

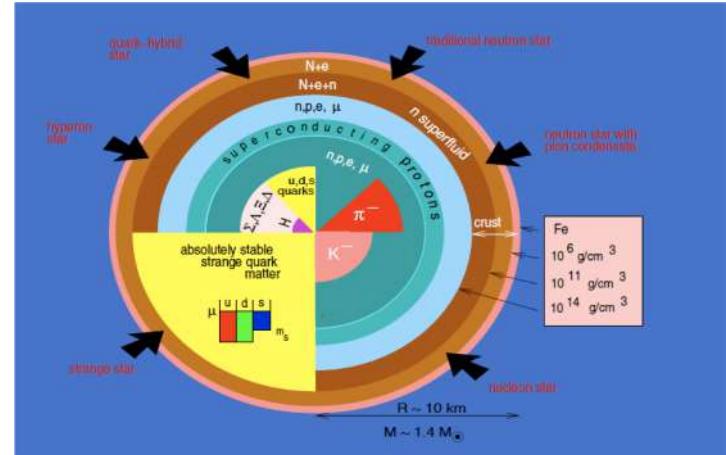
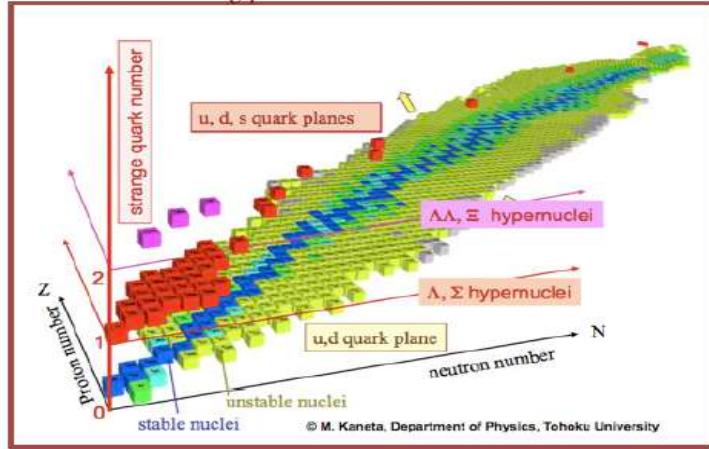
Strangeness in Nuclei and Neutron Stars

Laura Tolós

Institute of
Space Sciences

CSIC **IEEC**^R

based on
 Laura Tolos and Laura Fabbietti,
 Prog. Part. Nucl. Phys. 112 (2020) 103770, 2002.09223 [nucl-ex]



Strangeness Hyperons in Nuclei and Neutron Stars

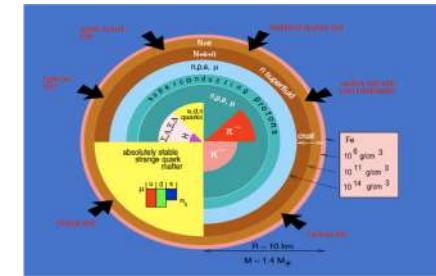
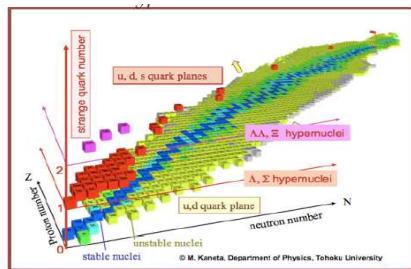
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Outline



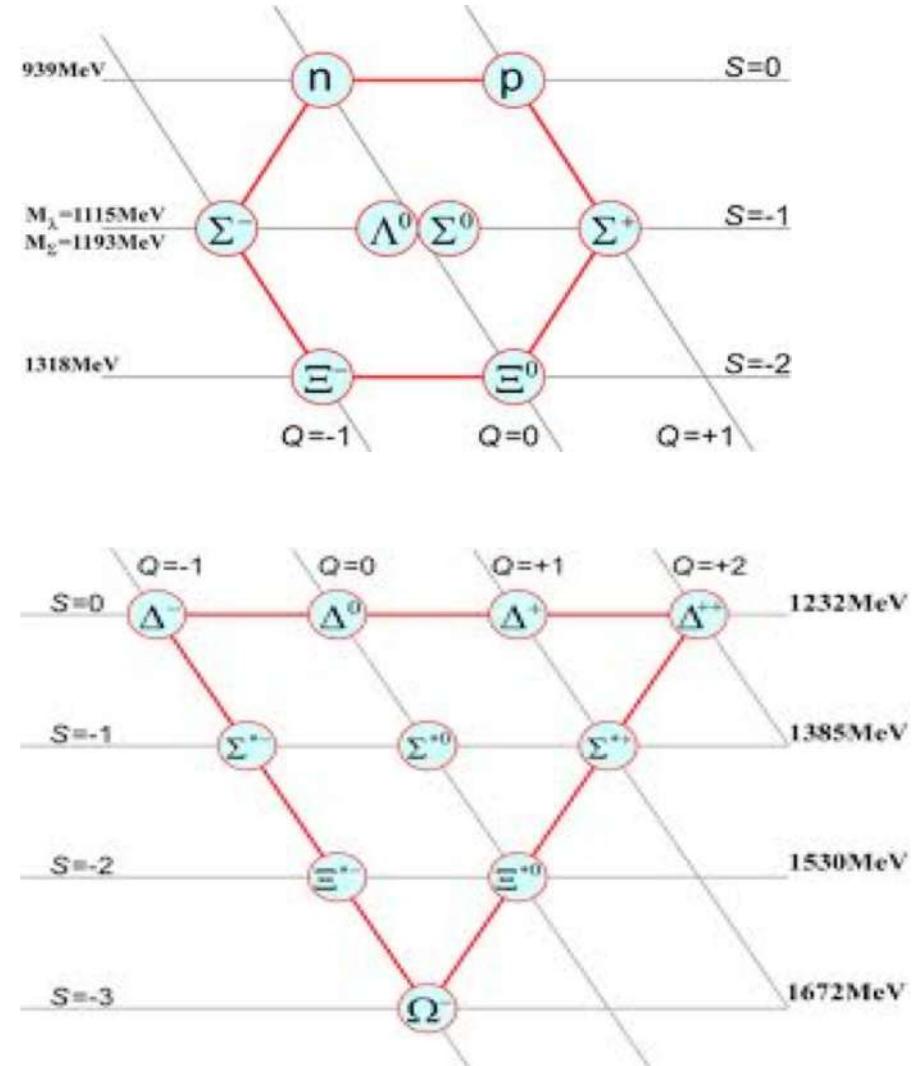
- Hyperons and where to find them
- YN and YY interactions
- Hypernuclei
- Hyperons in matter
- Hyperons and Neutron Stars
- Present and Future

Hyperons and where to find them

A **hyperon** is a baryon containing one or more strange quarks

Hyperon	Quarks	$I(J^P)$	Mass (MeV)
Λ	uds	0(1/2 ⁺)	1115
Σ^+	uus	1(1/2 ⁺)	1189
Σ^0	uds	1(1/2 ⁺)	1193
Σ^-	dds	1(1/2 ⁺)	1197
Ξ^0	uss	1/2(1/2 ⁺)	1315
Ξ^-	dss	1/2(1/2 ⁺)	1321
Ω^-	sss	0(3/2 ⁺)	1672

credit: I. Vidana

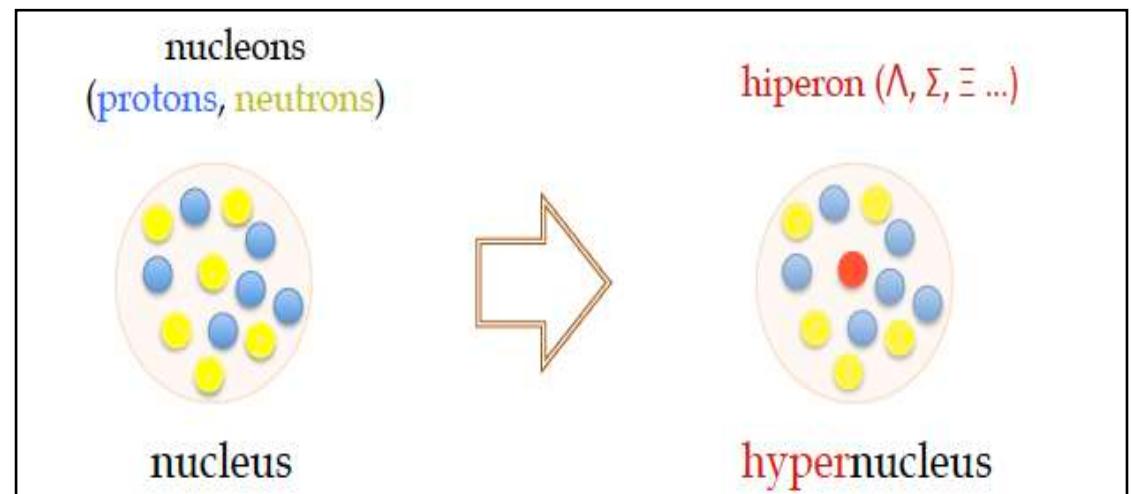


On Earth: Hypernuclei

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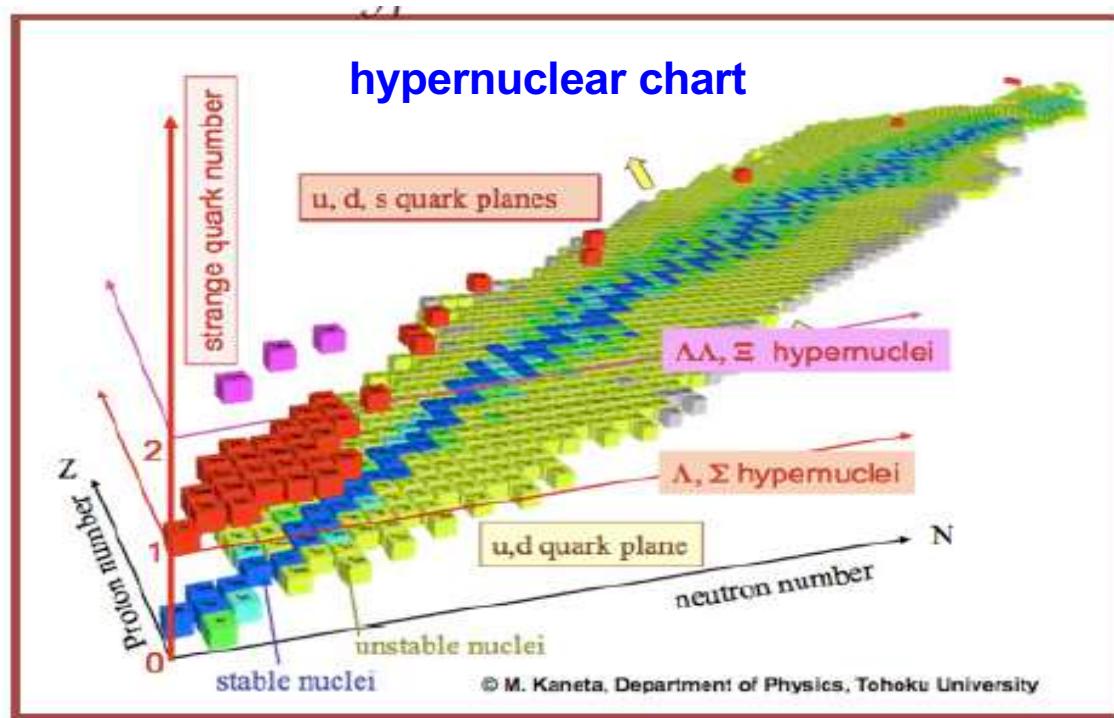
credit: I. Vidana



