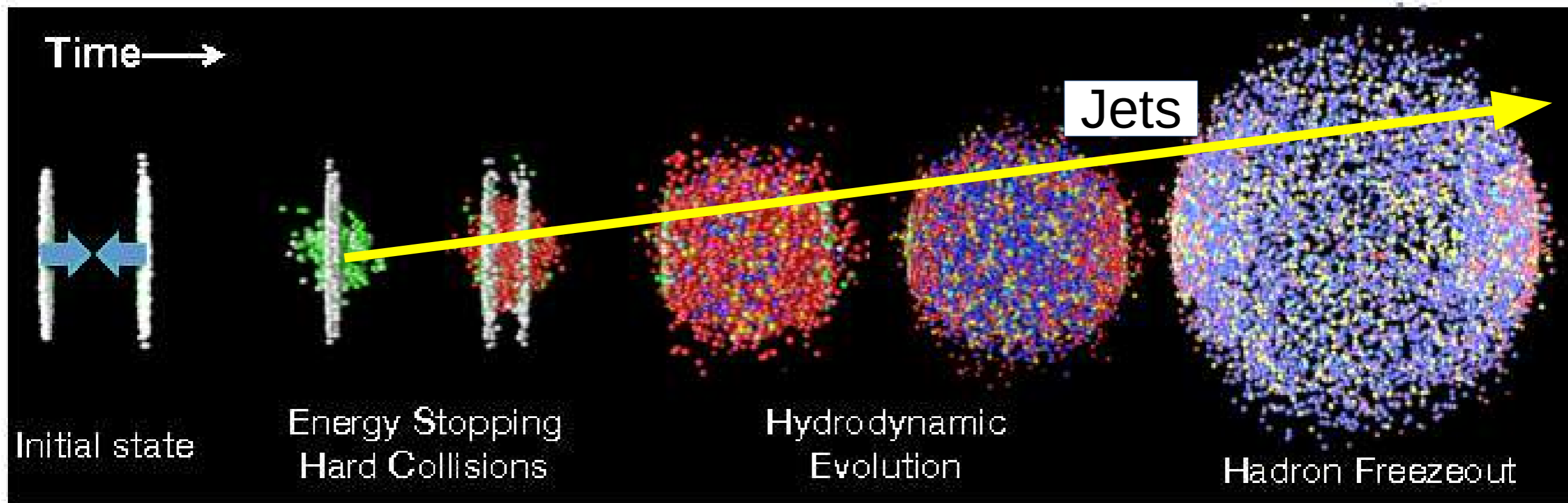


The Future

NEXT EXIT



Future Experiments

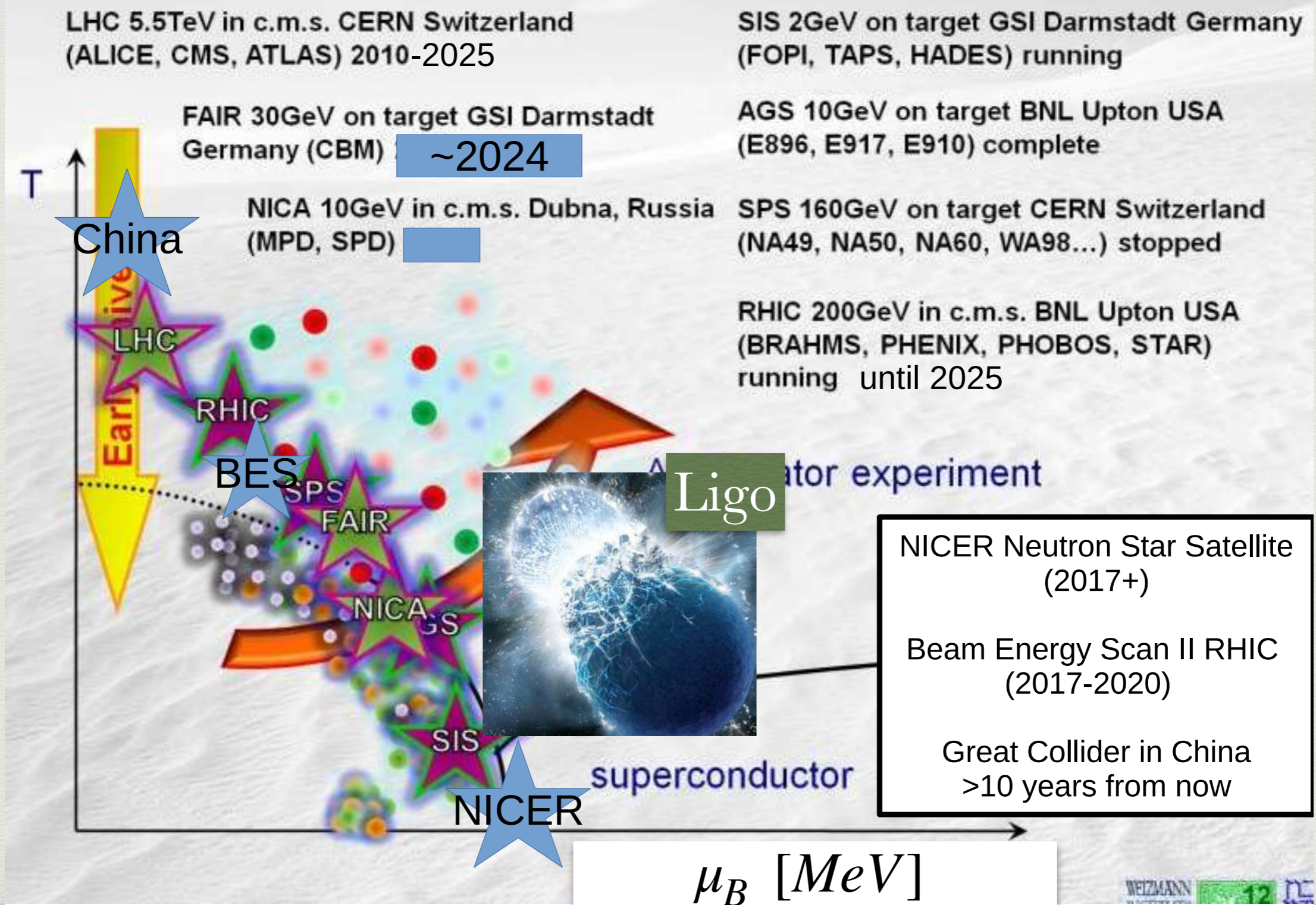


Electron Ion Collider
(EIC) -
Nucleon/Nuclei
Structure affect the
initial state (important
for small systems)
>2025

sPHENIX/LHC -
Jets probe shorter scales
i.e. a QGP microscope
2018-2025

Beam Energy Scan
(RHIC)/FAIR – High
baryon densities,
hadron gas phase
2018-2020, >2024

Mapping the QCD phase diagram



Summary

- Relativistic viscous hydrodynamics provides an extremely successful description of the Quark Gluon Plasma
- Much to come in the future on the addition of conserved charges, critical fluctuations, magnetohydrodynamics, jets coupled to hydrodynamics...
- STAR BESII, sPHENIX, EIC, FAIR, LIGO crucial to a further understanding of the Quark Gluon Plasma/
Quantum Chromodynamics