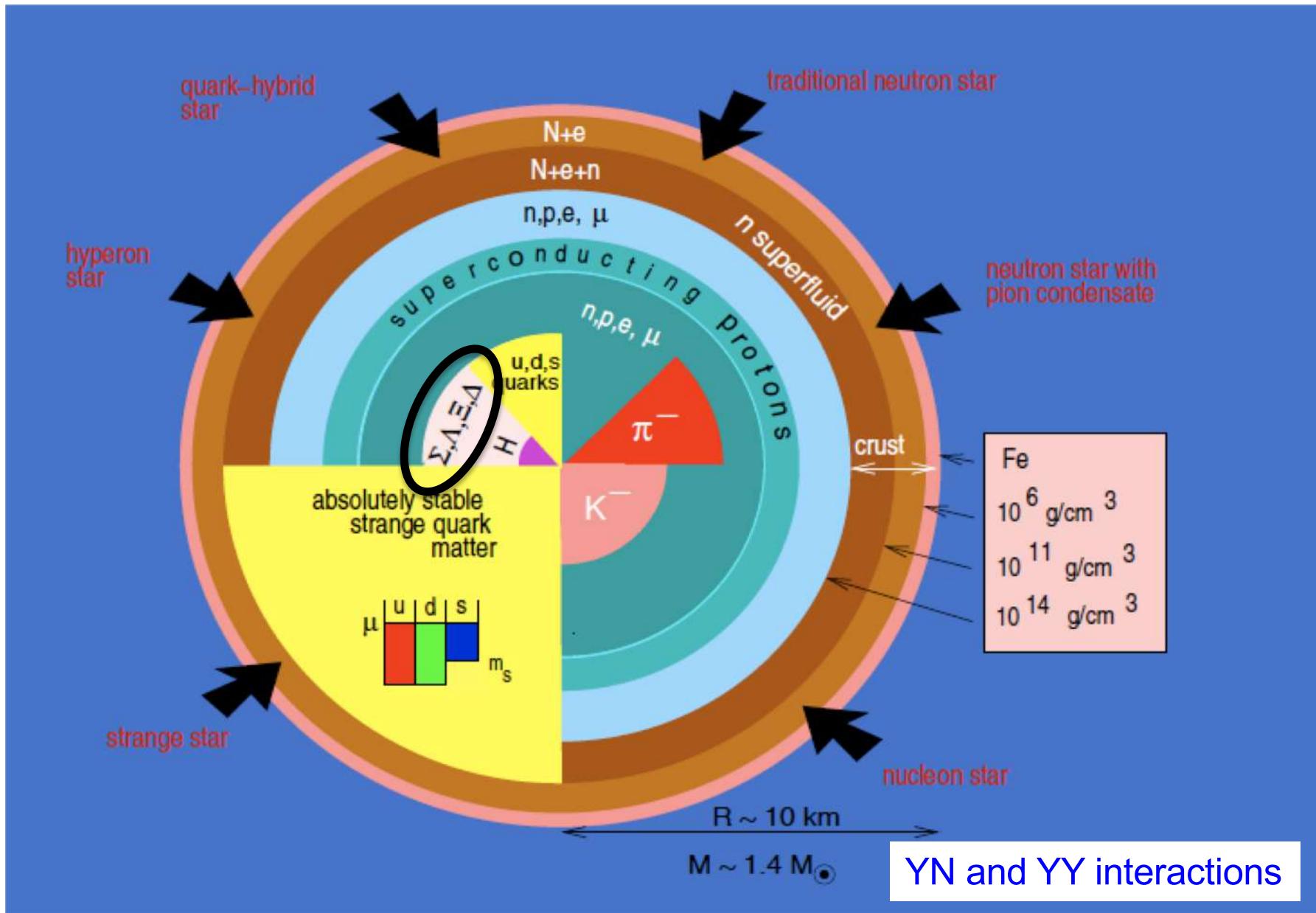
credit: A. Parreno

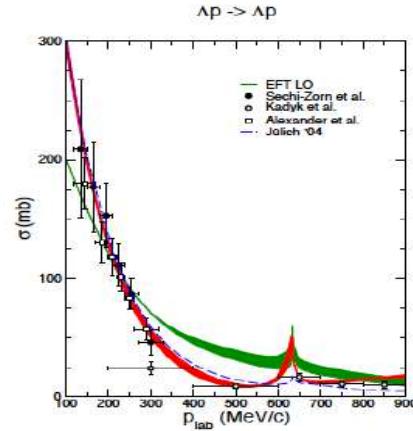
The study of hypernucleus allows for

- new spectroscopy
- information on strong and weak interactions between hyperons and nucleons



In Neutron Stars

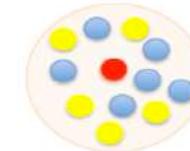




YN and YY interactions

- Study **strangeness** in nuclear physics
- Provide input for **hypernuclear physics** and astrophysics

hiperon ($\Lambda, \Sigma, \Xi \dots$)



hypernucleus

Scarce YN scattering data due to the short life of hyperons and the low-density beam fluxes

ΛN and ΣN : < 50 data points
 ΞN very few events

NN : > 5000 data
for $E_{\text{lab}} < 350$ MeV

Data from hypernuclei:

- more than 40 Λ -hypernuclei (ΛN attractive)
- few $\Lambda \Lambda$ - hypernuclei ($\Lambda \Lambda$ weak attraction)
- few Ξ -hypernuclei (ΞN attractive)
- no evidence of Σ -hypernuclei (ΣN repulsive)

Theoretical approaches to YN and YY

- Meson exchange models (Juelich/Nijmegen models)

To build YN and YY from a NN meson-exchange model imposing SU(3)_{flavor} symmetry

Juelich: Holzenkamp, Holinde, Speth '89; Haidenbauer and Meißner '05

Nijmegen: Maesen, Rijken, de Swart '89; Rijken, Nagels and Yamamoto '10

- Chiral effective field theory approach (Juelich-Bonn-Munich group)

To build YN and YY from a chiral effective Lagrangian similarly to NN interaction

Juelich-Bonn-Munich: Polinder, Haidenbauer and Meißner '06; Haidenbauer, Petschauer, Kaiser, Meißner, Nogga and Weise '13
Kohno '10; Kohno '18

- Quark model potentials

To build YN and YY within constituent quark models

Fujiwara, Suzuki, Nakamoto '07

Garcilazo, Fernandez-Carames and Valcarce '07 '10

- $V_{\text{low } k}$ approach

To calculate a “universal” effective low-momentum potential for YN and YY using RG techniques

Schaefer, Wagner, Wambach, Kuo and Brown '06

- Lattice calculations (HALQCD/NPLQCD)

To solve YN and YY interactions on the lattice

HALQCD: Ishii, Aoki, Hatsuda '07; Aoki, Hatsuda and Ishii '10; Aoki et al '12

NPLQCD: Beane, Orginos and Savage '11; Beane et al '12

Theoretical approaches to YN and YY

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NPLQCD: Beane, Orginos and Savage '11; Beane et al '12

YN (and YY) meson-exchange models

Built from a NN meson-exchange model imposing SU(3)_{flavor} symmetry

NIJMEGEN

(Nagels, Rijken, de Swart, Timmermans, Maessen..)

- ✓ Based on Nijmegen NN potential
- ✓ Momentum and Configuration Space
- ✓ Exchange of pseudoscalar, vector and scalar nonets
- ✓ SU(3) symmetry to relate YN to NN vertices
- ✓ Gaussian form factors

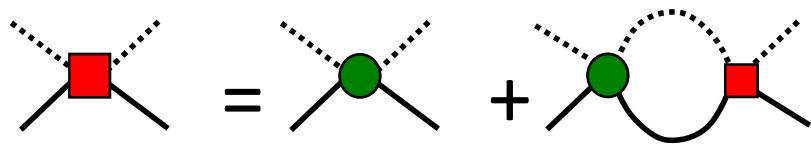
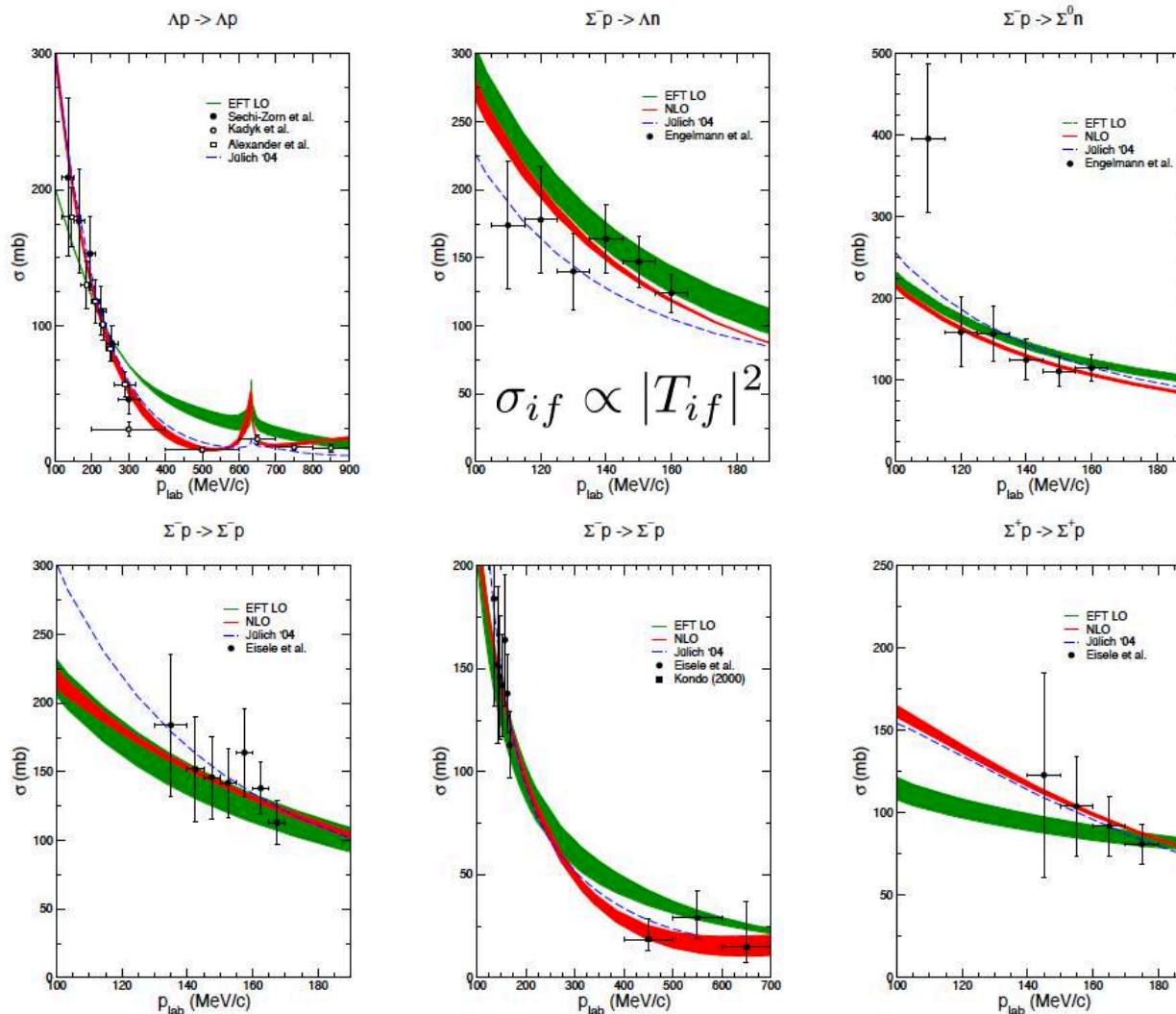
JUELICH

(Holzenkamp, Reube, Holinde, Speth, Haidenbauer, Meissner, Melnitchouck..)

- ✓ Based on Bonn NN potential
- ✓ Momentum Space, Full Energy Dependence & Non-localities
- ✓ Exchange of single mesons and higher order processes
- ✓ SU(6) symmetry to relate YN to NN vertices
- ✓ Dipolar form factors

ΛN and ΣN scattering

LO: H. Polinder, J.H., U. Mei β nner, NPA 779 (2006) 244
 NLO: J.H., N. Kaiser, et al., NPA 915 (2013) 24
 J \ddot{u} lich '04: J.H., U.-G. Mei β nner, PRC 72 (2005) 044005

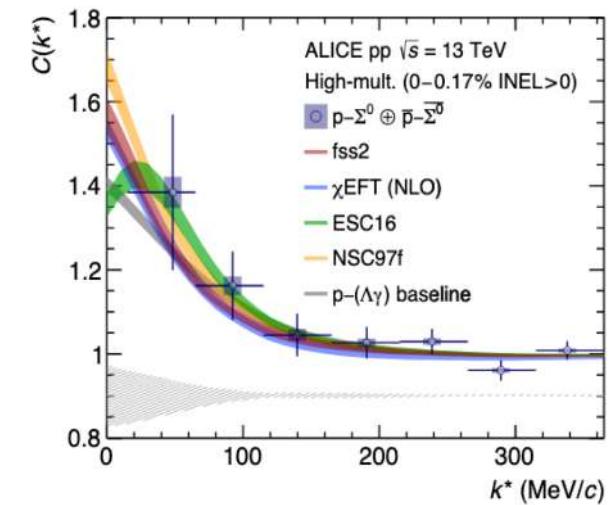


$$T = V + V \frac{1}{E_0 - H_0 + i\eta} T$$

New results from
femtoscopy for $\Sigma^0 p$

$$C(k^*) = \mathcal{N} \times \frac{N_{\text{same}}(k^*)}{N_{\text{mixed}}(k^*)}$$

$$k^* = \frac{1}{2} \times |\mathbf{p}_1^* - \mathbf{p}_2^*|$$



S. Acharya et al. 2019

YN (and YY) interactions in χ EFT

Baryon-Baryon interaction in SU(3) χ EFT a la Weinberg (1990);

- power counting allowing for a systematic improvement by going to higher order
- derivation of two- and three-baryon forces in a consistent way

Degrees of freedom: octet of baryons (N, Λ, Σ, Ξ) & pseudoscalar mesons (π, K, η)

Diagrams: pseudoscalar-meson exchanges and contact terms

credit: Haidenbauer

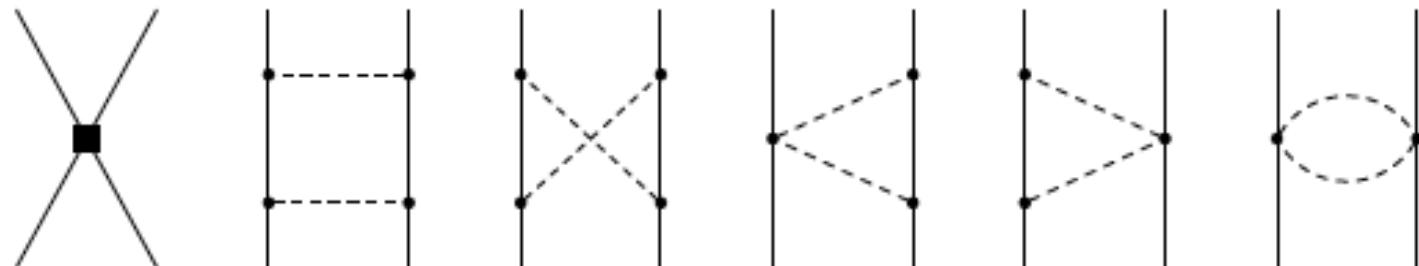
LO :



$$\nu = 2 - B + 2L + \sum_i v_i \Delta_i ,$$
$$\Delta_i = d_i + \frac{1}{2} b_i - 2 ,$$

B: number of incoming (outgoing) baryons
L: number of Goldstone boson loops
 v_i : number of vertices with dimension Δ_i
 d_i : derivatives
 b_i : number of internal baryons at vertex

NLO :

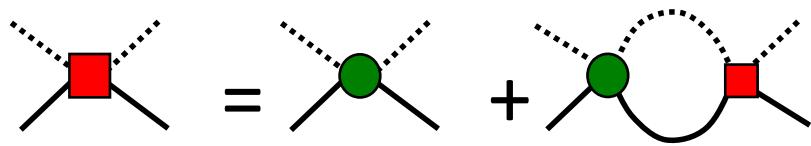
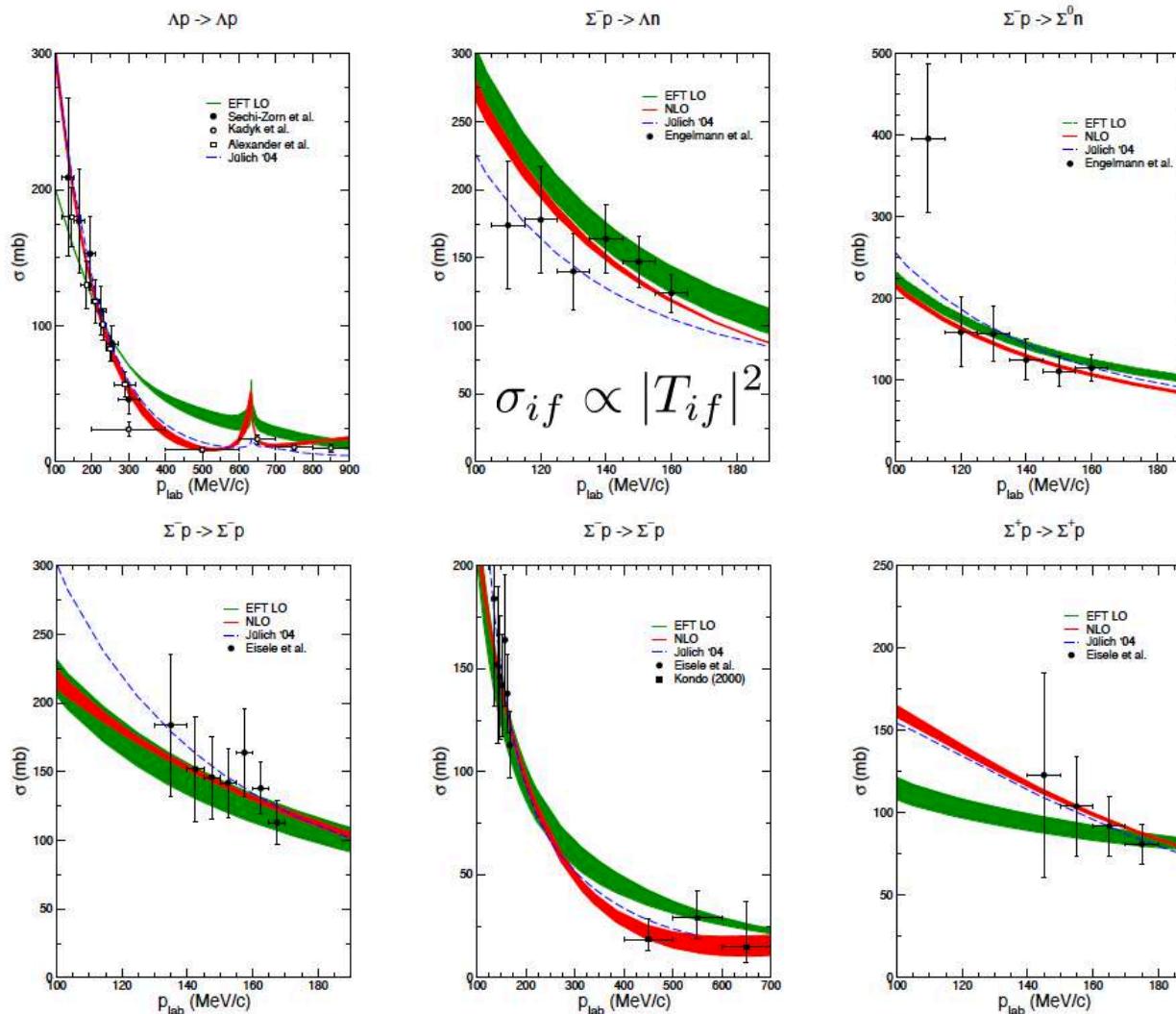


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ΛN and ΣN scattering

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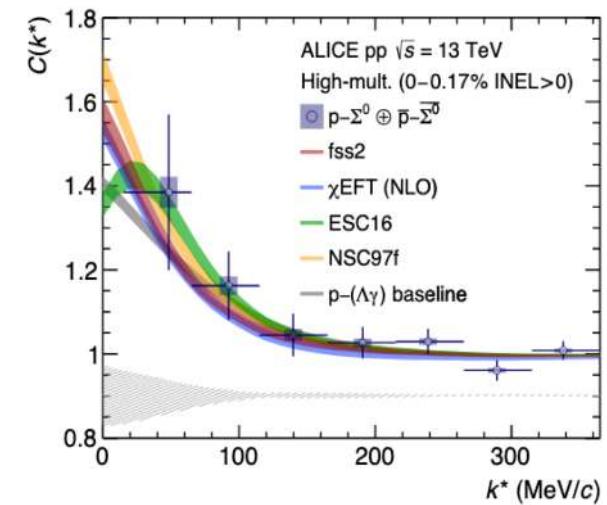


$$T = V + V \frac{1}{E_0 - H_0 + i\eta} T$$

New results from
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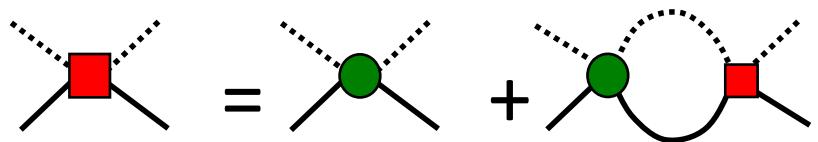
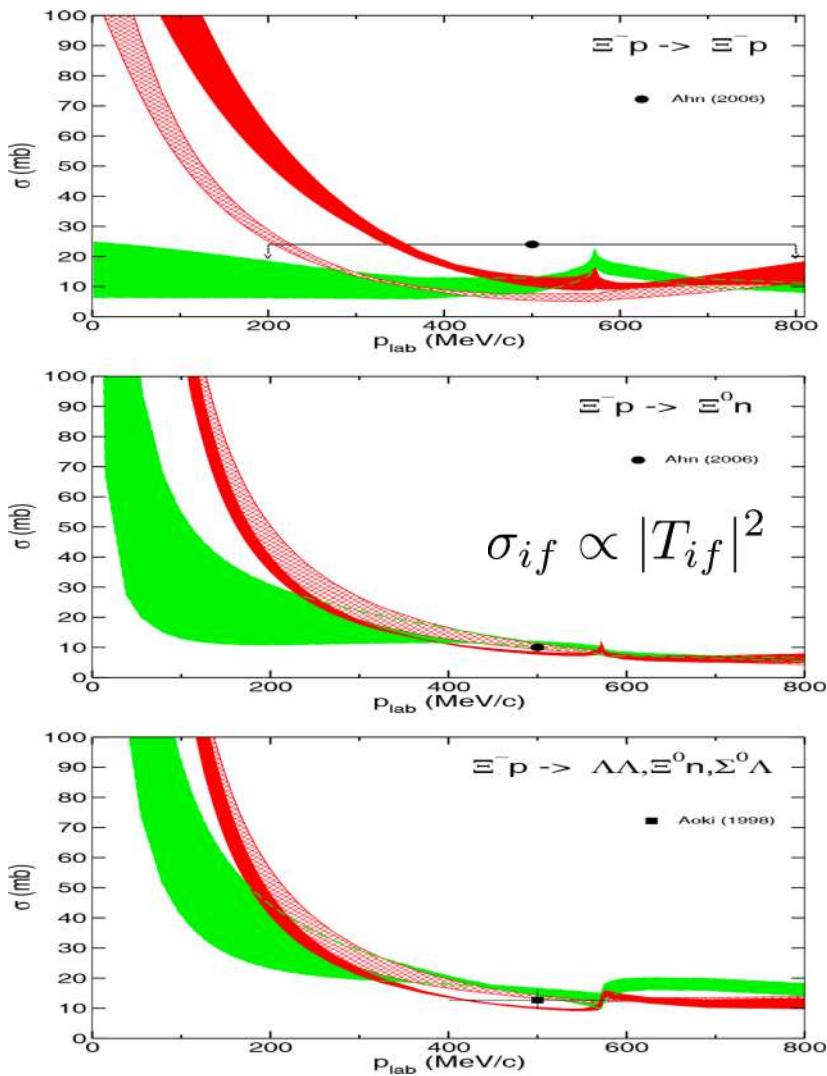
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$$k^* = \frac{1}{2} \times |\mathbf{p}_1^* - \mathbf{p}_2^*|$$



S. Acharya et al. 2019

ΞN scattering

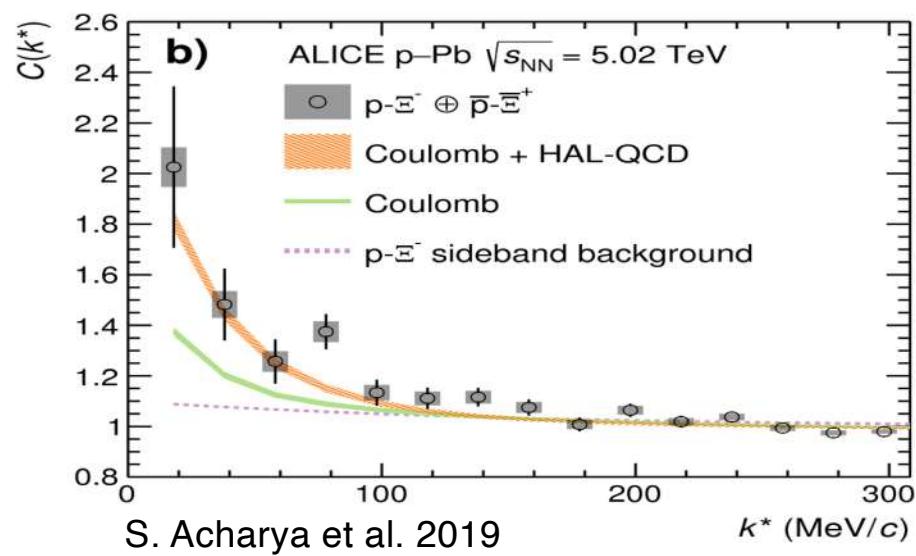


$$T = V + V \frac{1}{E_0 - H_0 + i\eta} T$$

ΞN cross sections are small

J. Haidenbauer and
U.G. Meißner EPJA 55 (2019) 23

Scarce experimental information.
New results from femtoscopy



S. Acharya et al. 2019

Hypernuclei

Λ hypernuclei

PRODUCTION REACTIONS

Strangeness exchange: $n(K^-, \pi^-)\Lambda$ CERN, BNL, KEK
 $p(K^-, \pi^\pm)\Sigma^\mp$ FINUDA@DAPHNE

$K^- \{ \bar{u} \leftrightarrow \bar{d} \} \pi^-$

$n \{ d \leftrightarrow u \} \Lambda$

$K^- + {}^{12}C \rightarrow \pi^- + {}_{\Lambda}^{12}C$

Associated production: $n(\pi^+, K^+)\Lambda$ BNL, KEK

$\pi^+ \{ u \leftrightarrow \bar{d} \} K^+$

$n \{ d \leftrightarrow u \} \Lambda$

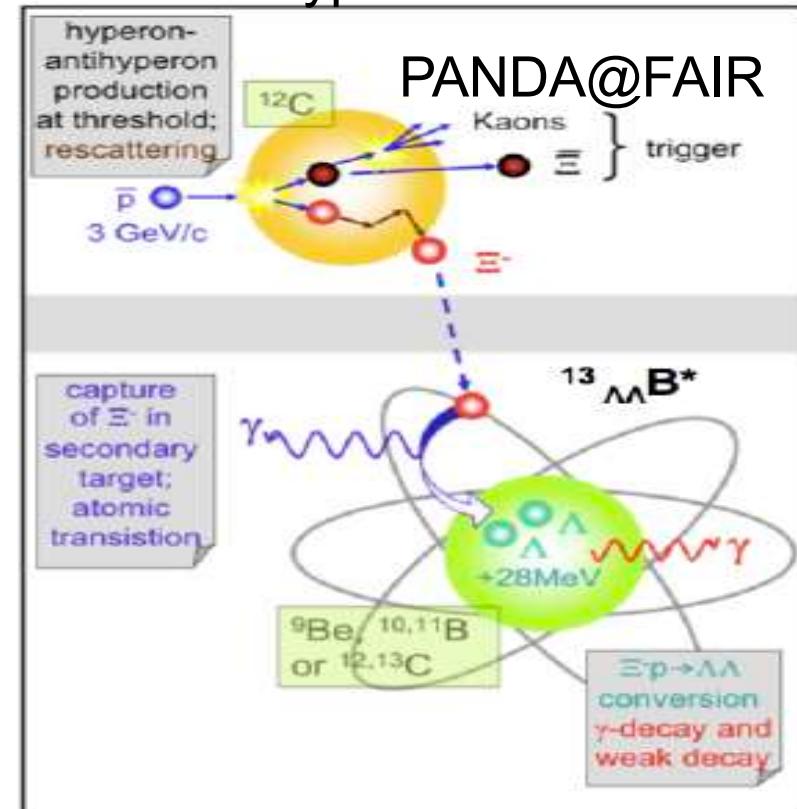
$\pi^+ + {}^{12}C \rightarrow K^+ + {}_{\Lambda}^{12}C$

Electroproduction: $p(\gamma, K^+)\Lambda$ Jlab, MAMI-C
 $p(e, e' K^+)\Lambda$

$e + {}^{12}C \rightarrow e' + K^+ + {}_{\Lambda}^{12}B$

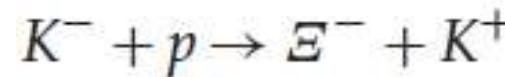
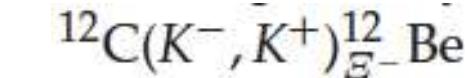
credit: A. Parreno

Double Λ hypernuclei



credit: A. Sanchez-Lorente

Also Ξ hypernuclei @ BNL, KEK



Laboratories:

BNL, CERN, KEK, JLab, DAΦNE, GSI, FAIR



Reactions:

Emulsion data

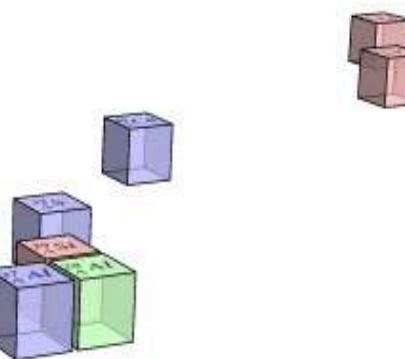
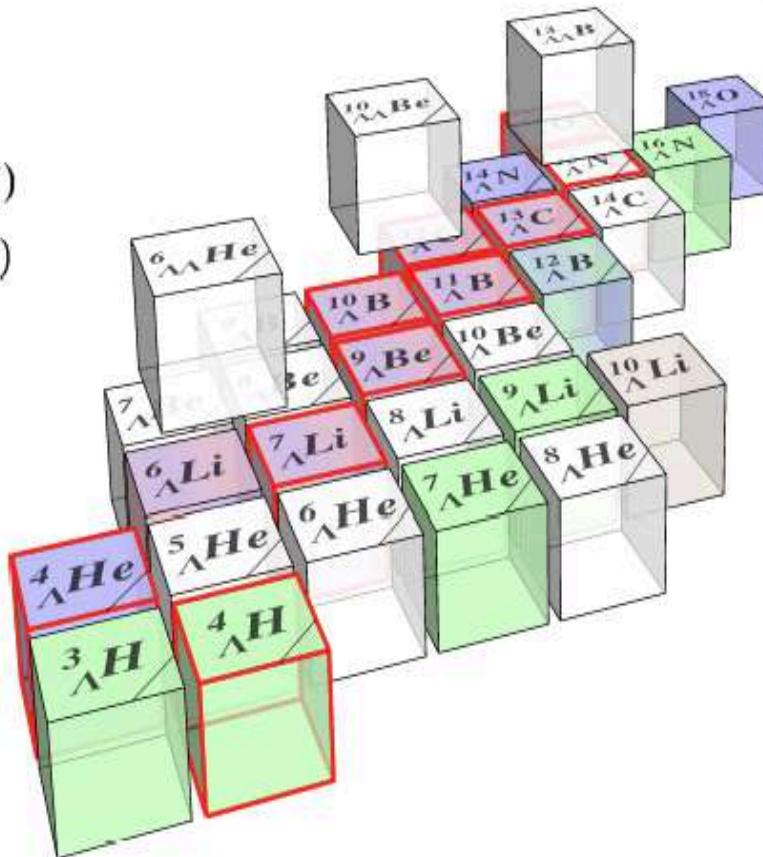
γ -ray data

(K^-, π^-)
 (K_{stop}^-, π^-)
 (K_{stop}^-, π^0)

$(e, e' K^+)$

(π^+, K^+)

(π^-, K^+)

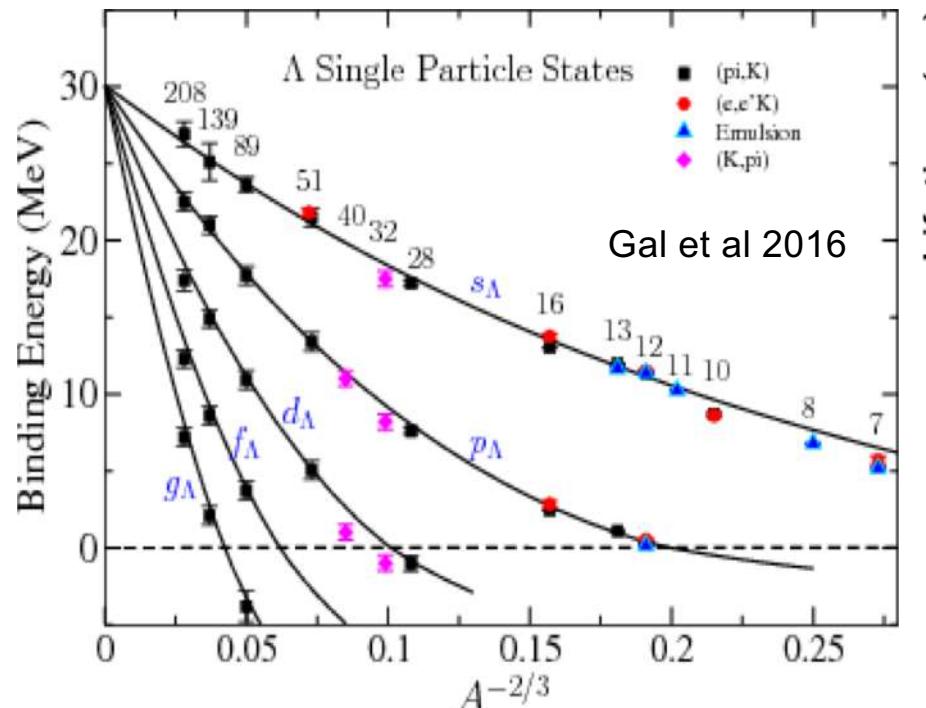


Physics that can be addressed:

- YN and YY interactions
- YN \rightarrow NN weak decay
- Hypernuclear structure

credit: Axel Perez-Obiol

Binding energy of Λ hypernuclei

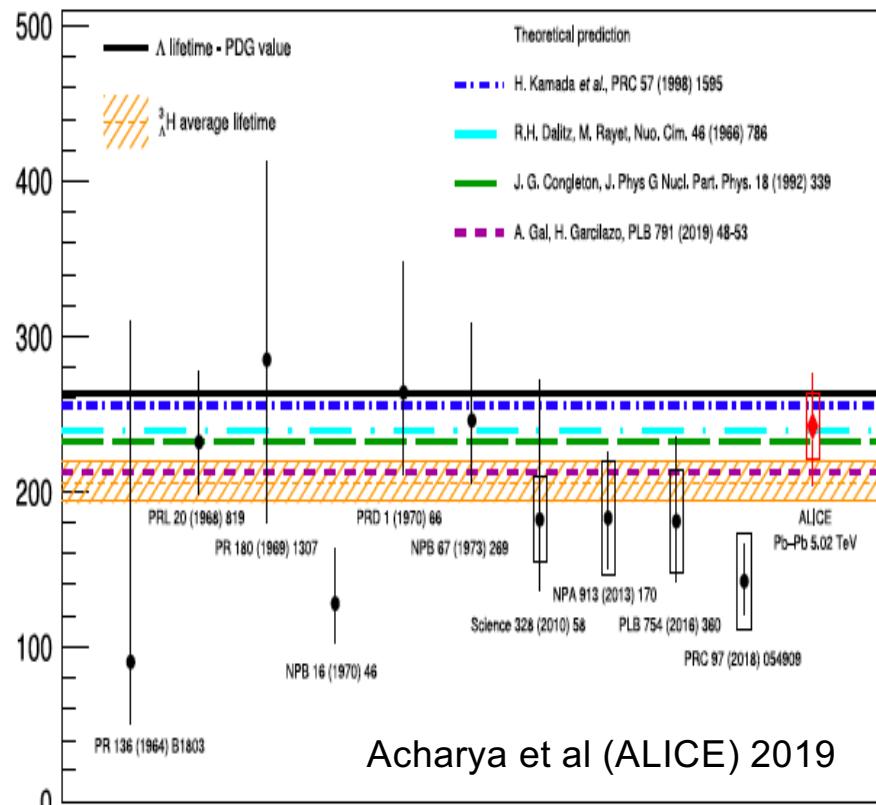


Binding energy of different hypernuclei as function of the mass number

Binding energy saturates at about -30 MeV for large nuclei

Single-particle model reproduces the data quite well Gal et al 2016

Hypertriton lifetime puzzle



Expected $\tau(^3\Lambda H) = \tau(\Lambda)$?
 ⇔ observed: $\tau(^3\Lambda H) < \tau(\Lambda)$?

Conflicting measurements by STAR and ALICE of the hypertriton lifetime triggered the revived experimental and theoretical interest

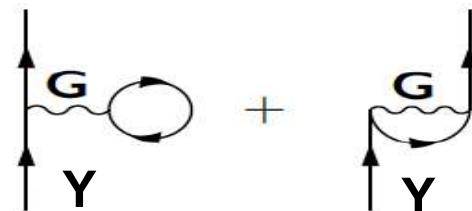
Hyperons in matter

Λ and Σ in dense matter

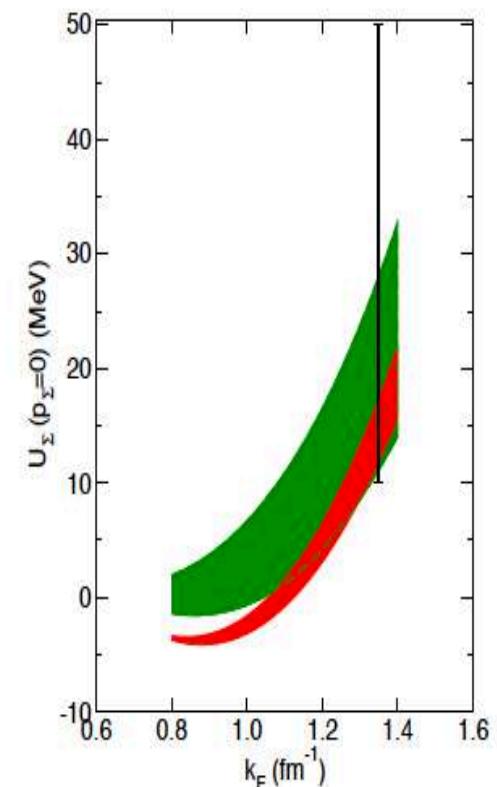
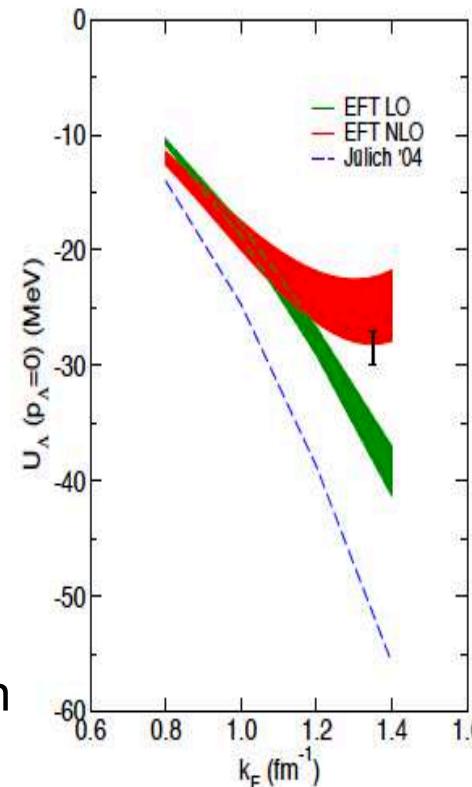
$$k_F = 1.35 \text{ fm}^{-1} (\rho_0 = 0.166 \text{ fm}^{-3})$$

	EFT LO	EFT NLO
Λ [MeV]	550 ... 700	500 ... 650
$U_\Lambda(0)$	-38.0 ... -34.4	-28.2 ... -22.4
$U_\Sigma(0)$	28.0 ... 11.1	17.3 ... 11.9

- Empirical value of Λ binding in nuclear matter $\sim 27\text{-}30$ MeV
- ΣN ($I=3/2$): 3S_1 - 3D_1 decisive for Σ properties in nuclear matter. YN data can be reproduced with attractive and repulsive 3S_1 - 3D_1 interaction. It is chosen to be repulsive in accordance to data on Σ^- atoms and (π^-, K^+) inclusive spectra for Σ^- formation in heavy nuclei.
Lattice* supports repulsion!



$$G = V + V \frac{Q_{\text{pauli}}}{E_0 - H_0} G$$

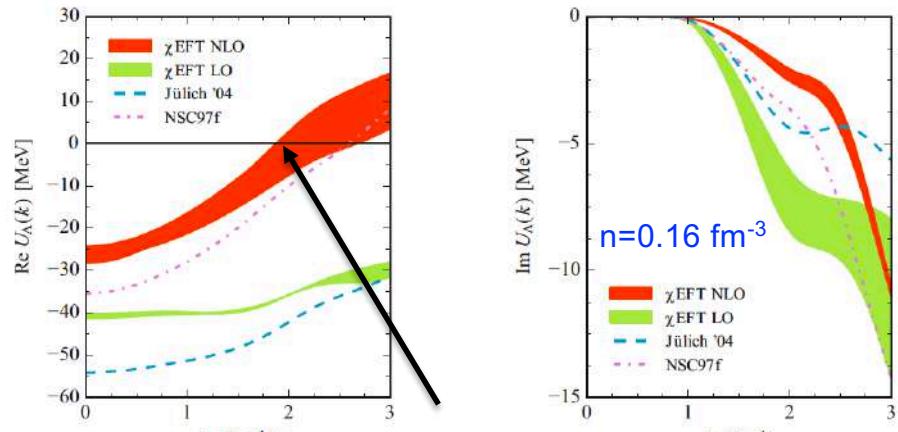


Haidenbauer and Mei  ner , NPA 936 (2015) 29

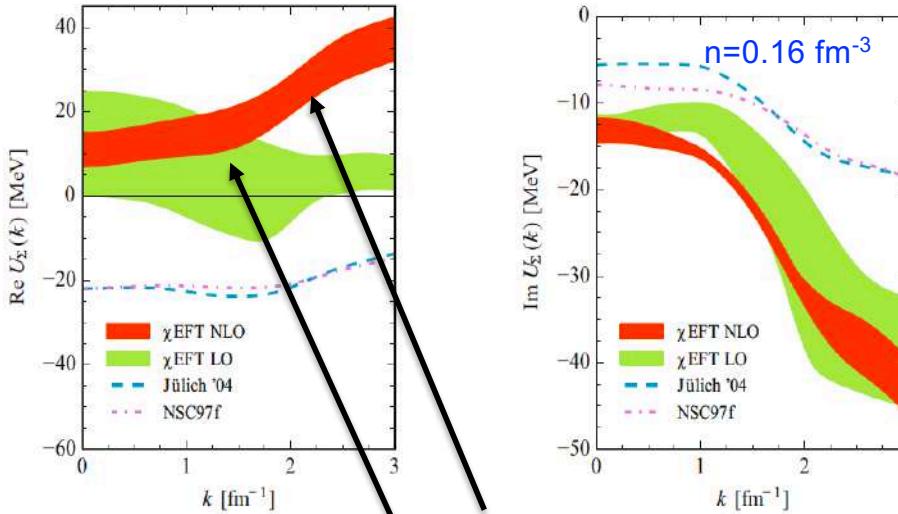
Improving on the calculation by using χ EFT NN interaction and continuous choice in Brueckner-Hartree-Fock approach while investigating isospin-asymmetric matter

S. Petschauer, J. Haidenbauer,
N. Kaiser, U.G. Meißner and
W. Weise EPJA 52 (2016) 15

symmetric nuclear matter

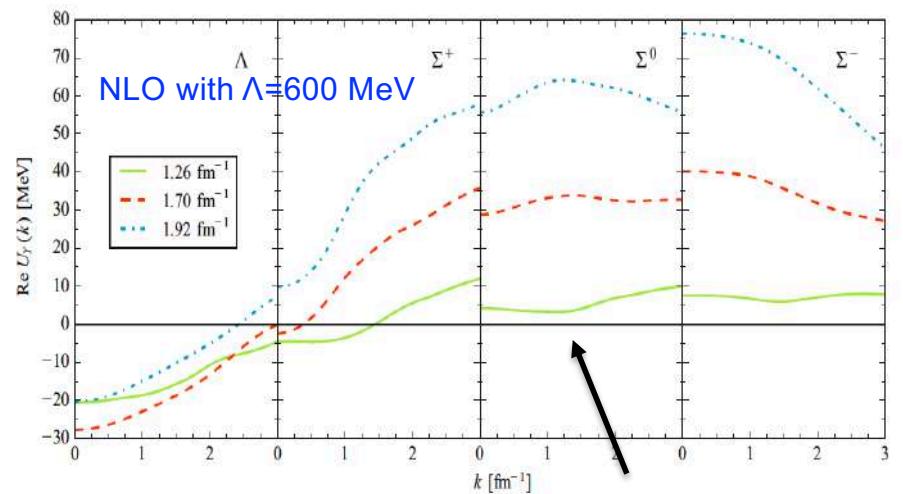


Λ single-particle potential
at NLO turns repulsive $k \sim 2 \text{ fm}^{-1}$

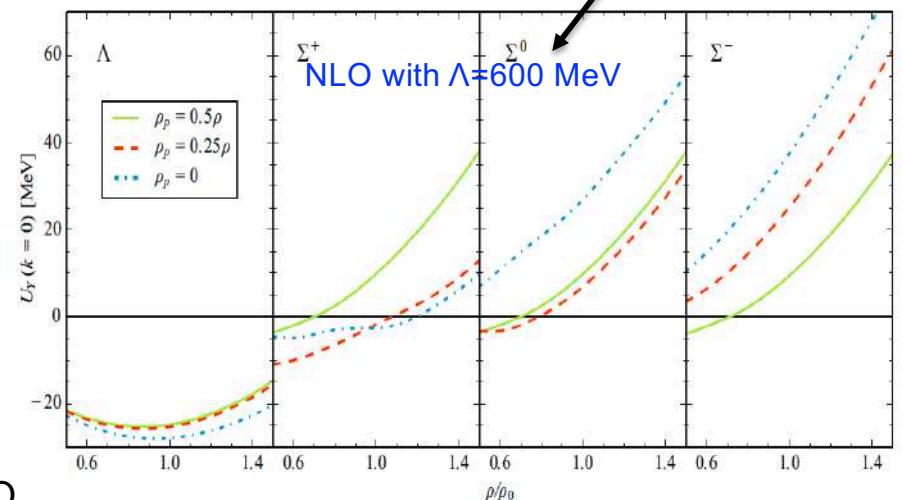


Σ -nuclear potential is moderately repulsive for LO and NLO

neutron matter



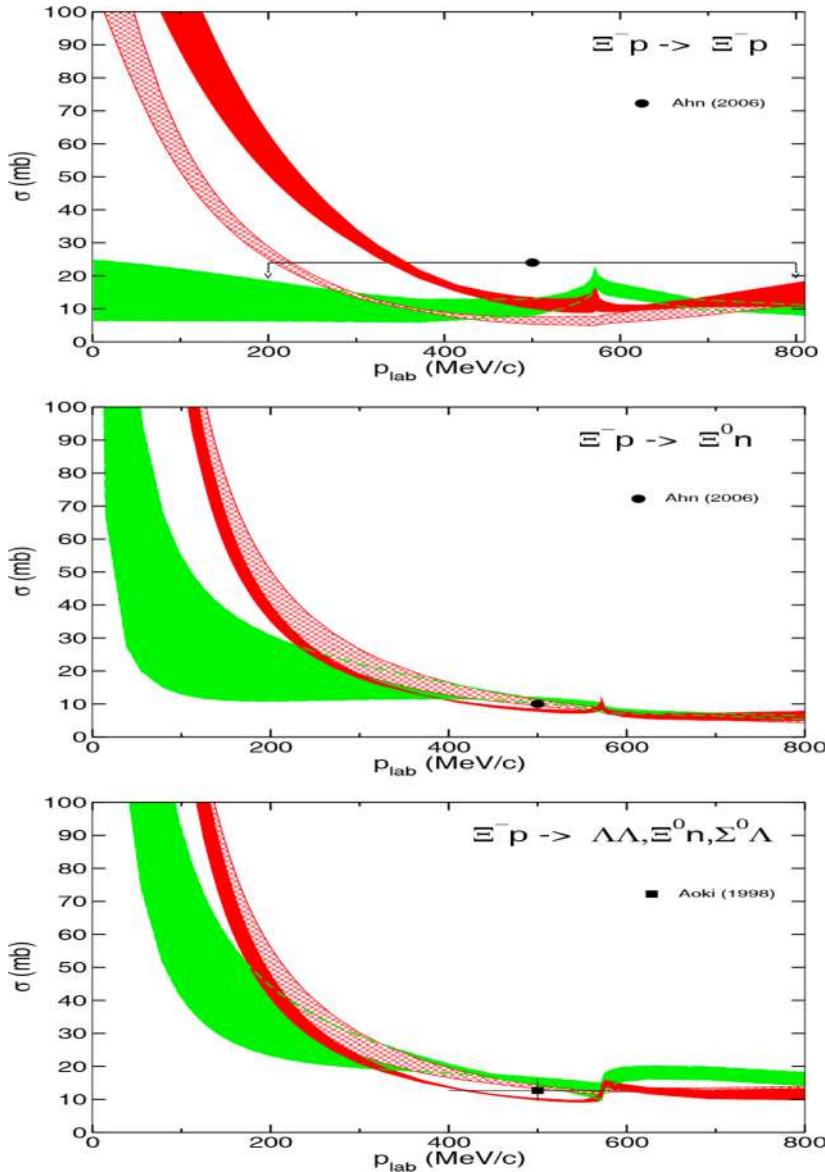
strong isospin dependence of the ΣN interaction



NLO with $\Lambda=600 \text{ MeV}$

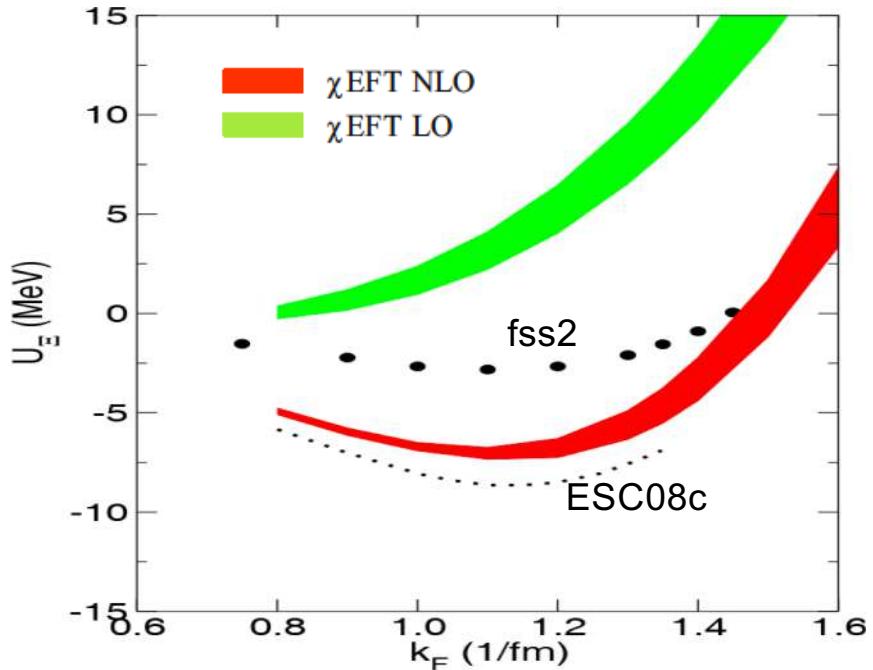
Ξ in dense matter

ΞN cross sections are small



J. Haidenbauer and
U.G. Meißner EPJA 55 (2019) 23

Ξ in dense matter

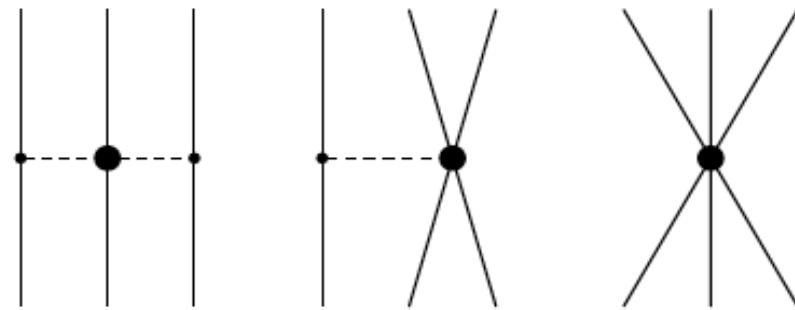


Moderately attractive Ξ -nuclear interaction,
with $U_\Xi(0, k_{F0}) \sim -3$ to -5 MeV.
Smaller than $U_\Xi(n_0) \sim -14$ MeV Khaustov et al'00
and in line with other BHF studies with
phenomenological ΞN potentials

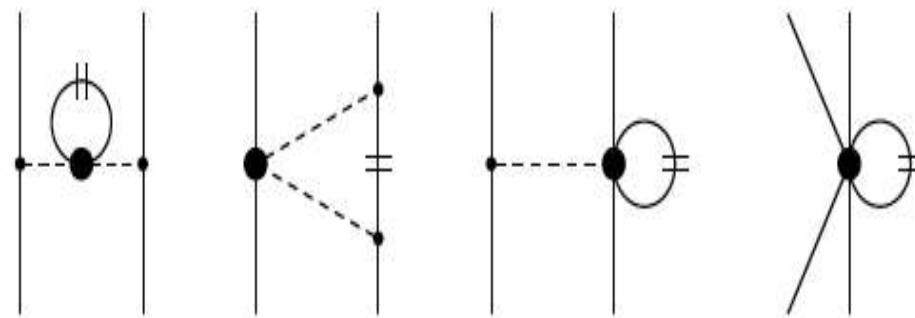
Λ in dense matter: including three-body forces

Three-body forces are required to reproduce few-nucleon binding energies, scattering observables and nuclear saturation in non-relativistic many-body approaches

Three-body force (nominally at N²LO)



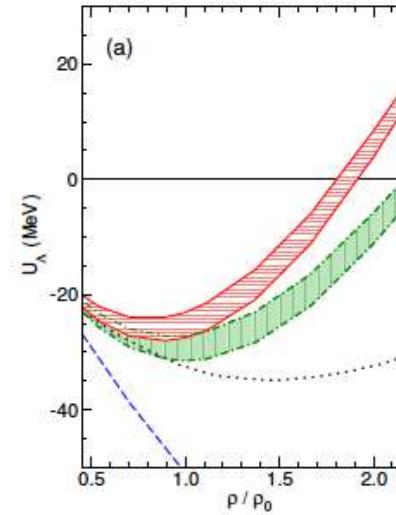
To use it in many-body calculations, such as BHF, one has to construct a density-dependent two-body interaction



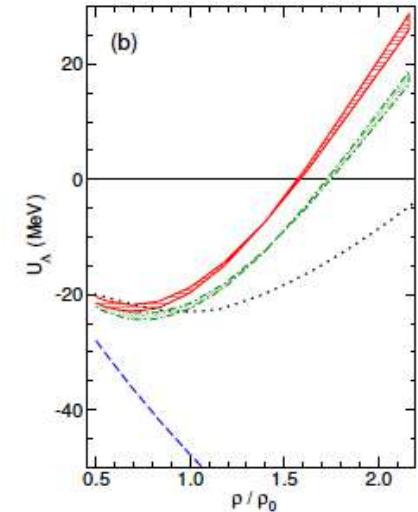
closing two baryon lines summing over the Fermi sea
credit: Haidenbauer

Λ in dense matter

symmetric matter



neutron matter



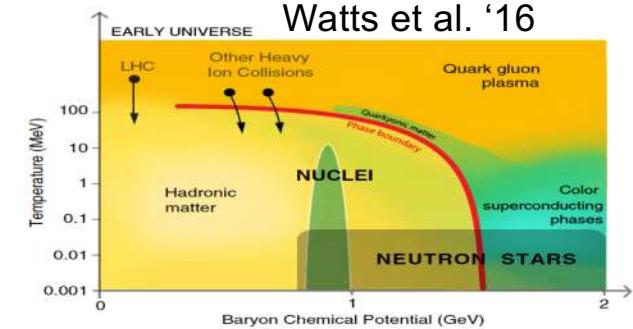
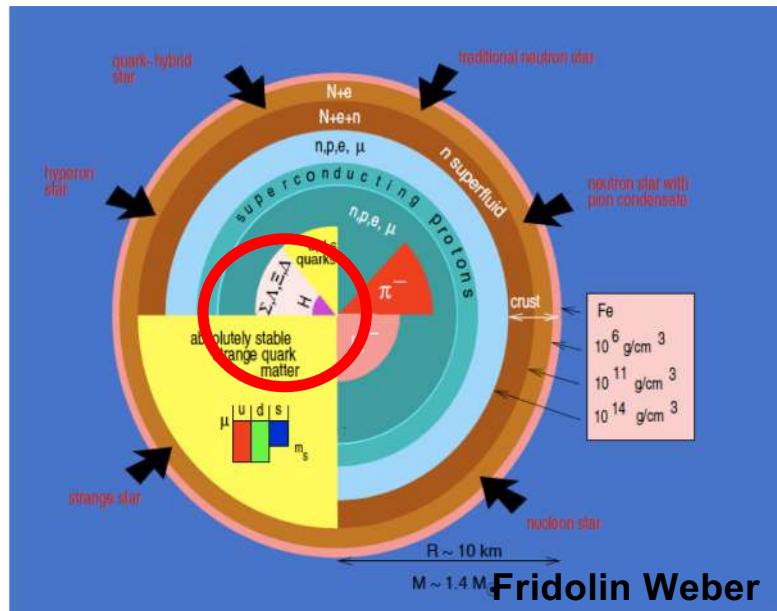
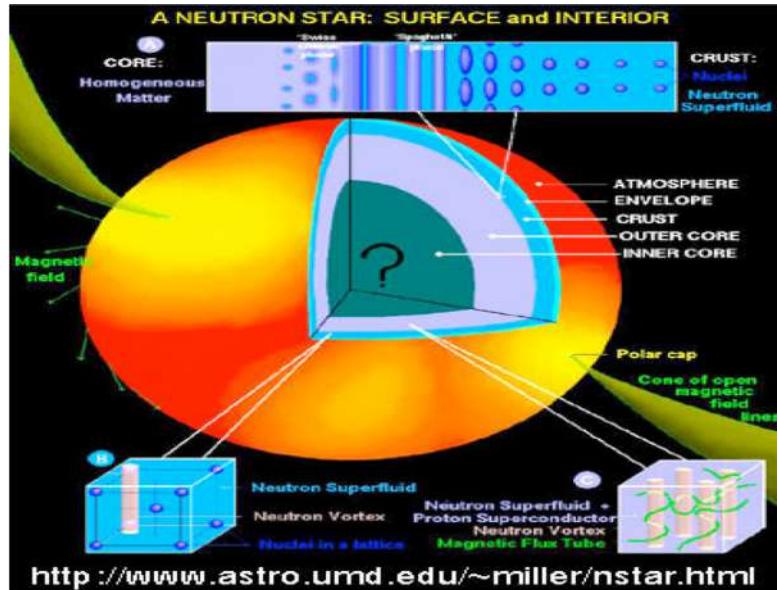
χ EFT gives little attraction or even repulsion for $n > n_0$

In neutron stars, hyperons will appear at high density!!

Solution of the Hyperon Puzzle?

J. Haidenbauer, U.G. Meißner, N. Kaiser and W. Weise EPJA 53 (2017) 121

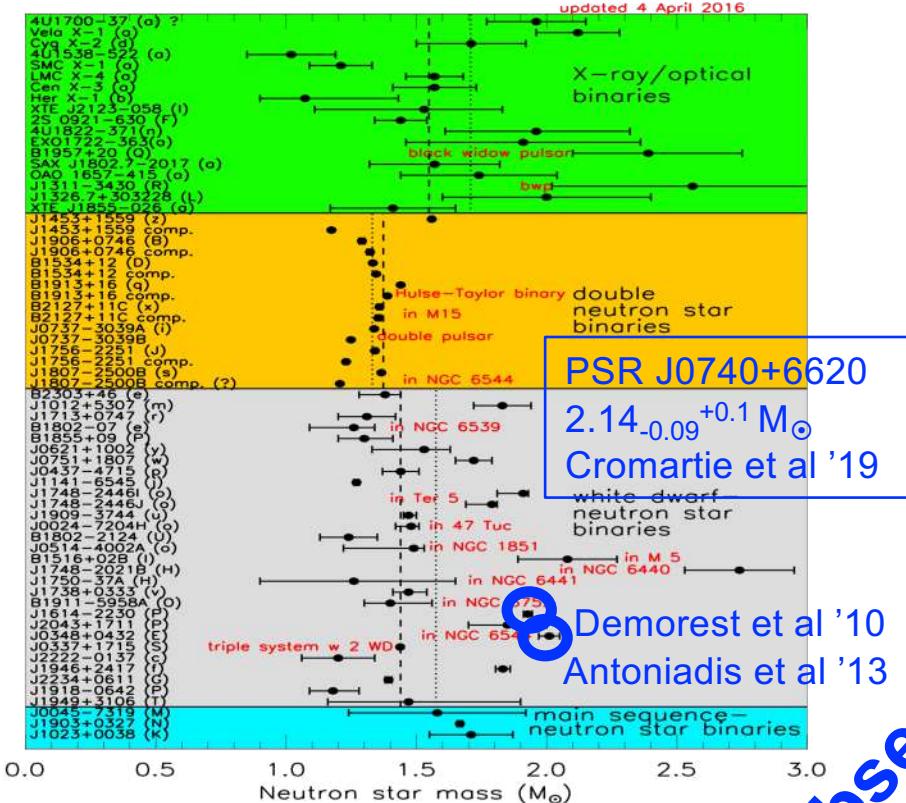
Hyperons and Neutron Stars



- produced in **core collapse supernova** explosions, usually observed as pulsars
- usually refer to compact objects with $M \approx 1\text{-}2 M_\odot$ and $R \approx 10\text{-}12 \text{ Km}$
- extreme densities up to $5\text{-}10 \rho_0$ ($n_0 = 0.16 \text{ fm}^{-3} \Rightarrow \rho_0 = 3 \cdot 10^{14} \text{ g/cm}^3$)
- magnetic field : $B \sim 10^{8\text{-}16} \text{ G}$
- temperature: $T \sim 10^{6\text{-}11} \text{ K}$
- observations: masses, radius (?), gravitational waves, cooling...

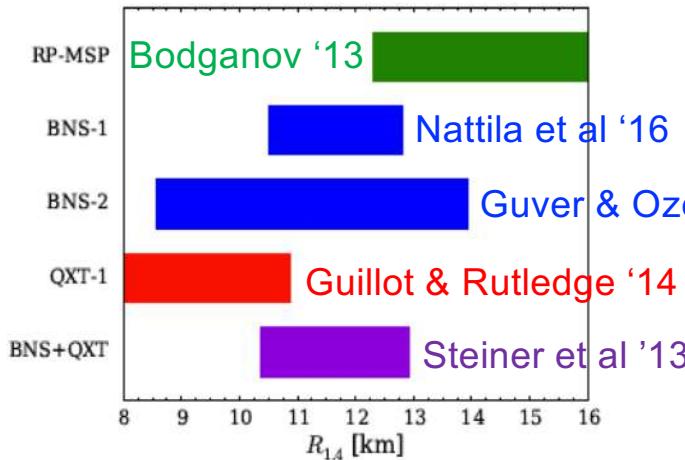
Masses

Lattimer '16

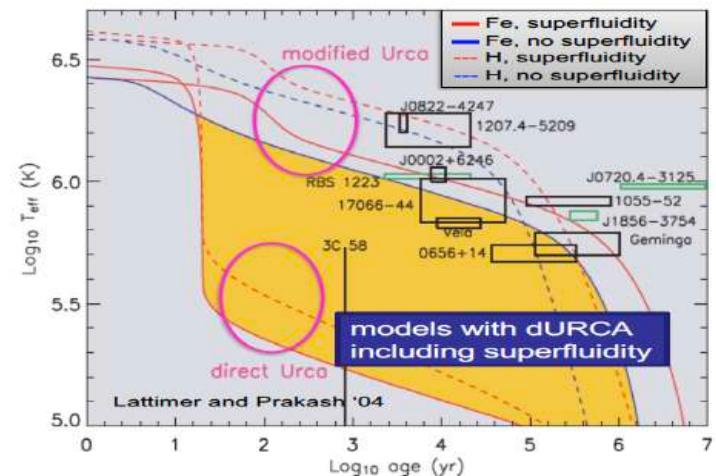


Radius

Fortin et al. '15



Cooling



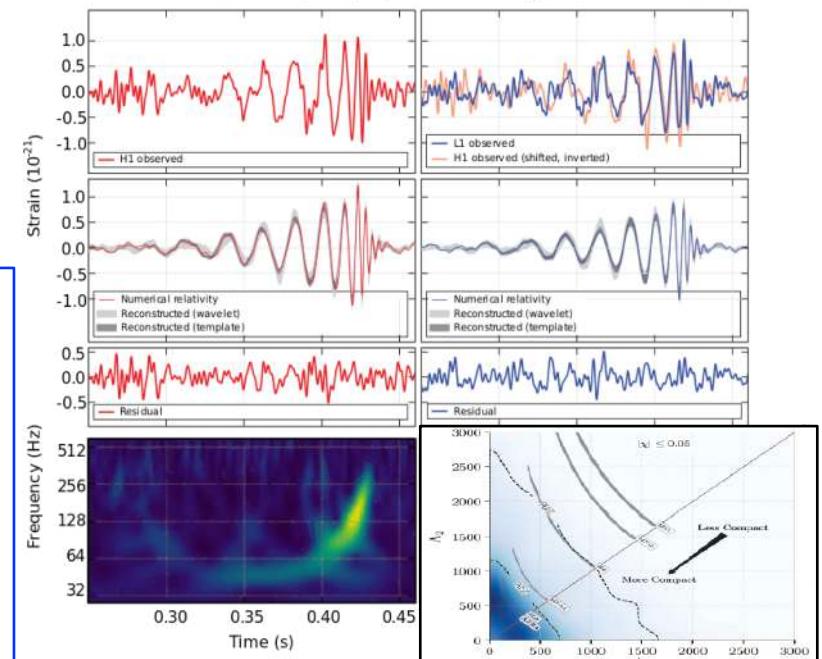
Observations

GW170817

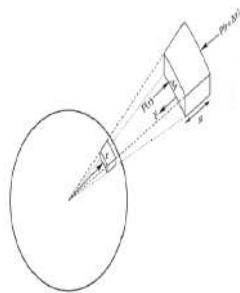
Abbot et al. (LIGO-VIRGO) '17 '18

Hanford, Washington (H1)

Livingston, Louisiana (L1)

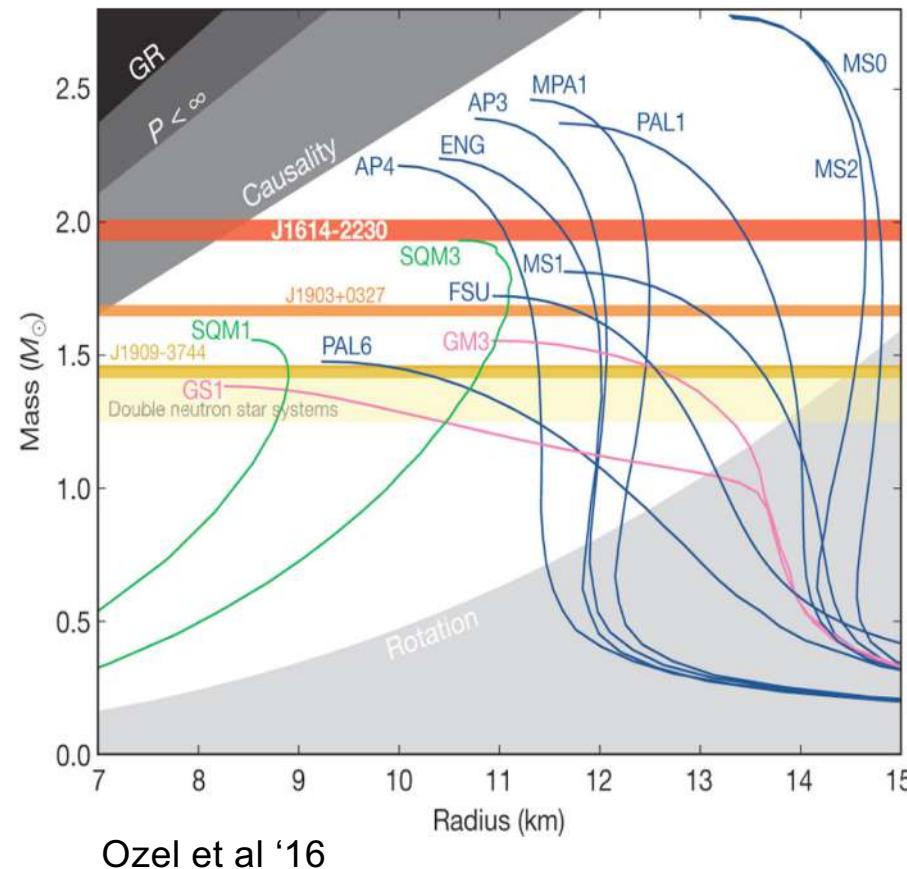


Mass-Radius Relation



$$\frac{dP}{dr} = -\frac{Gm\epsilon}{c^2r^2} \left(1 + \frac{P}{\epsilon}\right) \left(1 + \frac{4\pi r^3 P}{c^2 m}\right) \left(1 - \frac{2Gm}{c^2 r}\right)^{-1}$$

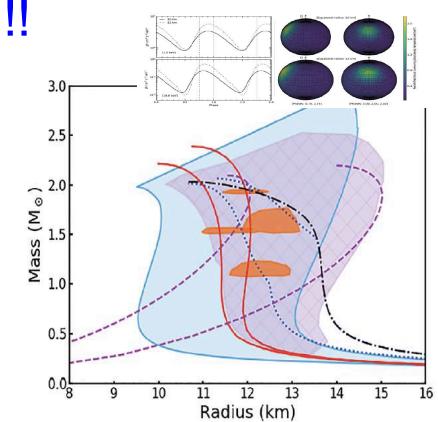
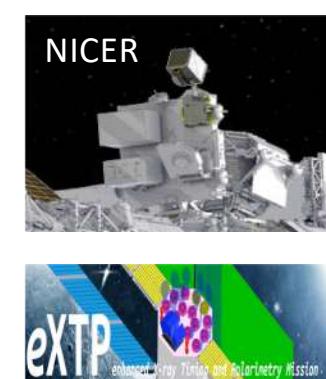
$$\frac{dm}{dr} = \frac{4\pi r^2 \epsilon}{c^2}$$



- primary ingredient:
EoS: $\epsilon(n)$, $P(n)$, $P(\epsilon)$
in charge neutral β -stable matter

- some constraints:
 - Schwarzschild limit (GR)
 $R \geq 2 GM/c^2$
 - causality limit for EoS
 $R \geq 2.9 GM/c^2$
 - mass-shedding limit
 $R < (GM/2\pi)^{1/3}/\nu^{2/3}$

Need of simultaneous
mass-radius measurements to
constrain EoS !!!



Watts et al. (LT) '19

The Nucleonic Equation of State

The Equation of State (EoS) is a relation between thermodynamic variables describing the state of matter

Microscopic Ab-initio Approaches:

based on solving the many-body problem starting from two- and three-body interactions

- *Variational method: APR, CBF,..*
- *Quantum Montecarlo : AFDMC..*
- *Coupled cluster expansion*
- *Diagrammatic: BBG (BHF), SCGF..*
- *Relativistic DBHF*
- *RG methods: SRG from χ EFT..*
- *Lattice methods*

Advantage: systematic addition of higher-order contributions

Disadvantage: applicable up to? (*SRG from χ EFT $\sim 1-2 n_0$*)

Phenomenological Approaches:

based on density-dependent interactions adjusted to nuclear observables and neutron star observations

- *Non-relativistic EDF: Skyrme..*
- *Relativistic Mean-Field (RMF) and Relativistic Hartree-Fock (RHF)*
- *Liquid Drop Model: BPS, BBP,..*
- *Thomas-Fermi model: Shen*
- *Statistical Model: HWN, RG, HS..*

Advantage: applicable to high densities beyond n_0

Disadvantage: not systematic

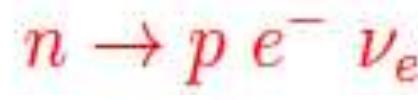
What about Hyperons?

credit: Vidana

First proposed in 1960 by Ambartsumyan & Saakyan

Hyperon	Quarks	I(J ^P)	Mass (MeV)
Λ	uds	0(1/2 ⁺)	1115
Σ^+	uus	1(1/2 ⁺)	1189
Σ^0	uds	1(1/2 ⁺)	1193
Σ^-	dds	1(1/2 ⁺)	1197
Ξ^0	uss	1/2(1/2 ⁺)	1315
Ξ^-	dss	1/2(1/2 ⁺)	1321
Ω^-	sss	0(3/2 ⁺)	1672

Traditionally neutron stars were modeled by a uniform fluid of neutron rich matter in β -equilibrium



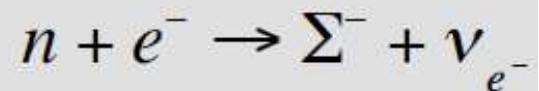
but more exotic degrees of freedom are expected, such as **hyperons**, due to:

- high value of density at the center and
- the rapid increase of the nucleon chemical potential with density

Hyperons might be present at $n \sim (2-3)n_0$!!!

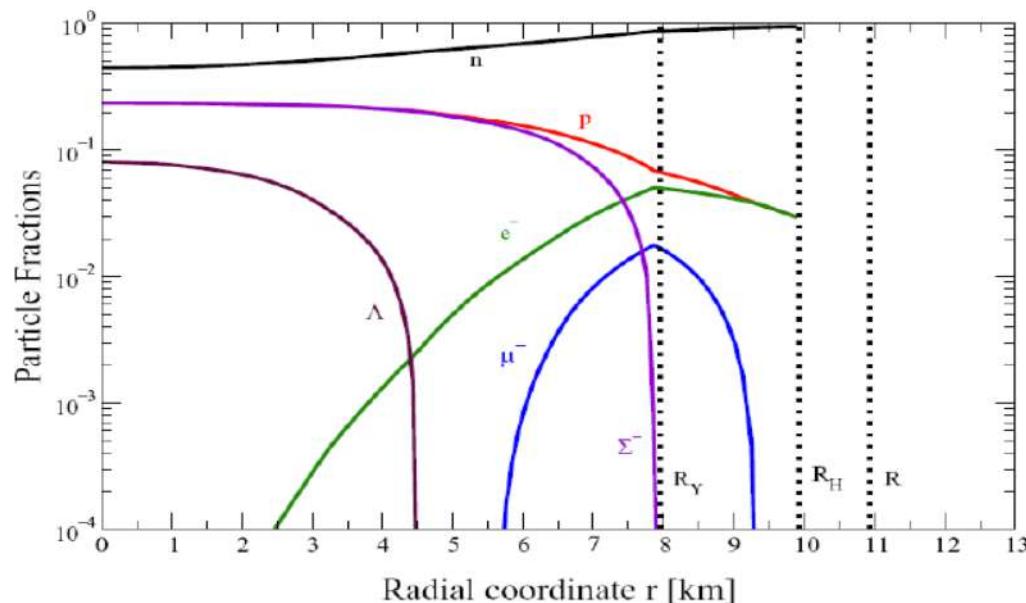
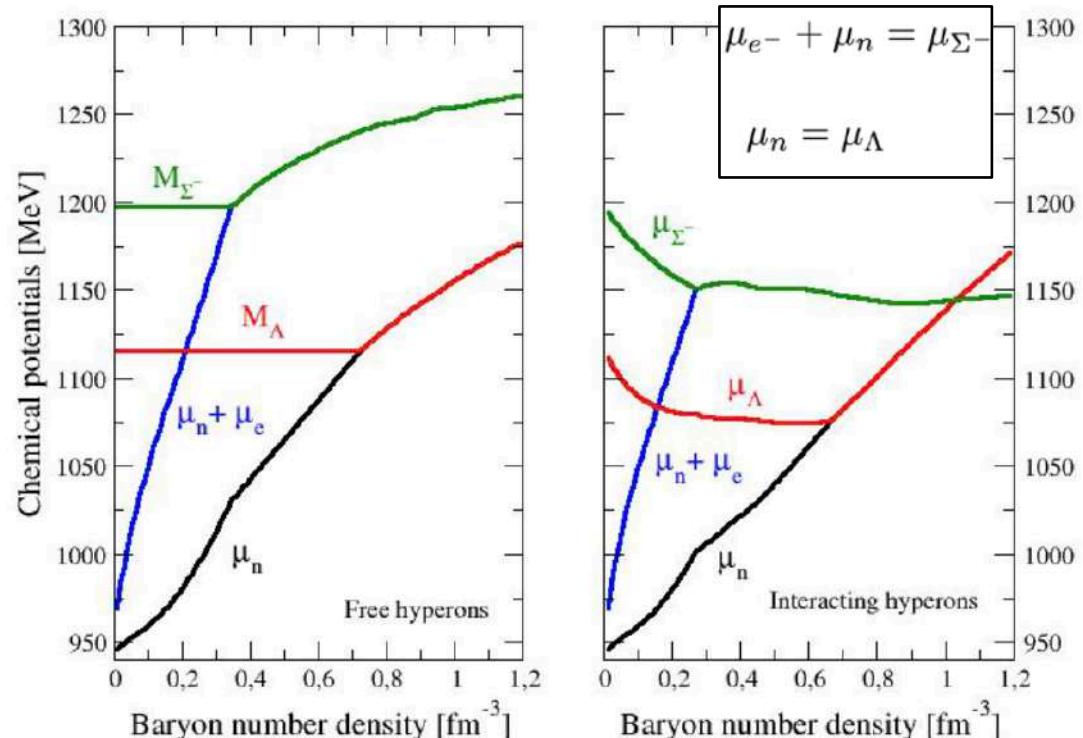
β -stable hyperonic matter

μ_N is large enough to make N->Y favorable



$$\mu_i = b_i \mu_n - q_i \mu_e$$

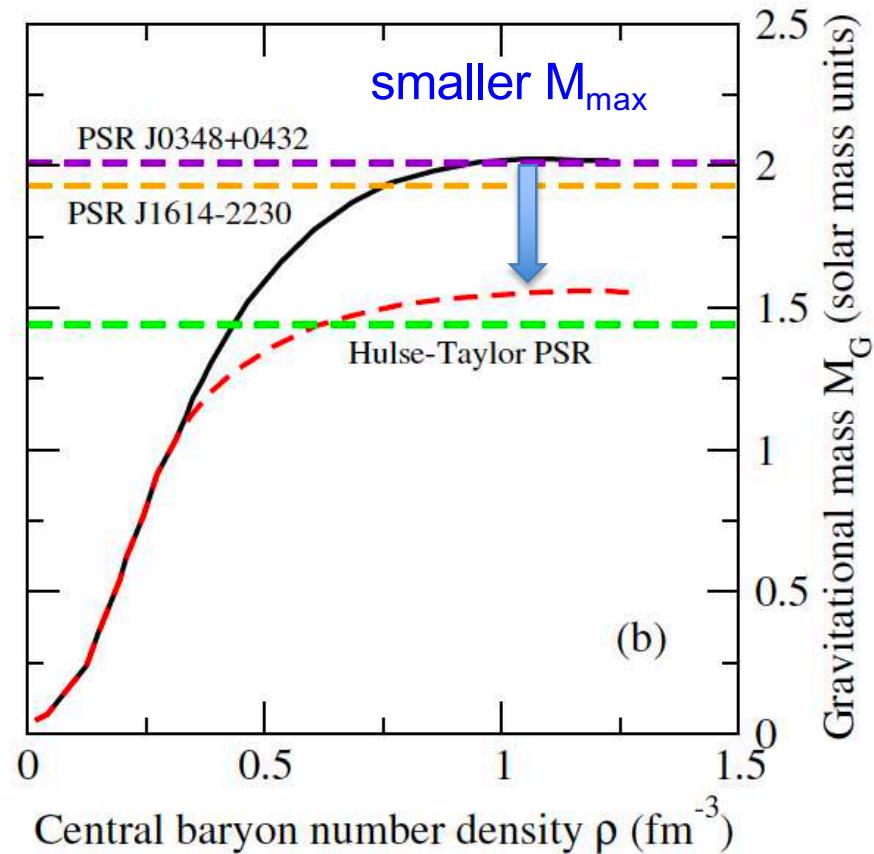
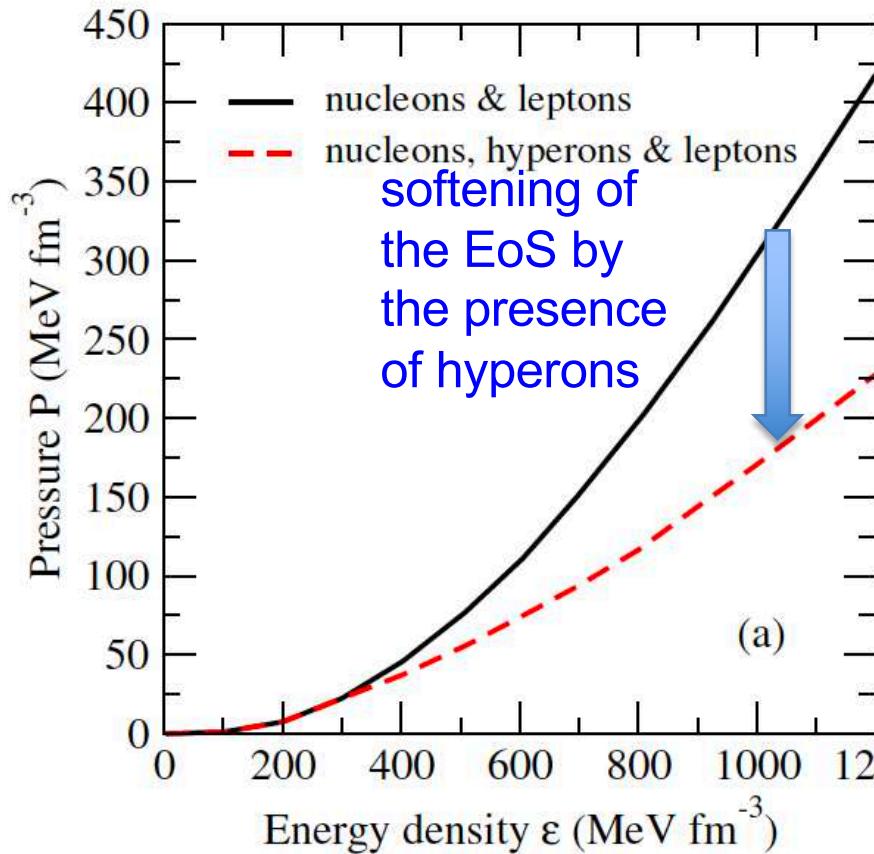
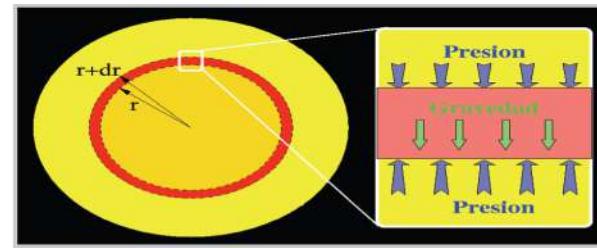
$$\sum_i x_i q_i = 0$$



credit: Vidana

Inclusion of hyperons....

Credit:
Dani P. Page



..... induces a strong softening of the EoS
that leads to $M_{\max} < 2M_{\odot}$

Chatterjee and Vidana '16
Vidana '18

The Hyperon Puzzle



The Hyperon Puzzle



Scarce experimental information:

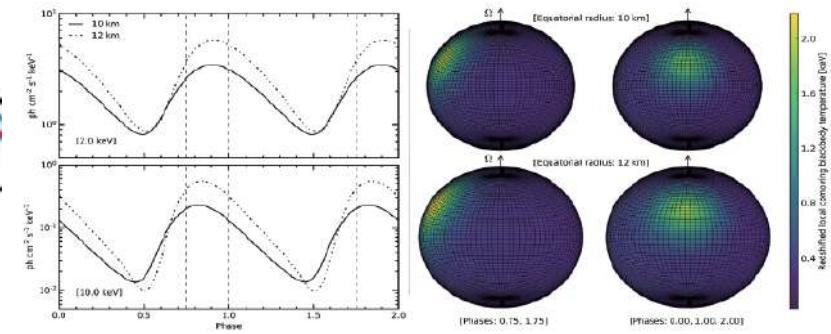
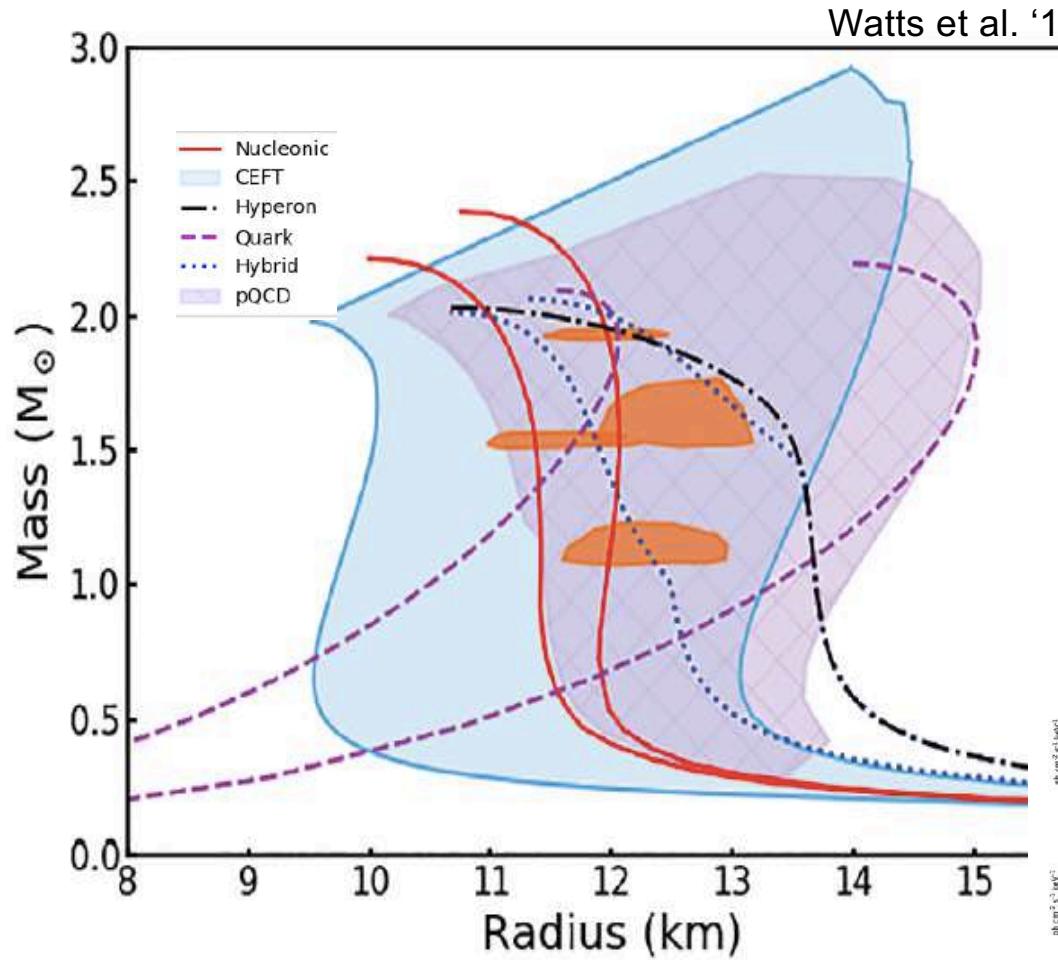
- data from several single Λ - and few Ξ - hypernuclei, and few double Λ hypernuclei
- few YN scattering data (~ 50 points) due to difficulties in preparing hyperon beams and no hyperon targets available
- YN data from femtoscopy

The presence of hyperons in neutron stars is energetically probable as density increases. However, it induces a strong softening of the EoS that leads to **maximum neutron star masses $< 2M_{\odot}$**

Solution?

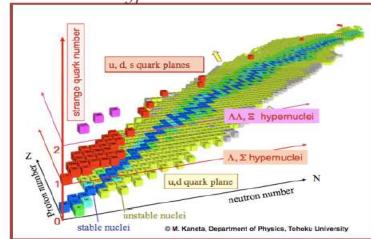
- stiffer YN and YY interactions
- hyperonic 3-body forces
- push of Y onset by Δ -isobars or meson condensates
- quark matter below Y onset
- dark matter, modified gravity theories...

Future: space missions to study the interior of NS

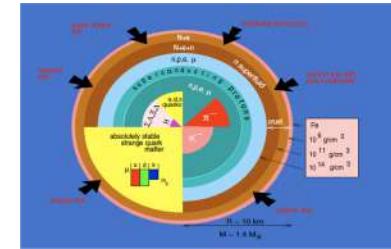


Constraints from pulse profile modelling of rotation-powered pulsars with eXTP

and multimessenger astronomy



Present and Future



A lot of experimental, observational and theoretical effort has been invested to understand **hyperons in nuclei and neutron stars**

Hyperon-nucleon and hyperon-hyperon interactions are crucial for hypernuclear physics and the physics of compact objects, such neutron stars

Neutron stars provide a unique scenario for testing hyperons at extreme densities

The **future** of hyperon physics relies on **particle** and nuclear experiments as well as X-ray and multimessenger astronomy

